

# Determination of the magnetic structure from powder neutron diffraction

Vladimir Pomjakushin

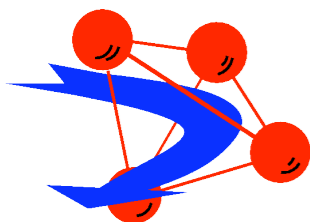
*Laboratory for Neutron Scattering, ETHZ and PSI*

A Hands-on Workshop on X-rays, Synchrotron Radiation and Neutron Diffraction  
Techniques

June 18-22, 2008, Paul Scherrer Institut, Villigen, Switzerland

Lecture notes:

<http://sinq.web.psi.ch/sinq/instr/hrpt/praktikum>



PSI

**ETH**

# Literature on (magnetic) neutron scattering

## Neutron scattering (general)

S.W. Lovesey, “*Theory of Neutron Scattering from Condensed Matter*”, Oxford Univ. Press, 1987. Volume 2 for magnetic scattering. **Definitive formal treatment**

G.L. Squires, “*Intro. to the Theory of Thermal Neutron Scattering*”, C.U.P., 1978, Republished by Dover, 1996. **Simpler version of Lovesey.**

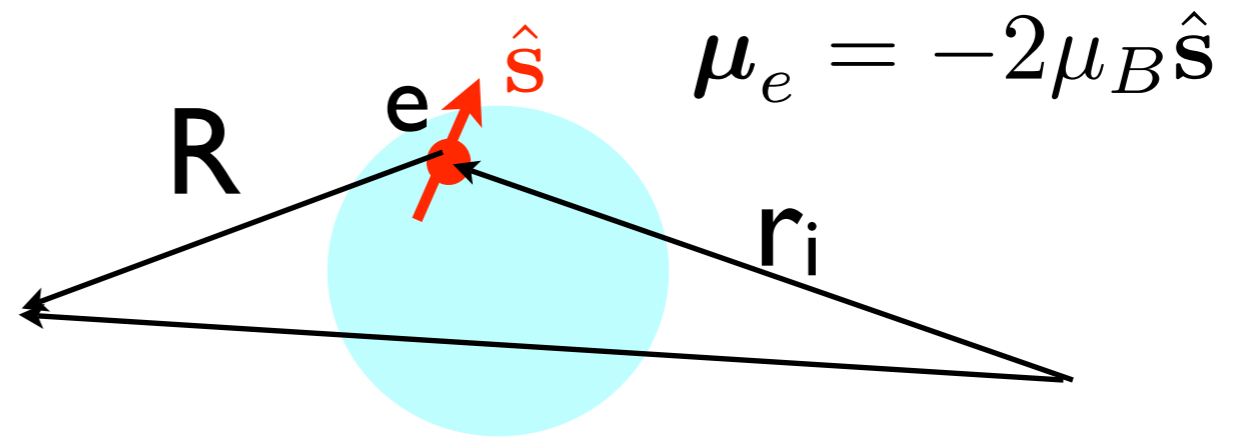
All you need to know about magnetic neutron diffraction. Symmetry, representation analysis

Yu.A. Izyumov, V. E. Naish and R. P. Ozerov, “*Neutron diffraction of magnetic materials*”, New York [etc.]: Consultants Bureau, 1991.

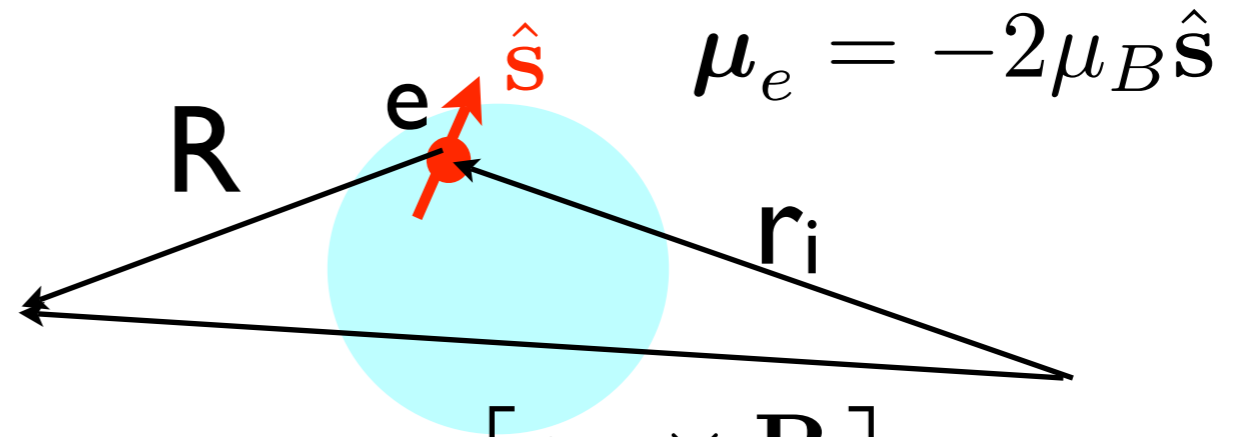
# Overview of Lecture

- Principles of magnetic neutron scattering/diffraction
- Types of magnetic structures
- Description of all possible magnetic structures. k-vector formalism for classifying the magnetic modes
- Real example of magnetic structure determination

# Magnetic neutron scattering on an atom



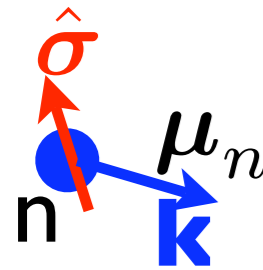
# Magnetic neutron scattering on an atom



Magnetic field from an electron

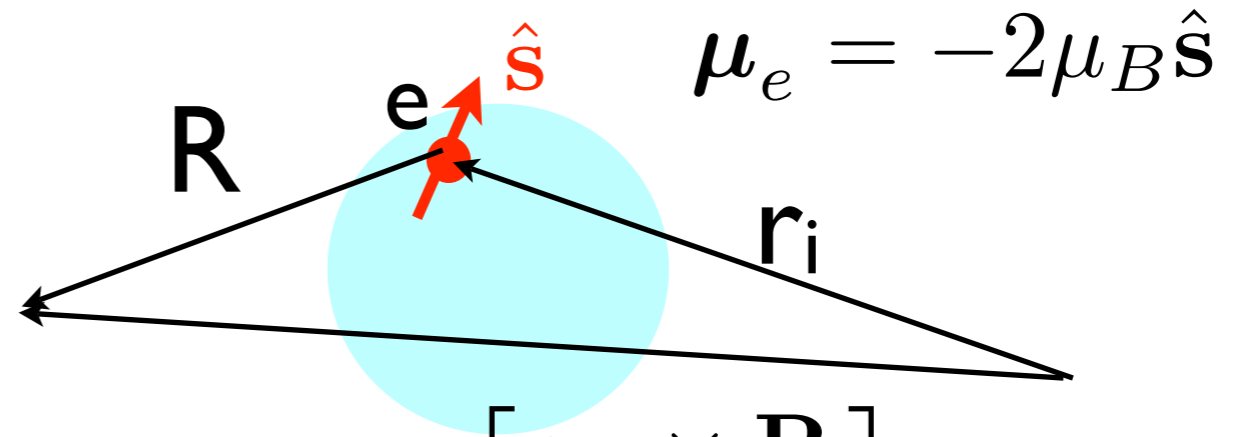
$$\mathbf{H}(\mathbf{R}) = -\text{rot} \left[ \frac{\mu_e \times \mathbf{R}}{|\mathbf{R}|^3} \right] + \text{transl. part}$$

