

by V. Pomjakushin (2014)

Maple script for calculations related to the practicum (<http://sing.web.psi.ch/sing/instr/hrpt/praktikum.html>). Also includes dot-product and cross-product functions in anybasis.

> with(LinearAlgebra) :

▼ Metric tensor $G1_{\{ij\}}$

Scalar product should be defined through the metric tensor G .

> $G := \langle a[1], a[2], a[3] \rangle \cdot \langle | \rangle (a[1], a[2], a[3])$;

$$G := \begin{bmatrix} a_1^2 & a_1 a_2 & a_1 a_3 \\ a_1 a_2 & a_2^2 & a_2 a_3 \\ a_1 a_3 & a_2 a_3 & a_3^2 \end{bmatrix}$$

(1.1)

> $m_{1,1} := a_1 a_1$:

$m_{2,1} := a_1 a_2 \cos(\alpha_3)$:

$m_{3,1} := a_1 a_3 \cos(\alpha_2)$:

$m_{2,2} := a_2 a_2$:

$m_{3,2} := a_3 a_2 \cos(\alpha_1)$:

$m_{3,3} := a_3 a_3$:

$G1 := Matrix(3, 3, \{(1, 1) = m_{1,1}, (1, 2) = m_{2,1}, (1, 3) = m_{3,1}, (2, 1) = m_{2,1}, (2, 2) = m_{2,2}, (2, 3) = m_{3,2}, (3, 1) = m_{3,1}, (3, 2) = m_{3,2}, (3, 3) = m_{3,3}\})$;

(1.2)

$$G1 := \begin{bmatrix} a_1^2 & a_1 a_2 \cos(\alpha_3) & a_1 a_3 \cos(\alpha_2) \\ a_1 a_2 \cos(\alpha_3) & a_2^2 & a_3 a_2 \cos(\alpha_1) \\ a_1 a_3 \cos(\alpha_2) & a_3 a_2 \cos(\alpha_1) & a_3^2 \end{bmatrix} \quad (1.2)$$

Volume² is det(G1).

$$\begin{aligned} > Vol := \text{RealDomain}:-\text{simplify}\left(\text{Determinant}(G1)^{\frac{1}{2}}\right); \\ Vol := \sqrt{1 + 2 \cos(\alpha_1) \cos(\alpha_3) \cos(\alpha_2) - \cos(\alpha_1)^2 - \cos(\alpha_3)^2 - \cos(\alpha_2)^2} |a_1 a_2 a_3| \end{aligned} \quad (1.3)$$

We define the dual basis metric tensor

$$\begin{aligned} > G2 := \text{MatrixInverse}(G1) : \\ & \text{unassign('i'); unassign('j', k);} \\ & d2 := \text{add}(\text{add}((G2[i, j] \cdot h[i] \cdot h[j]), j = 1 .. 3), i = 1 .. 3) : \\ & \text{simplify}(d2); \end{aligned} \quad (1.4)$$

$$\begin{aligned} & \left(\cos(\alpha_2)^2 a_1^2 a_3^2 h_2^2 - 2 \cos(\alpha_2) \cos(\alpha_3) a_1^2 a_2 a_3 h_2 h_3 - 2 \cos(\alpha_2) \cos(\alpha_1) a_1 a_2 a_3^2 h_1 h_2 + \cos(\alpha_3)^2 a_1^2 a_2^2 h_3^2 \right. \\ & \quad \left. - 2 \cos(\alpha_3) \cos(\alpha_1) a_1 a_2^2 a_3 h_1 h_3 + \cos(\alpha_1)^2 a_2^2 a_3^2 h_1^2 + 2 \cos(\alpha_2) a_1 a_2^2 a_3 h_1 h_3 + 2 \cos(\alpha_3) a_1 a_2 a_3^2 h_1 h_2 + 2 \cos(\alpha_1) \right. \\ & \quad \left. a_1^2 a_2 a_3 h_2 h_3 - h_3^2 a_1^2 a_2^2 - h_2^2 a_1^2 a_3^2 - h_1^2 a_2^2 a_3^2 \right) / \left(a_1^2 a_2^2 a_3^2 \left(-2 \cos(\alpha_1) \cos(\alpha_3) \cos(\alpha_2) + \cos(\alpha_2)^2 + \cos(\alpha_3)^2 + \cos(\alpha_1)^2 \right. \right. \\ & \quad \left. \left. - 1 \right) \right) \end{aligned}$$

Intro

Praktikum for neutron diffraction. Determination of long range antiferromagnetic order by powder neutron diffraction.

From the analysis of the nuclear and magnetic Bragg peak intensities and positions we will verify the crystal and magnetic structures of manganese sulfide and determine the size of the magnetic moment on manganese.

Integral intensity for powder Bragg scattering.