# Magnetic neutron scattering on an atom



A blue dot labeled 'n' represents a neutron. A red arrow labeled  $\hat{\sigma}$  points upwards and to the right, representing the spin. A blue arrow labeled  $\mathbf{k}$  points to the right, representing the wave vector. A blue arrow labeled  $\mu_n$  points to the right, representing the magnetic moment.

$$\mu_n = 2\gamma\mu_n \frac{\hat{\sigma}}{2}$$



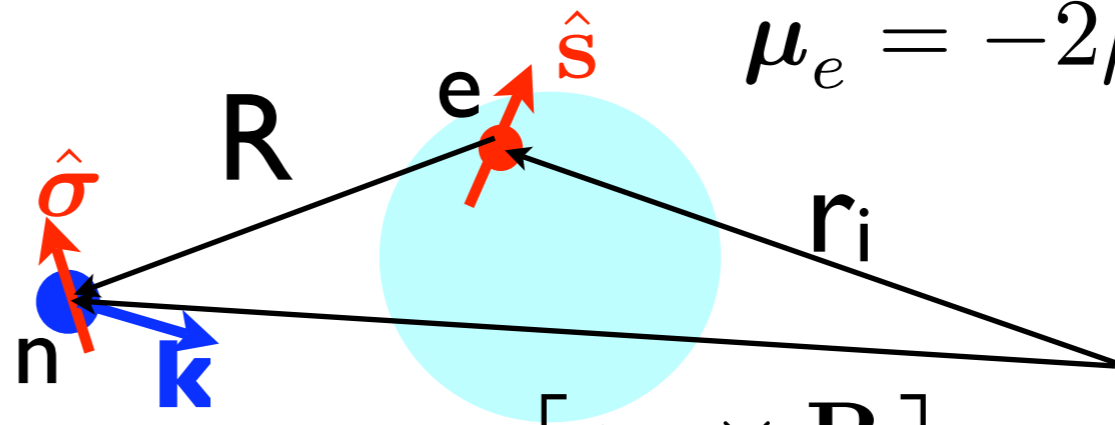
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$$\mu_e = -2\mu_B \hat{S}$$



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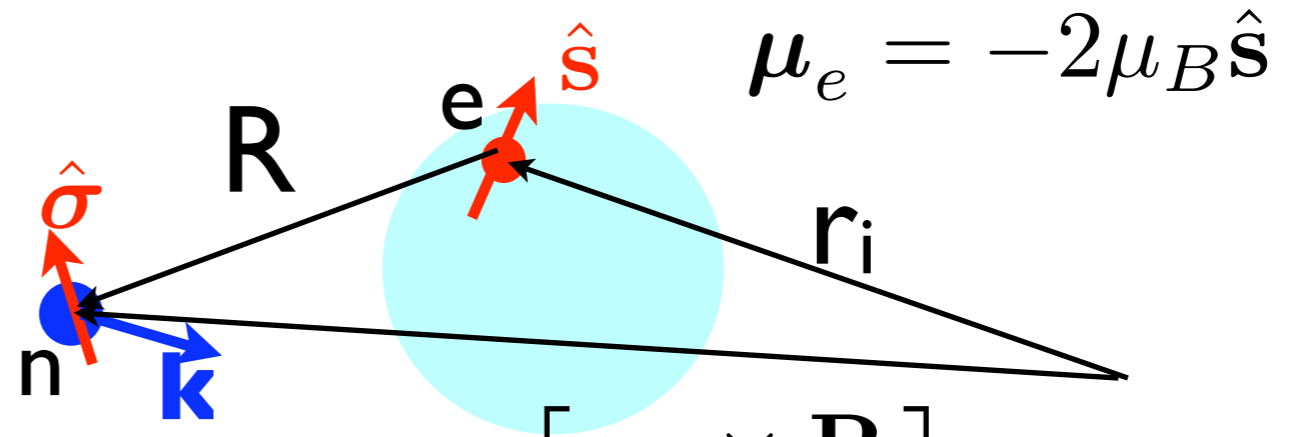
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neutron-electron dipole interaction

$$V(\mathbf{R}) = -\gamma\mu_n \hat{\sigma} \mathbf{H}(\mathbf{R})$$

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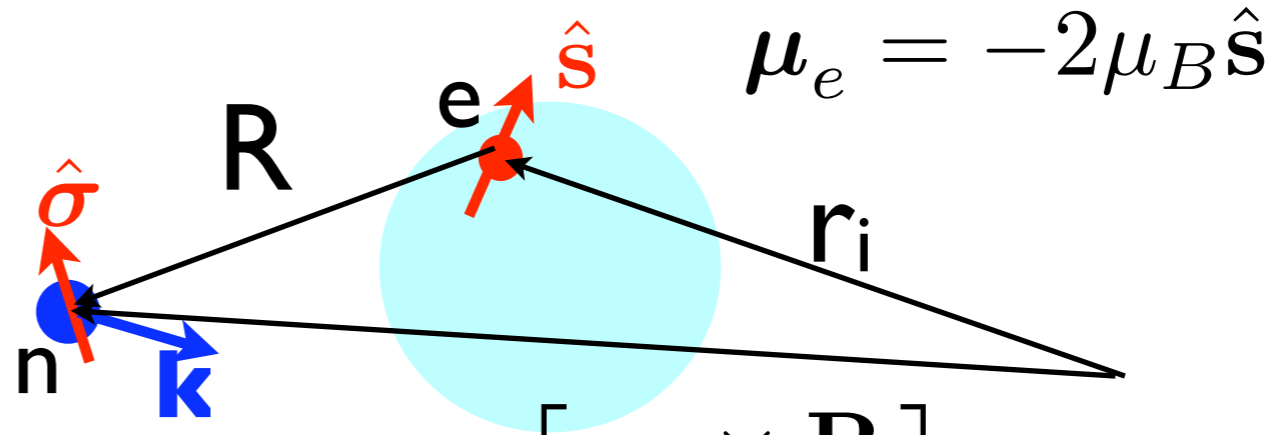
averaging over neutron coordinates

$$\langle \mathbf{k}' | V(\mathbf{R}) | \mathbf{k} \rangle_{\mathbf{q} = \mathbf{k}' - \mathbf{k}} = \gamma r_e \hat{\boldsymbol{\sigma}} \frac{1}{q^2} [\mathbf{q} \times [\hat{\mathbf{s}}_i e^{i\mathbf{q}\mathbf{r}_i} \times \mathbf{q}]]$$



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magnetic interaction operator  $\hat{\mathbf{Q}}_{\perp}$

# Magnetic neutron scattering on an atom

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x-ray scattering length:  $Z r_e$

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### Comparison of neutron scattering lengths (fm)

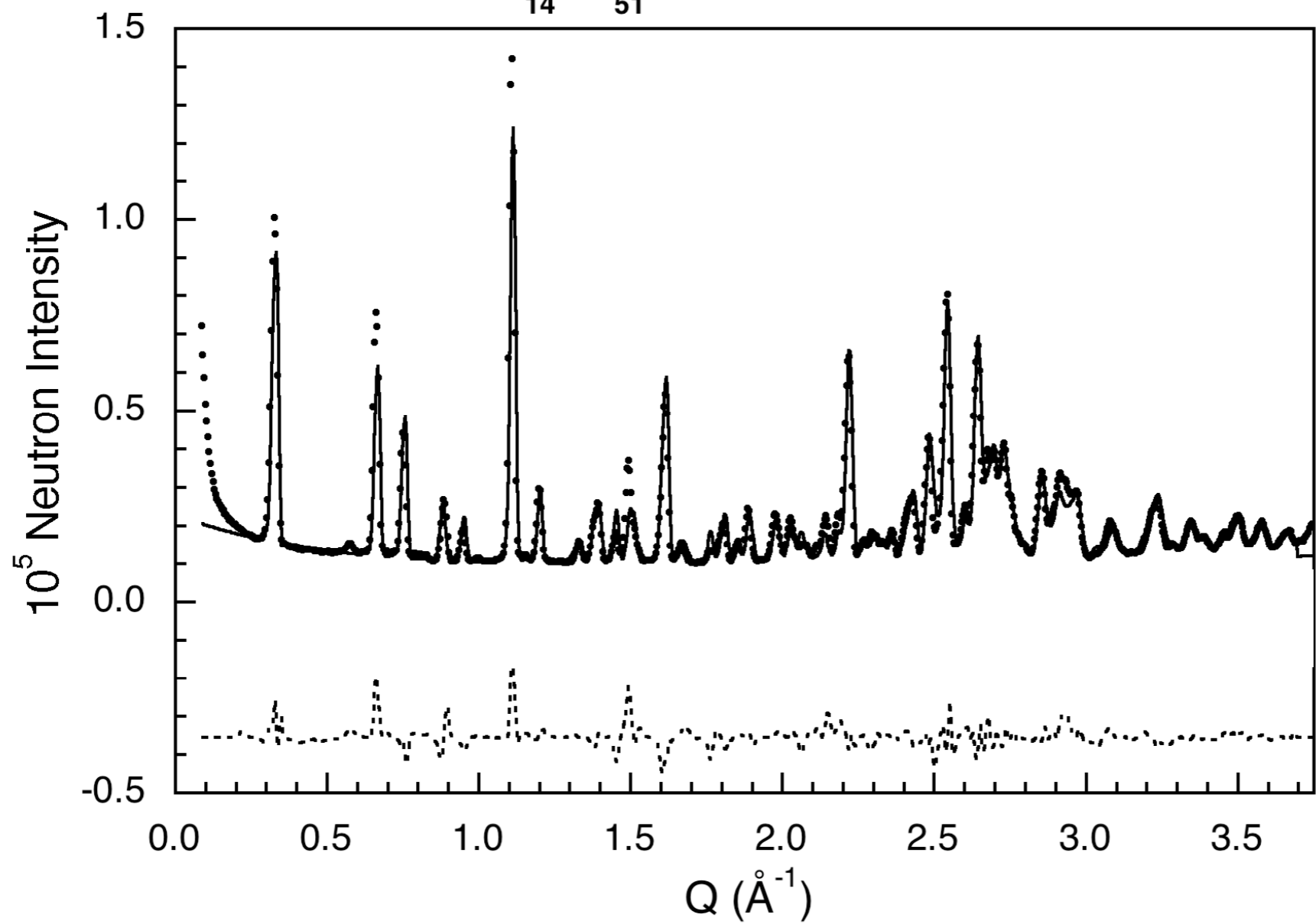
**magnetic**

**Mn<sup>3+</sup> (S=2): -10.8, Cu<sup>2+</sup> (S=1/2): -2.65**

**nuclear**

**Mn : -3.7, Cu: 7.7**

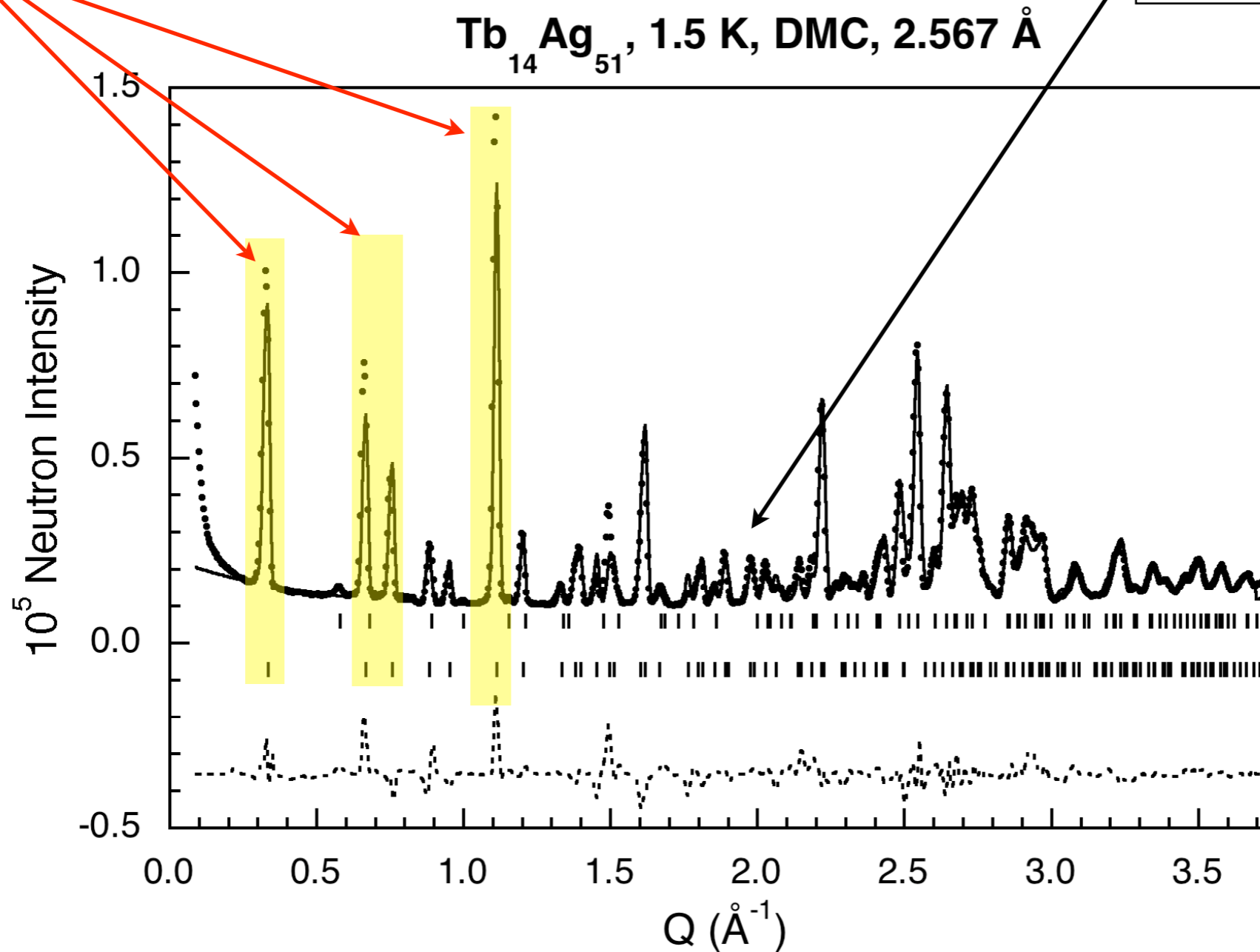
Tb<sub>14</sub>Ag<sub>51</sub>, 1.5 K, DMC, 2.567 Å



# magnetic scattering intensity can be larger than the nuclear one

magnetic

nuclear