1. Lorentz factor

$$L := \theta \mapsto \frac{1}{\sin(\theta) \sin(2\theta)}$$

$$L := \theta \mapsto \frac{1}{\sin(\theta) \sin(2\theta)} \quad (2.1)$$

2. Lattice sum over Bravais lattice. It is important for calculations only if one uses different elementary cell for magnetic and nuclear structure factor calculations.

$$LS := \frac{N(2\pi)^3}{v_0}$$

$$LS := \frac{8N\pi^3}{v_0} \quad (2.2)$$

Define zeroth cell with centering translations & the lattice

Define F-centering translations

$$\begin{aligned} \#q &:= \langle | \rangle (h, k, l); \\ cf[1] &:= \langle | \rangle (0, 0, 0); \\ cf[2] &:= \langle | \rangle \left(\frac{1}{2}, \frac{1}{2}, 0 \right); \\ cf[3] &:= \langle | \rangle \left(\frac{1}{2}, 0, \frac{1}{2} \right); \end{aligned}$$

$$cf[4] := \langle | \rangle \left(0, \frac{1}{2}, \frac{1}{2} \right);$$

$$cf_1 := \begin{bmatrix} 0 & 0 & 0 \end{bmatrix}$$

$$cf_2 := \begin{bmatrix} \frac{1}{2} & \frac{1}{2} & 0 \end{bmatrix}$$

$$cf_3 := \begin{bmatrix} \frac{1}{2} & 0 & \frac{1}{2} \end{bmatrix}$$

$$cf_4 := \begin{bmatrix} 0 & \frac{1}{2} & \frac{1}{2} \end{bmatrix}$$

(3.1)

2theta from H

$$> G1cube := RealDomain:-simplify \left(subs \left(\left[a_1 = c, \alpha_1 = \frac{\pi}{2} \right], \left[a_2 = c, \alpha_2 = \frac{\pi}{2} \right], \left[a_3 = c, \alpha_3 = \frac{\pi}{2} \right], G1 \right) \right);$$

$$d2cube := RealDomain:-simplify \left(subs \left(\left[a_1 = c, \alpha_1 = \frac{\pi}{2} \right], \left[a_2 = c, \alpha_2 = \frac{\pi}{2} \right], \left[a_3 = c, \alpha_3 = \frac{\pi}{2} \right], d2^{\frac{1}{2}} \right) \text{ assuming } c > 0 \right);$$

$$G1cube := \begin{bmatrix} c^2 & 0 & 0 \\ 0 & c^2 & 0 \\ 0 & 0 & c^2 \end{bmatrix}$$

$$d2cube := \frac{\sqrt{h_1^2 + h_2^2 + h_3^2}}{|c|}$$

(3.2)

$$> qvec := 2 \cdot \text{Pi} \cdot d2cube;$$

$$\text{evalf}\left(\frac{2 \cdot \text{Pi}}{5.209}\right);$$

$$qvec := \frac{2 \pi \sqrt{h_1^2 + h_2^2 + h_3^2}}{|c|} \\ 1.206217183 \quad (3.3)$$

Bragg law

$$> \text{Bragg} := 2 \cdot \sin(\theta) = \lambda \cdot d2cube;$$

$$\text{Bragg} := 2 \sin(\theta) = \frac{\lambda \sqrt{h_1^2 + h_2^2 + h_3^2}}{|c|} \quad (3.4)$$

theta as a function of H.

$$> t2 := \text{solve}(\text{Bragg}, \theta);$$

$$t2 := (\text{subs}([\lambda = 1.886, c = 5.209], \text{solve}(\text{Bragg}, \theta)));$$

$$t2 := \arcsin\left(\frac{\lambda \sqrt{h_1^2 + h_2^2 + h_3^2}}{2 |c|}\right)$$

$$t2 := \arcsin\left(\frac{0.9430000000 \sqrt{h_1^2 + h_2^2 + h_3^2}}{|5.209|}\right) \quad (3.5)$$

qvec as a ...

$$> qvec := \text{evalf}\left(\text{subs}\left([\lambda = 1.886], \frac{4 \cdot \pi \cdot \sin(t2)}{\text{lambda}}\right)\right);$$

$$qvec := 1.206217183 \sqrt{h_1^2 + h_2^2 + h_3^2} \quad (3.6)$$

Set the atoms

Set up Mn atoms, both positions and spins

> for j from 1 to 4 do

$b[j] := b_{Mn} :$

$r[j] := \langle | \rangle (0, 0, 0) + cf[j];$

$s[j] := \langle | \rangle (s0, -s0, 0) \cdot \exp\left(2 \cdot \text{Pi} \cdot I \cdot \left\langle \frac{1}{2} \middle| \frac{1}{2} \middle| \frac{1}{2} \right\rangle \cdot r[j]\right);$

end do; unassign('j');