# Magnetic neutron scattering on an atom

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## 2. q-dependence

$$\text{“magnetic scattering amplitude”} = \gamma r_e \langle \hat{\mathbf{Q}}_{\perp} \rangle, \\ \frac{1}{q^2} [\mathbf{q} \times \hat{\mathbf{Q}} \times \mathbf{q}]$$

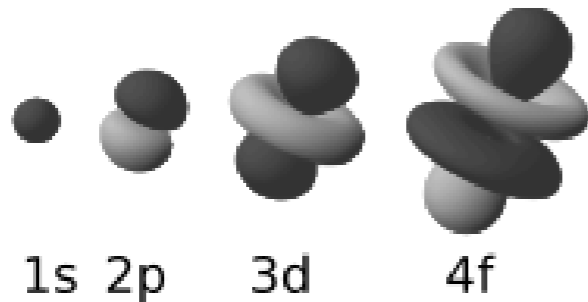
# Magnetic neutron scattering on an atom

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$$\langle \hat{\mathbf{Q}} \rangle = \left\langle \sum_i \hat{\mathbf{s}}_i e^{i\mathbf{q}\mathbf{r}_i} \right\rangle = \mathbf{S} \int d\mathbf{r} \rho_s(\mathbf{r}) e^{i\mathbf{q}\mathbf{r}}$$



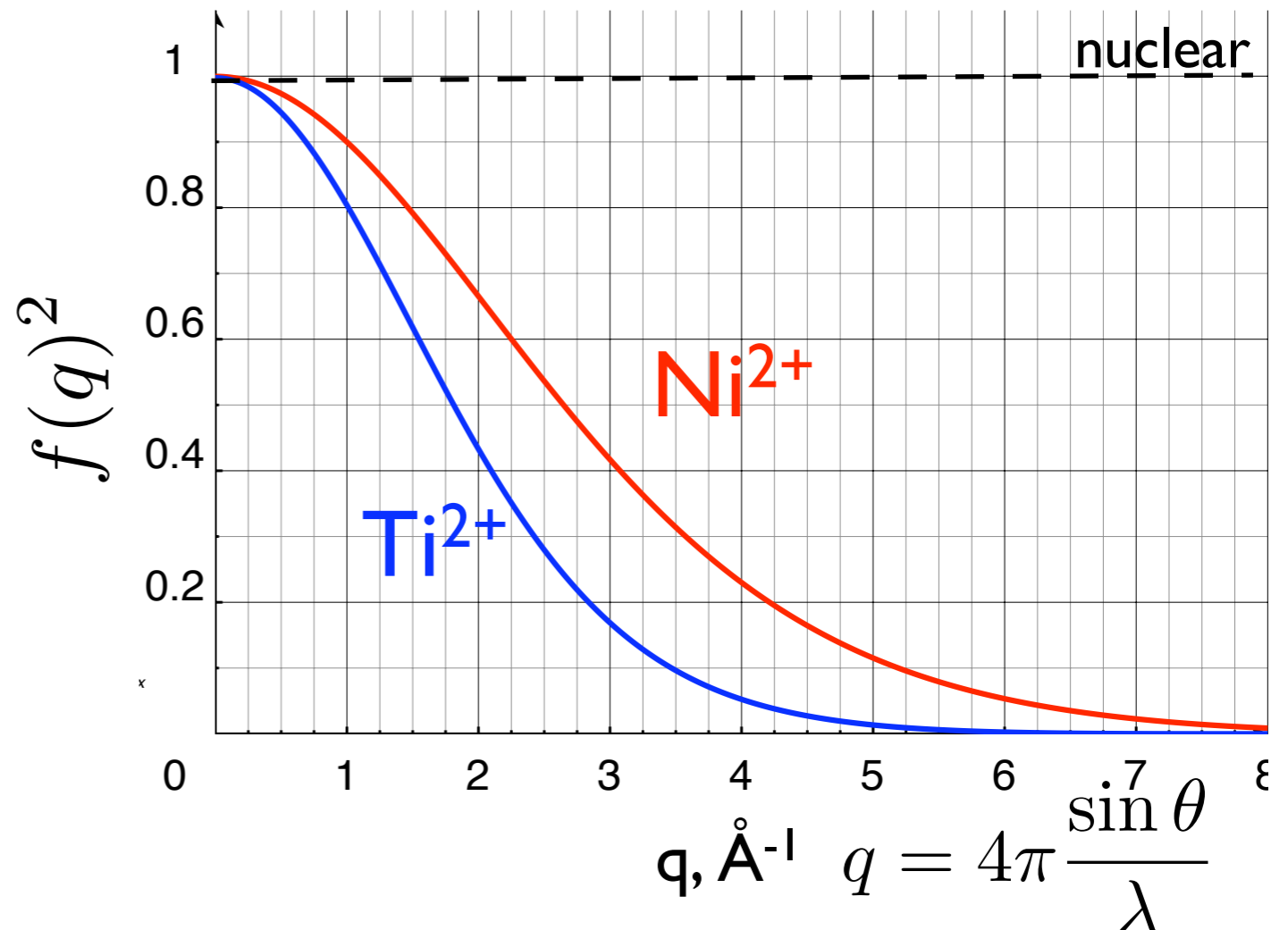
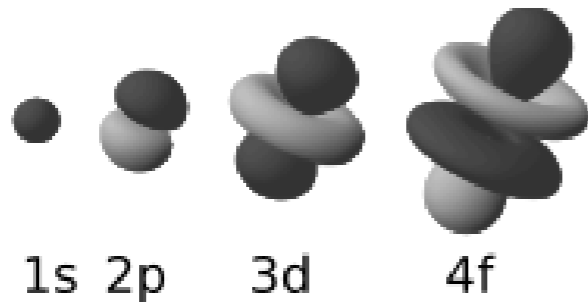
# Magnetic neutron scattering on an atom

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“magnetic scattering amplitude” =  $\gamma r_e \langle \hat{\mathbf{Q}}_{\perp} \rangle$ ,

Fourier image of the spin density in atom  
or magnetic form-factor

$$\langle \hat{\mathbf{Q}} \rangle = \left\langle \sum_i \hat{\mathbf{s}}_i e^{i\mathbf{q}\mathbf{r}_i} \right\rangle = \mathbf{S} \int d\mathbf{r} \rho_s(\mathbf{r}) e^{i\mathbf{q}\mathbf{r}} = \mathbf{S} f(q)$$



# Magnetic neutron scattering on an atom

$$\text{“magnetic scattering amplitude”} = \gamma r_e \langle \hat{\mathbf{Q}}_{\perp} \rangle$$

$$\mathbf{Q}_{\perp} = \tilde{\mathbf{q}} \times \mathbf{Q} \times \tilde{\mathbf{q}} = [\tilde{\mathbf{q}} \times \mathbf{S} \times \tilde{\mathbf{q}}] f(q)$$

$$\tilde{\mathbf{q}} = \mathbf{q}/q$$

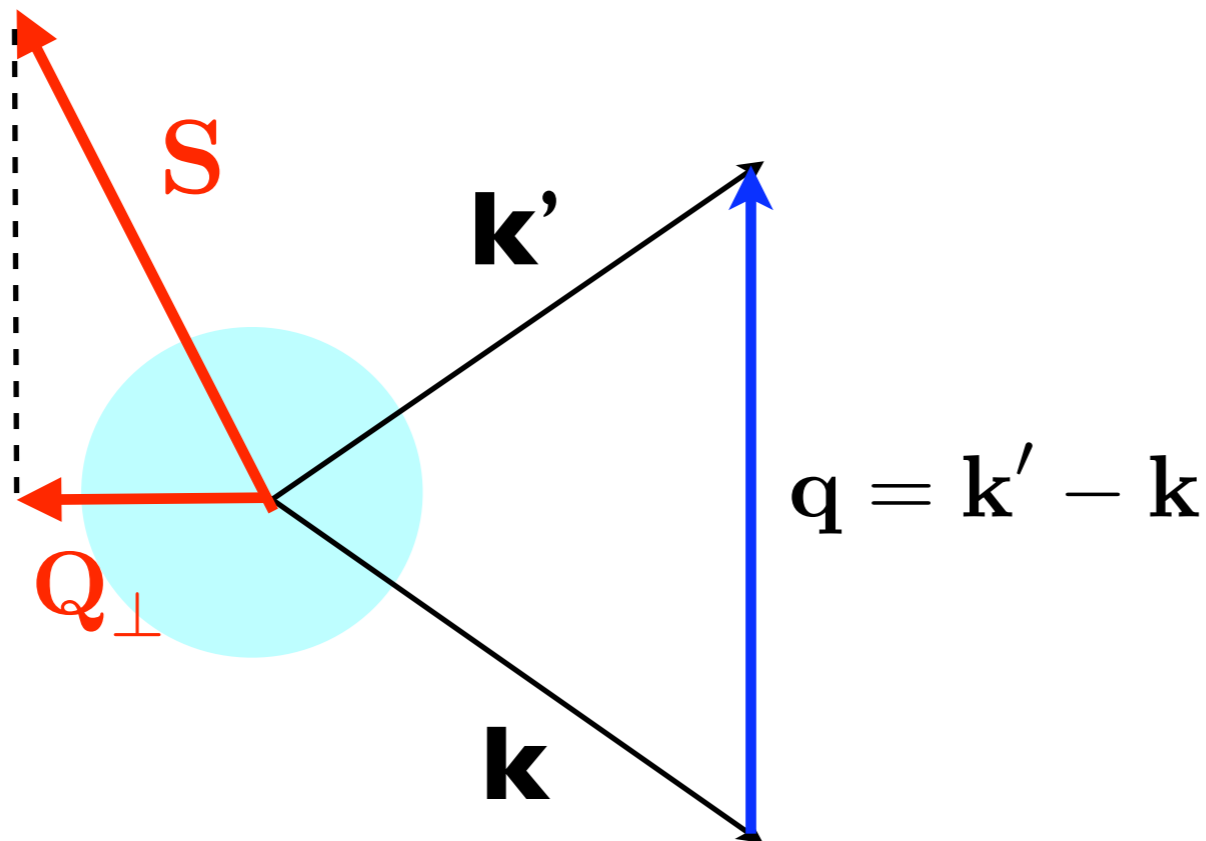
# Magnetic neutron scattering on an atom