$$b_1 := b_{Mn}$$

$$r_1 := \begin{bmatrix} 0 & 0 & 0 \end{bmatrix}$$

$$s_1 := \begin{bmatrix} s0 & -s0 & 0 \end{bmatrix}$$

$$b_2 := b_{Mn}$$

$$r_2 := \begin{bmatrix} \frac{1}{2} & \frac{1}{2} & 0 \end{bmatrix}$$

$$s_2 := \begin{bmatrix} -s0 & s0 & 0 \end{bmatrix}$$

$$b_3 := b_{Mn}$$

$$r_3 := \begin{bmatrix} \frac{1}{2} & 0 & \frac{1}{2} \end{bmatrix}$$

$$s_3 := \begin{bmatrix} -s0 & s0 & 0 \end{bmatrix}$$

$$b_4 := b_{Mn}$$

$$r_4 := \begin{bmatrix} 0 & \frac{1}{2} & \frac{1}{2} \end{bmatrix}$$

$$s_4 := \begin{bmatrix} -s0 & s0 & 0 \end{bmatrix}$$

(4.1)

=> Magnetic form-factor coeff. (A,a,B,b,C,c,D)

MMN3 0.4198 14.2829 0.6054 5.4689 0.9241 -0.0088 -0.9498

> MMN3 := ⟨0.4198, 14.2829, 0.6054, 5.4689, 0.9241, -0.0088, -0.9498, 0⟩ :

From international tables mMn(q) ff of Mn3+ of $q = \frac{4\pi \sin(\theta)}{\lambda}$

> #j0 := (s) → A·exp(-a·s²) + B·exp(-b·s²) + C·exp(-c·s²) + D;

j0 := (s, A :: Vector) → add(A[i]·exp(-A[i+1]·s²), i = seq(1 .. 7, 2));

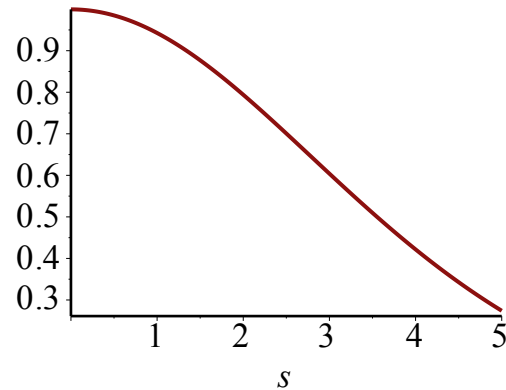
mMn := (s) → j0($\frac{s}{4 \cdot \text{Pi}}$, MMN3);

$$j0 := (s, A::Vector) \mapsto \text{add}\left(A_i e^{-A_{i+1} s^2}, i = \text{seq}(1..7, 2)\right)$$

$$mMn := s \mapsto j0\left(\frac{s}{4\pi}, MMN3\right)$$

(4.2)

> plot(mMn(s), s = 0 .. 5);



Set up S-atoms

> **for** j **from** 1 **to** 4 **do**

$b[j + 4] := b_S :$

$r[j + 4] := \langle | \rangle \left(\frac{1}{2}, 0, 0 \right) + cf[j];$

end do: *unassign('j');*

eval(r);

$$\text{table} \left(\left[1 = \begin{bmatrix} 0 & 0 & 0 \end{bmatrix}, 2 = \begin{bmatrix} \frac{1}{2} & \frac{1}{2} & 0 \end{bmatrix}, 3 = \begin{bmatrix} \frac{1}{2} & 0 & \frac{1}{2} \end{bmatrix}, 4 = \begin{bmatrix} 0 & \frac{1}{2} & \frac{1}{2} \end{bmatrix}, 5 = \begin{bmatrix} \frac{1}{2} & 0 & 0 \end{bmatrix}, 6 = \begin{bmatrix} 1 & \frac{1}{2} & 0 \end{bmatrix}, 7 = \begin{bmatrix} 1 & 0 & \frac{1}{2} \end{bmatrix}, 8 \right. \right. \\ \left. \left. = \begin{bmatrix} \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix} \right] \right) \quad (4.3)$$

▼ Nuclear intensity and peak positions

▼ **First calculate numerically, in 10^{-12} [cm]**

#q := (<, > (h, k, l)) %T

Nuclear structure factor

> F := (q) → add(b[j]·exp(2·Pi·I·(r[j]·q)), j = 1 .. 8);
 eval(F(<h, k, l>));
 factor(%);

$$F := q \mapsto \text{add}\left(b_j e^{2i\pi(\text{Typesetting: } -\text{delayDotProduct}(r_j, q))}, j = 1 \dots 8\right)$$

$$b_{Mn} + b_{Mn} e^{2i\pi\left(\frac{h}{2} + \frac{k}{2}\right)} + b_{Mn} e^{2i\pi\left(\frac{h}{2} + \frac{l}{2}\right)} + b_{Mn} e^{2i\pi\left(\frac{k}{2} + \frac{l}{2}\right)} + b_S e^{i\pi h} + b_S e^{2i\pi\left(h + \frac{k}{2}\right)} + b_S e^{2i\pi\left(h + \frac{l}{2}\right)}$$