## 3. geometry

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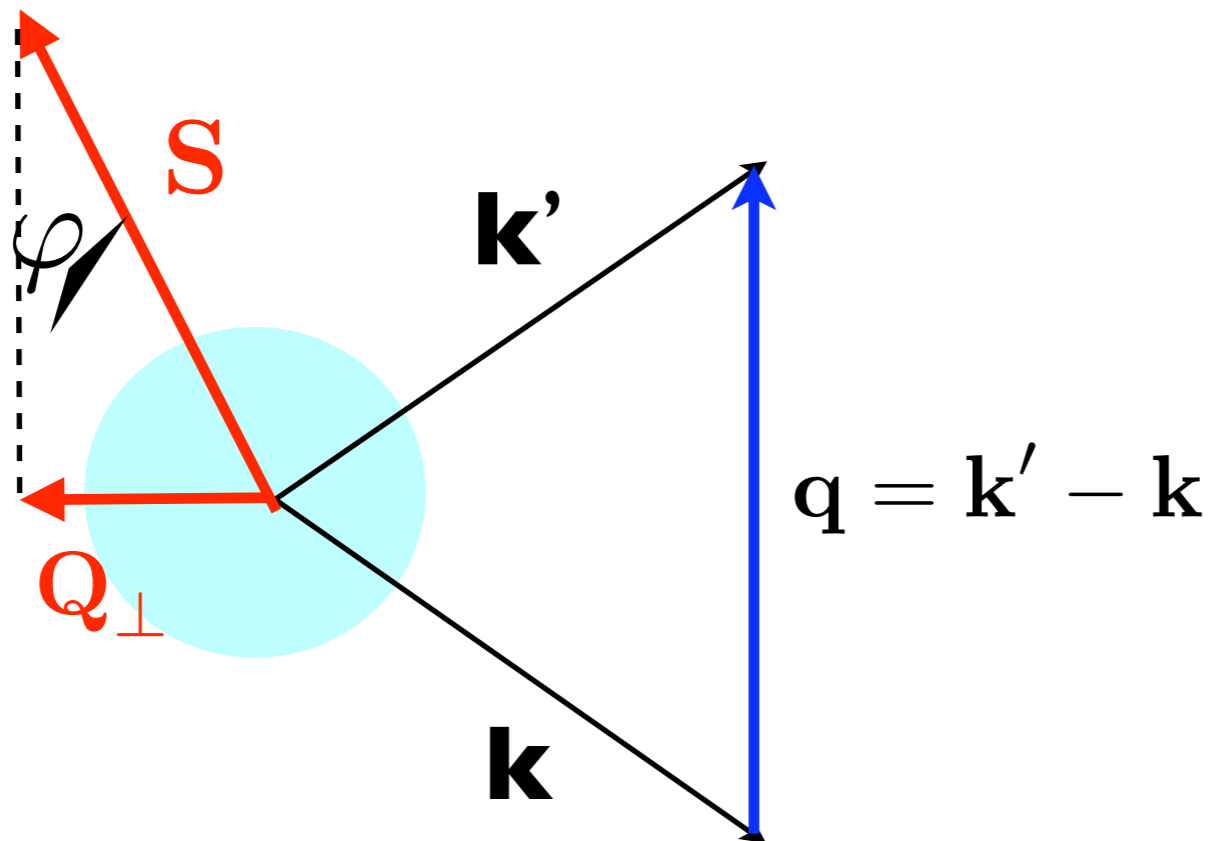
# Magnetic neutron scattering on an atom

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$$|\mathbf{Q}_{\perp}| = |\mathbf{S}| \sin(\varphi)$$

# Elastic scattering intensity

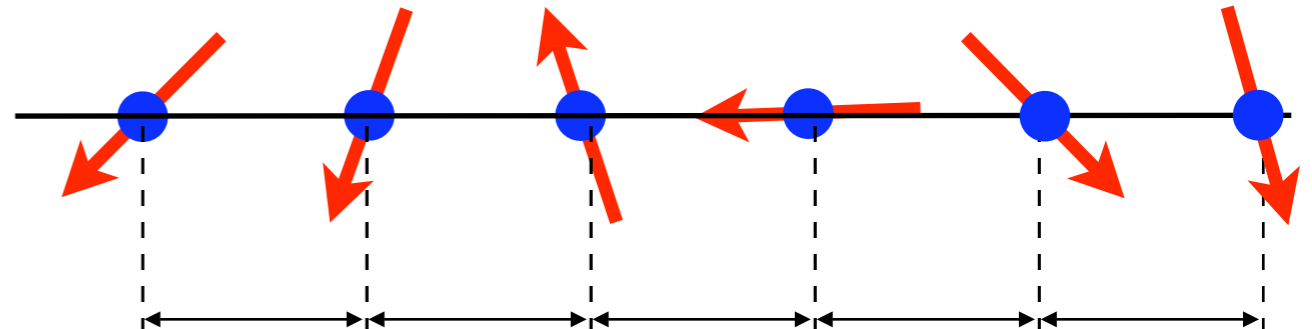
Neutron scattering cross-section  
(for unpolarized neutron beam)

$$\frac{d\sigma}{d\Omega} \propto |\mathbf{Q}_{\perp}|^2$$

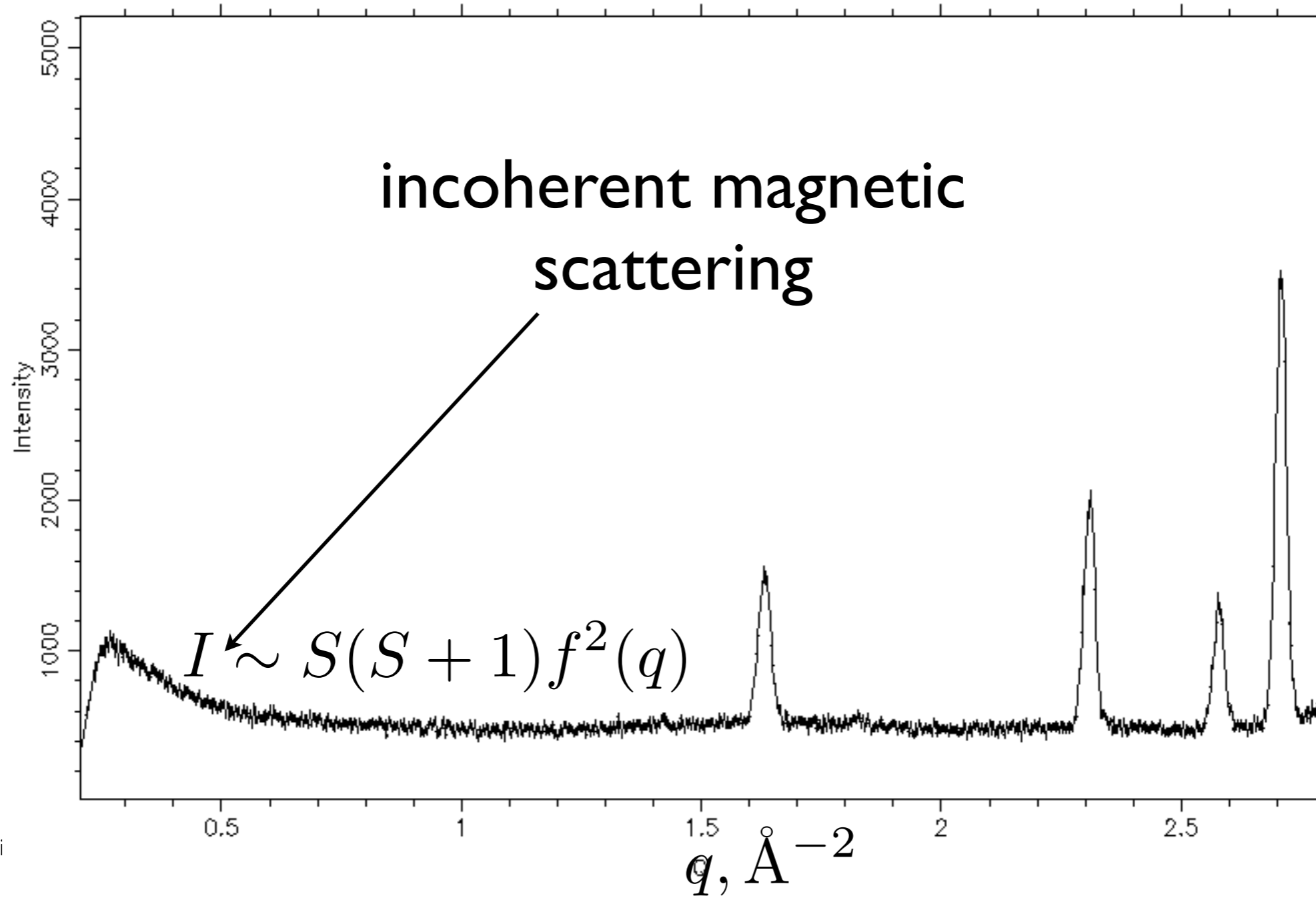
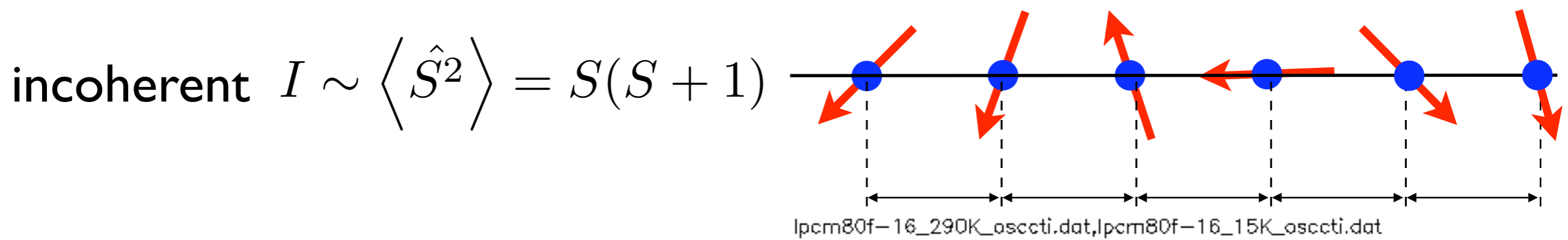