$$+ b_S e^{2i\pi\left(\frac{h}{2} + \frac{k}{2} + \frac{l}{2}\right)}$$

$$b_{Mn} + b_{Mn} e^{i\pi(h+k)} + b_{Mn} e^{i\pi(h+l)} + b_{Mn} e^{i\pi(k+l)} + b_S e^{i\pi h} + b_S e^{i\pi(2h+k)} + b_S e^{i\pi(2h+l)} + b_S e^{i\pi(h+k+l)} \quad (5.1.1)$$

$$b_{Mn} + b_{Mn} e^{i\pi(h+k)} + b_{Mn} e^{i\pi(h+l)} + b_{Mn} e^{i\pi(k+l)} + b_S e^{i\pi h} + b_S e^{i\pi(2h+k)} + b_S e^{i\pi(2h+l)} + b_S e^{i\pi(h+k+l)} \quad (5.1.2)$$

$$b_{Mn} + b_{Mn} e^{i\pi(h+k)} + b_{Mn} e^{i\pi(h+l)} + b_{Mn} e^{i\pi(k+l)} + b_S e^{i\pi h} + b_S e^{i\pi(2h+k)} + b_S e^{i\pi(2h+l)} + b_S e^{i\pi(h+k+l)} \quad (5.1.2)$$

▼ **Extinction rules**

The sum is either 4b_Mn +/- 4b_S or 0.

assume(h, even); F(<h, 0, 0>); simplify(%)

$$4 b_{Mn} + 4 b_S$$

$$4 b_{Mn} + 4 b_S \quad (5.1.1.1)$$

assume(h, even, k, even); F(<h, k, 0>); simplify(%)

$$4 b_{Mn} + 4 b_S$$

$$4 b_{Mn} + 4 b_S \quad (5.1.1.2)$$

assume(h, even, k, odd); F(⟨h, k, 0⟩); simplify(%)

$$\frac{b_S + b_S (-1)^{2h\sim + k\sim}}{0}$$

(5.1.1.3)

assume(h, even, k, even, l, even); F(⟨h, k, l⟩); simplify(%)

$$\frac{4 b_{Mn} + 4 b_S}{4 b_{Mn} + 4 b_S}$$

(5.1.1.4)

assume(h, even, k, even, l, odd); F(⟨h, k, l⟩); simplify(%)

$$\frac{b_S + b_S (-1)^{2h\sim + l\sim}}{0}$$

(5.1.1.5)

assume(h, odd, k, odd, l, even); F(⟨h, k, l⟩); simplify(%)

$$\frac{b_S + b_S (-1)^{2h\sim + k\sim}}{0}$$

(5.1.1.6)

assume(h, odd, k, odd, l, odd); F(⟨h, k, l⟩); simplify(%)

$$\frac{4 b_{Mn} - 2 b_S + b_S (-1)^{2h\sim + k\sim} + b_S (-1)^{2h\sim + l\sim}}{4 b_{Mn} - 4 b_S}$$

(5.1.1.7)

F(⟨1, 1, 1⟩); F(⟨2, 0, 0⟩)

$$\frac{4 b_{Mn} - 4 b_S}{4 b_{Mn} + 4 b_S}$$

(5.1.1.8)

lh := [⟨1, 1, 1⟩, ⟨2, 0, 0⟩, ⟨2, 2, 0⟩, ⟨3, 1, 1⟩];

$$lh := \left[\begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}, \begin{bmatrix} 2 \\ 0 \\ 0 \end{bmatrix}, \begin{bmatrix} 2 \\ 2 \\ 0 \end{bmatrix}, \begin{bmatrix} 3 \\ 1 \\ 1 \end{bmatrix} \right] \quad (5.1.3)$$

Numerical F(h) in 10^{-12} [cm]

$$bnum := [b_{Mn} = -0.373, b_S = +0.2847];$$

$$bnum := [b_{Mn} = -0.373, b_S = 0.2847] \quad (5.1.4)$$

the scattering lengths are in 10^{-12} [cm]

$$subs(bnum, eval(F(\langle 1, 1, 1 \rangle)));$$

$$subs(bnum, eval(F(\langle 2, 0, 0 \rangle)));$$

$$subs(bnum, eval(F(\langle 2, 2, 0 \rangle)));$$

$$subs(bnum, eval(F(\langle 3, 1, 1 \rangle)));$$

$$-2.6308$$

$$-0.3532$$

$$-0.3532$$

$$-2.6308$$

(5.1.5)

Calculate the peak positions for the hkl-list. First nuclear.

$$\text{> } lh := [\langle 1, 1, 1 \rangle, \langle 2, 0, 0 \rangle, \langle 2, 2, 0 \rangle, \langle 3, 1, 1 \rangle];$$

$$mult := [8, 6, 12, 24];$$

$$seq\left(\left[HKL = h, evalf\left(2 \cdot \frac{t2 \cdot 180}{\pi}\right) \right], h \in lh\right);$$

$$mult := [8, 6, 12, 24]$$

$$\left[\begin{bmatrix} HKL = \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}, 36.54757372 \end{bmatrix}, \begin{bmatrix} HKL = \begin{bmatrix} 2 \\ 0 \\ 0 \end{bmatrix}, 42.45421864 \end{bmatrix}, \begin{bmatrix} HKL = \begin{bmatrix} 2 \\ 2 \\ 0 \end{bmatrix}, 61.59937241 \end{bmatrix}, \begin{bmatrix} HKL = \begin{bmatrix} 3 \\ 1 \\ 1 \end{bmatrix}, 73.79967604 \end{bmatrix} \right] \quad (5.1.6)$$

Then magnetic peak positions.

> $lh := \left[\left\langle \frac{1}{2}, \frac{1}{2}, \frac{1}{2} \right\rangle, \left\langle \frac{3}{2}, \frac{1}{2}, \frac{1}{2} \right\rangle, \left\langle \frac{3}{2}, \frac{3}{2}, \frac{1}{2} \right\rangle, \left\langle \frac{5}{2}, \frac{1}{2}, \frac{1}{2} \right\rangle \right];$

$seq\left(\left[HKL = h, evalf\left(\frac{t2 \cdot 2 \cdot 180}{Pi}\right) \right], h \in lh\right)$

$$\left[\left[\begin{matrix} HKL = \\ \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{matrix}, 18.03997726 \right], \left[\begin{matrix} HKL = \\ \frac{3}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{matrix}, 34.94031101 \right], \left[\begin{matrix} HKL = \\ \frac{3}{2} \\ \frac{3}{2} \\ \frac{1}{2} \end{matrix}, 46.47606041 \right], \left[\begin{matrix} HKL = \\ \frac{5}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{matrix}, 56.11235915 \right] \right] \quad (5.1.7)$$

#seq(h[i] = h || i, i = 1 ..3);

Intensity. We set the rtable ma[.].

> $lh := [\langle 1, 1, 1 \rangle, \langle 2, 0, 0 \rangle, \langle 2, 2, 0 \rangle, \langle 3, 1, 1 \rangle];$

$mult := [8, 6, 12, 24];$

> **for i from 1 to 4 do**

$ma[lh[i]] := mult[i]$

end do:

$$mult := [8, 6, 12, 24] \quad (5.1.8)$$

Lorenz factor

> $seq([HKL = h, evalf(L(t2))], h \in lh);$

$$\left[\left[\begin{matrix} HKL = \\ 1 \\ 1 \\ 1 \end{matrix}, 5.355592645 \right], \left[\begin{matrix} HKL = \\ 2 \\ 0 \\ 0 \end{matrix}, 4.091742821 \right], \left[\begin{matrix} HKL = \\ 2 \\ 2 \\ 0 \end{matrix}, 2.220194594 \right], \left[\begin{matrix} HKL = \\ 3 \\ 1 \\ 1 \end{matrix}, 1.734374830 \right] \right] \quad (5.1.9)$$

Structure factor

> $seq([HKL = h, evalf(F(h))], h \in lh);$

$$\left[\begin{array}{c} HKL = \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}, 4. b_{Mn} - 4. b_S \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 2 \\ 0 \\ 0 \end{bmatrix}, 4. b_{Mn} + 4. b_S \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 2 \\ 2 \\ 0 \end{bmatrix}, 4. b_{Mn} + 4. b_S \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 3 \\ 1 \\ 1 \end{bmatrix}, 4. b_{Mn} - 4. b_S \end{array} \right] \quad (5.1.10)$$

Numeric square of structure factor

> seq([HKL = h, evalf((subs(bnum, F(h)))²)], h ∈ lh);

$$\left[\begin{array}{c} HKL = \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}, 6.92110864 \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 2 \\ 0 \\ 0 \end{bmatrix}, 0.12475024 \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 2 \\ 2 \\ 0 \end{bmatrix}, 0.12475024 \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 3 \\ 1 \\ 1 \end{bmatrix}, 6.92110864 \end{array} \right] \quad (5.1.11)$$

Intensity without multiplicity

> seq([HKL = h, evalf(L(t2) · (subs(bnum, F(h)))²)], h ∈ lh);

$$\left[\begin{array}{c} HKL = \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}, 37.06663854 \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 2 \\ 0 \\ 0 \end{bmatrix}, 0.5104458990 \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 2 \\ 2 \\ 0 \end{bmatrix}, 0.2769698085 \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 3 \\ 1 \\ 1 \end{bmatrix}, 12.00379662 \end{array} \right] \quad (5.1.12)$$

Total Intensity

> seq([HKL = h, evalf(L(t2) · (subs(bnum, F(h)))² · ma[h])], h ∈ lh);

$$\left[\begin{array}{c} HKL = \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}, 296.5331083 \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 2 \\ 0 \\ 0 \end{bmatrix}, 3.062675394 \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 2 \\ 2 \\ 0 \end{bmatrix}, 3.323637701 \end{array} \right], \left[\begin{array}{c} HKL = \begin{bmatrix} 3 \\ 1 \\ 1 \end{bmatrix}, 288.0911189 \end{array} \right] \quad (5.1.13)$$

► Analytical version

▼ Magnetic intensity

▼ First: close to numerical calculation way. In units 10^{-12} [cm]

List of Q-vectors of the first magnetic Bragg peaks

$$\begin{aligned}
 > lh := \left[\left\langle \frac{1}{2}, \frac{1}{2}, \frac{1}{2} \right\rangle, \left\langle \frac{3}{2}, \frac{1}{2}, \frac{1}{2} \right\rangle, \left\langle \frac{3}{2}, \frac{3}{2}, \frac{1}{2} \right\rangle, \left\langle \frac{5}{2}, \frac{1}{2}, \frac{1}{2} \right\rangle \right] \\
 lh := \left[\begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}, \begin{bmatrix} \frac{3}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}, \begin{bmatrix} \frac{3}{2} \\ \frac{3}{2} \\ \frac{1}{2} \end{bmatrix}, \begin{bmatrix} \frac{5}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} \right]
 \end{aligned}$$

(6.1.1)

Double cross product.

eval vs. subs.

subs does just "syntactic substitution"

one should use delayed substitution! with '...'

$$\begin{aligned}
 > \#S := \langle s[1], s[2], s[3] \rangle; \\
 \#S := \langle s, -s, 0 \rangle; \\
 \#Q := lh[4];
 \end{aligned}$$

$$> s[1]$$

$$[s0 \ -s0 \ 0]$$

(6.1.2)

Double cross product Qp perpendicular $Q_{\perp} = [Q \times S \times Q] / (Q \cdot Q)$

$$> Qp := (S :: Vector, Q :: Vector) \rightarrow \frac{\&x'(Q, (\&x'(S, Q)))}{\cdot'(Q, Q)};$$

$$\#Qp := (S :: Vector, Q :: Vector) \rightarrow \frac{S \cdot (Q \cdot Q) - Q \cdot (Q \cdot S)}{(Q \cdot Q)};$$

$$Qp := (S::Vector, Q::Vector) \mapsto \frac{\text{LinearAlgebra:}\&x'(Q, \text{LinearAlgebra:}\&x'(S, Q))}{\text{LinearAlgebra:}\cdot'(Q, Q)}$$

(6.1.3)

> (eval('.(q, q), q = lh[1]));

$$\frac{3}{4}$$

(6.1.4)

> eval(Qp(s[1], <q₁, q₂, q₃>)) assuming real;
eval(Qp(s[1], lh[2])) assuming real;

$$\begin{bmatrix} \frac{q_2 (q_1 s_0 + q_2 s_0) + q_3^2 s_0}{q_1^2 + q_2^2 + q_3^2} \\ \frac{-q_1 (q_1 s_0 + q_2 s_0) - q_3^2 s_0}{q_1^2 + q_2^2 + q_3^2} \\ \frac{-q_1 q_3 s_0 + q_2 q_3 s_0}{q_1^2 + q_2^2 + q_3^2} \end{bmatrix}$$

$$\begin{bmatrix} \frac{5 s_0}{11} \\ -\frac{13 s_0}{11} \\ -\frac{2 s_0}{11} \end{bmatrix}$$

(6.1.5)

> (subs(q = lh[1], 'DotProduct'(q, q)));
eval(subs(q = lh[1], 'DotProduct'(q, q)));
eval(subs(q = lh[1], '.(q, q)));

eval(subs(q = lh[1], '&x'(q, q)));

$$\text{DotProduct} \left(\begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}, \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} \right)$$

$$\frac{3}{4}$$

$$\frac{3}{4}$$

$$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

(6.1.6)

sum and add are principally different. The sum is intended to evaluate summations and must be able to operate on expression generated outside the sum function. The summation index is not local to the sum. <http://www.mapleprimes.com/questions/43498-Sum-Versus-Add>

> $Fm := (Q :: Vector) \rightarrow \text{add}(r0 \cdot Qp(s[j], Q) \cdot \exp(2 \cdot \text{Pi} \cdot I \cdot (r[j] \cdot Q)), j = 1 .. 4);$
 $Phase := (Q :: Vector, r :: Vector) \rightarrow \exp(2 \cdot \text{Pi} \cdot I \cdot (r \cdot Q));$

$$Fm := Q :: Vector \mapsto \text{add} \left(r0 \cdot Qp(s_j, Q) \cdot e^{2 \cdot 1 \pi \cdot (Typesetting: -delayDotProduct(r_j, Q))}, j = 1 .. 4 \right)$$

$$Phase := (Q :: Vector, r :: Vector) \mapsto e^{2 \cdot 1 \pi \cdot (Typesetting: -delayDotProduct(r, Q))}$$

(6.1.7)

$r0$ is $\frac{1}{2} r_e \gamma$ in 10^{-12} cm, same units as for nuclear.

> $r0num := [r0 = 0.2695];$

$$r0num := [r0 = 0.2695]$$

(6.1.8)

> $\text{seq}(\text{eval}(Qp(s[i], lh[2])), i = 1 .. 4);$

$seq(eval(Phase(lh[2], r[i])), i = 1 ..4);$

$$\begin{bmatrix} \frac{5s0}{11} \\ -\frac{13s0}{11} \\ -\frac{2s0}{11} \end{bmatrix}, \begin{bmatrix} -\frac{5s0}{11} \\ \frac{13s0}{11} \\ \frac{2s0}{11} \end{bmatrix}, \begin{bmatrix} -\frac{5s0}{11} \\ \frac{13s0}{11} \\ \frac{2s0}{11} \end{bmatrix}, \begin{bmatrix} -\frac{5s0}{11} \\ \frac{13s0}{11} \\ \frac{2s0}{11} \end{bmatrix}$$

1, 1, 1, -1

(6.1.9)

$$\begin{aligned} > lh := \left[\left\langle \frac{1}{2}, \frac{1}{2}, \frac{1}{2} \right\rangle, \left\langle -\frac{1}{2}, -\frac{1}{2}, \frac{3}{2} \right\rangle, \left\langle \frac{3}{2}, -\frac{1}{2}, -\frac{1}{2} \right\rangle, \left\langle \frac{1}{2}, -\frac{3}{2}, \frac{1}{2} \right\rangle \right] \\ lh := \left[\begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}, \begin{bmatrix} -\frac{1}{2} \\ -\frac{1}{2} \\ \frac{3}{2} \end{bmatrix}, \begin{bmatrix} \frac{3}{2} \\ -\frac{1}{2} \\ -\frac{1}{2} \end{bmatrix}, \begin{bmatrix} \frac{1}{2} \\ -\frac{3}{2} \\ \frac{1}{2} \end{bmatrix} \right] \end{aligned}$$

(6.1.10)

$> seq\left(\left[HKL = h, evalf\left(j0\left(\frac{qvec}{4 \cdot \text{Pi}}, MMN3\right)\right)\right], h \in lh\right);$

$$\left[\begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}, 0.9376498075 \right], \left[\begin{bmatrix} -\frac{1}{2} \\ -\frac{1}{2} \\ \frac{3}{2} \end{bmatrix}, 0.7939001556 \right], \left[\begin{bmatrix} \frac{3}{2} \\ -\frac{1}{2} \\ -\frac{1}{2} \end{bmatrix}, 0.7939001556 \right], \left[\begin{bmatrix} \frac{1}{2} \\ -\frac{3}{2} \\ \frac{1}{2} \end{bmatrix}, \right]$$

(6.1.11)

0.7939001556

> seq([HKL = h, eval(Fm(h))], h ∈ lh);

$$HKL = \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}, \begin{bmatrix} 4 r0 s0 \\ -4 r0 s0 \\ 0 \end{bmatrix}, HKL = \begin{bmatrix} -\frac{1}{2} \\ -\frac{1}{2} \\ \frac{3}{2} \end{bmatrix}, \begin{bmatrix} 4 r0 s0 \\ -4 r0 s0 \\ 0 \end{bmatrix}, HKL = \begin{bmatrix} \frac{3}{2} \\ -\frac{1}{2} \\ -\frac{1}{2} \end{bmatrix}, \begin{bmatrix} -\frac{4 r0 s0}{11} \\ -\frac{28 r0 s0}{11} \\ \frac{16 r0 s0}{11} \end{bmatrix}, HKL = \begin{bmatrix} \frac{1}{2} \\ -\frac{3}{2} \\ \frac{1}{2} \end{bmatrix},$$

(6.1.12)

$$\begin{bmatrix} \frac{28 r0 s0}{11} \\ \frac{4 r0 s0}{11} \\ -\frac{16 r0 s0}{11} \end{bmatrix}$$

> eval(Fm(lh[1]), {r0 = 0.2695, s0 = 3});
eval(%); eval(Norm(%), Euclidean) · 0.93);

$$\begin{array}{cccc}
\begin{bmatrix} 0.606375 \\ -0.606375 \\ 0.0 \end{bmatrix} & \begin{bmatrix} -0.606375 \\ 0.606375 \\ 0.0 \end{bmatrix} & \begin{bmatrix} -0.606375 \\ 0.606375 \\ 0.0 \end{bmatrix} & \begin{bmatrix} -0.606375 \\ 0.606375 \\ 0.0 \end{bmatrix} \\
\begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} \cdot \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} & \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} \cdot \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} & \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} \cdot \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} & \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} \cdot \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} \\
\begin{bmatrix} 3.234000000000000 \\ -3.234000000000000 \\ 0. \\ 4.253416995 \end{bmatrix} & & &
\end{array}$$

(6.1.13)

Absolute values of Fmag: to be compared with fullprof

> seq([HKL = h, evalf(mMn(qvec) · Norm(eval(Fm(h), {r0 = 0.2695, s0 = 3}), Euclidean))], h ∈ lh);

$$\left[\begin{array}{l} \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}, 4.288403896, \begin{bmatrix} -\frac{1}{2} \\ -\frac{1}{2} \\ \frac{3}{2} \end{bmatrix}, 3.630955283, \begin{bmatrix} \frac{3}{2} \\ -\frac{1}{2} \\ -\frac{1}{2} \end{bmatrix}, 1.896204555, \begin{bmatrix} \frac{1}{2} \\ -\frac{3}{2} \\ \frac{1}{2} \end{bmatrix}, 1.896204555 \end{array} \right] \quad (6.1.14)$$

> seq([HKL = h, evalf((mMn(qvec) · Norm(eval(Fm(h), {r0 = 0.2695, s0 = 3}), Euclidean))^2)], h ∈ lh);

$$\left[\begin{array}{c} HKL = \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}, 18.39040798 \\ HKL = \begin{bmatrix} -\frac{1}{2} \\ -\frac{1}{2} \\ \frac{3}{2} \end{bmatrix}, 13.18383627 \\ HKL = \begin{bmatrix} \frac{3}{2} \\ -\frac{1}{2} \\ -\frac{1}{2} \end{bmatrix}, 3.595591711 \\ HKL = \begin{bmatrix} \frac{1}{2} \\ -\frac{3}{2} \\ \frac{1}{2} \end{bmatrix}, 3.595591711 \end{array} \right] \quad (6.1.15)$$

Total magnetic intensities

> seq([HKL = h, evalf(2·L(t2)·(mMn(qvec)·Norm(eval(Fm(h), {r0 = 0.2695, s0 = 3}), Euclidean))²)], h ∈ lh);

$$\left[\begin{array}{c} HKL = \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}, 757.5643460 \\ HKL = \begin{bmatrix} -\frac{1}{2} \\ -\frac{1}{2} \\ \frac{3}{2} \end{bmatrix}, 153.3570245 \\ HKL = \begin{bmatrix} \frac{3}{2} \\ -\frac{1}{2} \\ -\frac{1}{2} \end{bmatrix}, 41.82464306 \\ HKL = \begin{bmatrix} \frac{1}{2} \\ -\frac{3}{2} \\ \frac{1}{2} \end{bmatrix}, 41.82464306 \end{array} \right] \quad (6.1.16)$$

Lorentz factor

> seq([HKL = h, evalf(L(t2))], h ∈ lh);

$$\left[\begin{array}{c} HKL = \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}, 20.59672484 \\ HKL = \begin{bmatrix} -\frac{1}{2} \\ -\frac{1}{2} \\ \frac{3}{2} \end{bmatrix}, 5.816100162 \\ HKL = \begin{bmatrix} \frac{3}{2} \\ -\frac{1}{2} \\ -\frac{1}{2} \end{bmatrix}, 5.816100162 \\ HKL = \begin{bmatrix} \frac{1}{2} \\ -\frac{3}{2} \\ \frac{1}{2} \end{bmatrix}, 5.816100162 \end{array} \right] \quad (6.1.17)$$

> #assume(s0, real);

seq([HKL = h, evalf(2·L(t2)·(mMn(qvec)·Norm(eval(Fm(h), {r0 = 0.2695}), Euclidean))²)], h ∈ lh);

> h := lh[1];

$$h := \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix} \quad (6.1.18)$$

```
> simplify(eval(L(t2)), assume = real);
```

$$20.59672484 \quad (6.1.19)$$

```
> 2·L(t2)·(mMn(qvec))2·eval(Fm(h).Fm(h), {r0 = 0.2695}) :
> simplify(%, 'assume=real' );
```

$$84.17381632 \text{ s}^2 \quad (6.1.20)$$

Experimental magnetic/nuclear (Im/In) intensities for 1/2,1/2,1/2 and 1,1,1 is 5174/2095

In=C*2.6.53

Im=C*89.8*s0^2

So mag mom in muB

```
> evalf( (296.53 / 84.18 * 5174 / 2095 * 2)1/2 ); unassign('h');
```

$$4.171248414 \quad (6.1.21)$$

```
> #evalf( 2·L(t2)·(mMn(qvec)·Norm(eval(Fm(h), {r0 = 0.2695}), Euclidean))2 )
```

```
> V := <v[1], v[2], v[3]>;
Norm(V);
(Norm(V, Euclidean))2;
V · V;
```

► **Playing. Alternative way to create the list of intensities using inert form of Fm. All analytical!**

- ▶ **Calculate moment analitically**
- ▶ **Scalar product in anybasis**
- ▶ **Conversions between basis and dual basis (reciprocal)**
- ▶ **Calculating cross product in anybasis**
- ▶ **Own Levi-Civita**
- ▶ **some useful stuff About sum and add in maple**