# Elastic scattering on a lattice of spins

incoherent  $I \sim \langle \hat{S}^2 \rangle = S(S + 1)$



# Elastic scattering on a lattice of spins

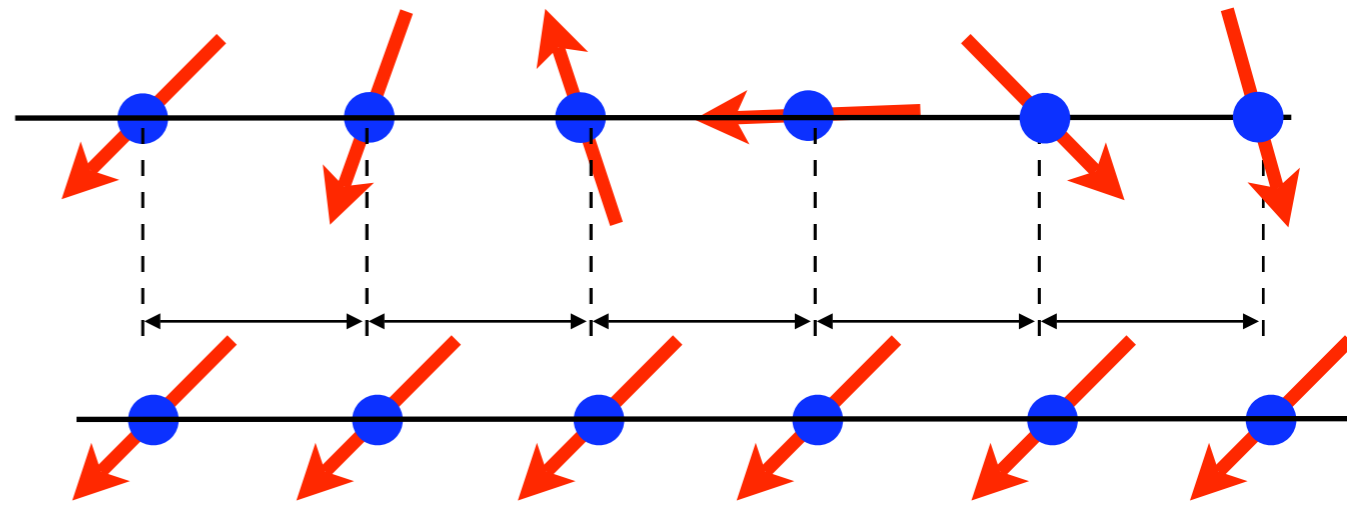


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coherent Bragg scattering

$$I \sim S^2 F_{HKL}^2$$

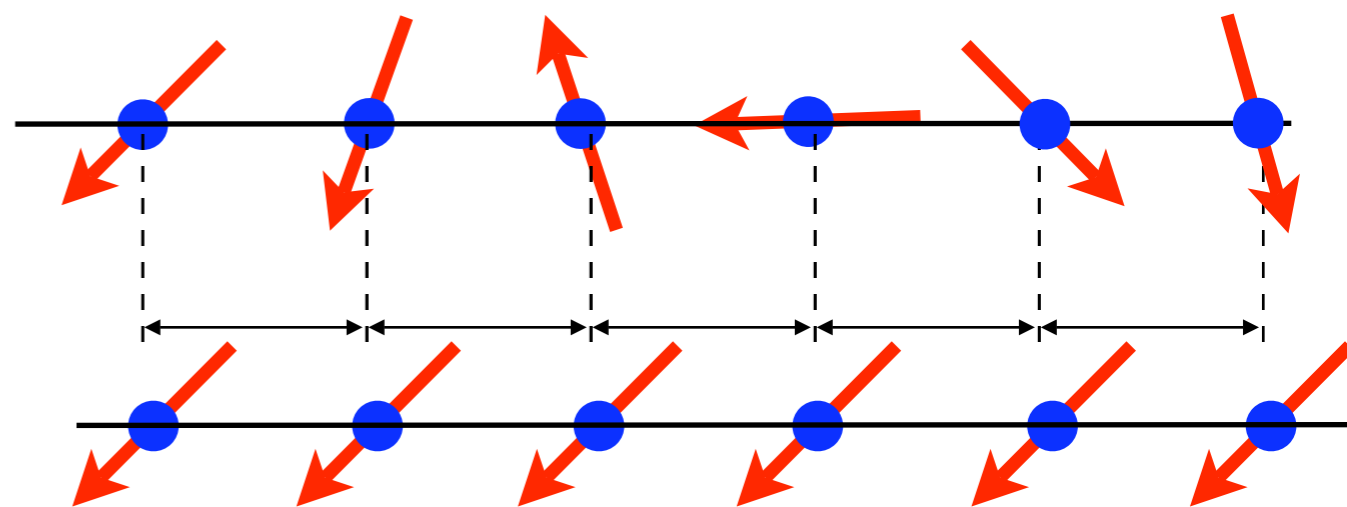


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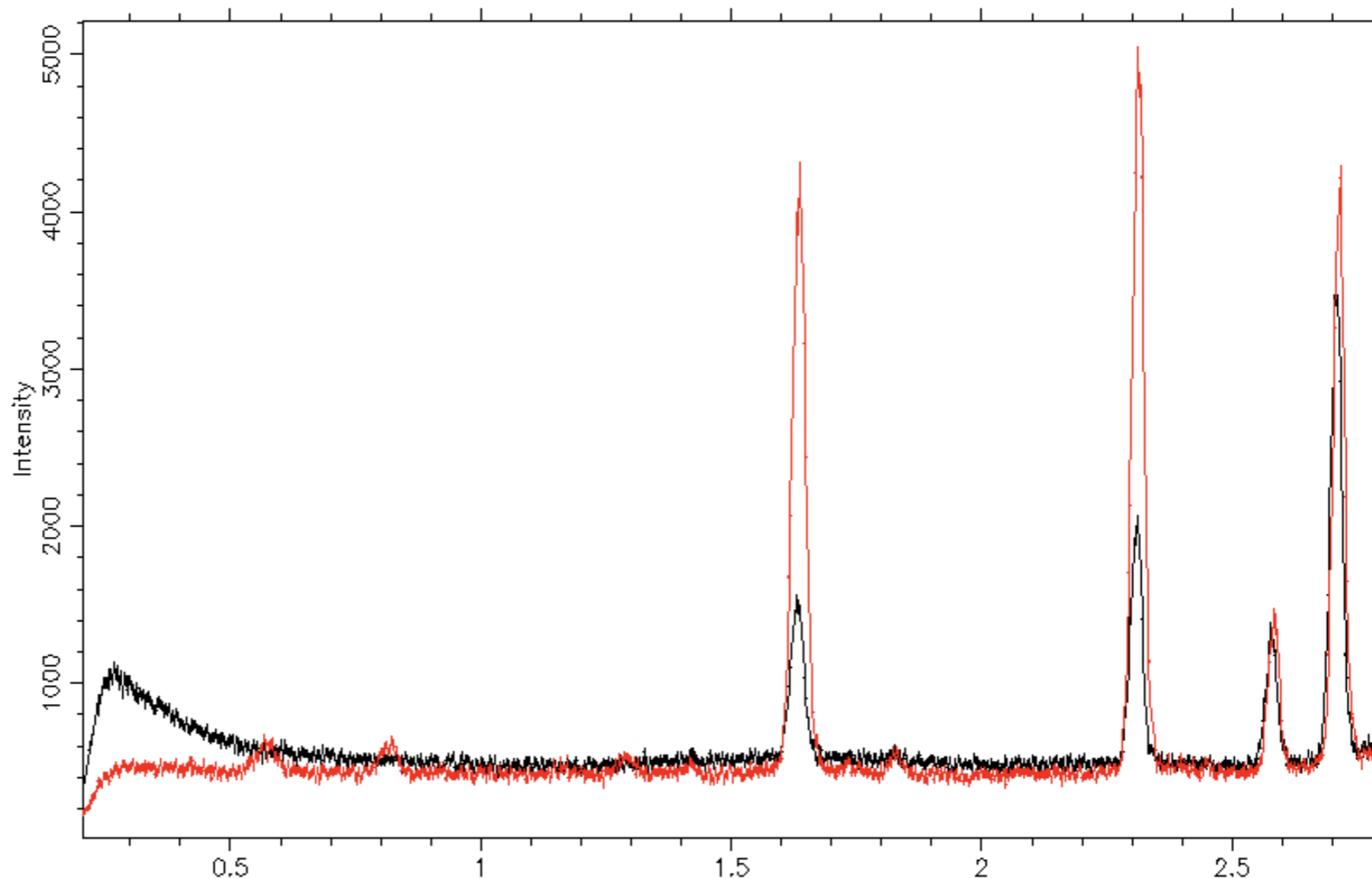
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lpcm80f-16\_290K\_osccti.dat, lpcm80f-16\_15K\_osccti.dat





# Non-polarized neutron diffraction

$$I^{++} \propto \langle |\mathbf{Q}_\perp \sigma_n + F|^2 \rangle_{\sigma_n}$$

average over neutron polarization

$$I \propto \langle (\mathbf{Q}_\perp \sigma_n)(\mathbf{Q}_\perp^* \sigma_n) + FF^* + \cancel{\sigma_n(F\mathbf{Q}_\perp^* + F^*\mathbf{Q}_\perp)} \rangle_{\sigma_n}$$

no magnetic/nuclear interference

$$I \propto |\mathbf{Q}_\perp|^2 + |F|^2$$

Magnetic and nuclear scattering are completely independent and can be treated as two independent phases in the Rietveld refinement

# Interference between nuclear and magnetic scattering

## General note:

When the magnetic unit cell is larger than the nuclear one (propagation vector  $k \neq 0$ ) the interference between nuclear and magnetic scattering is absent in any (un)polarized neutron diffraction experiment.

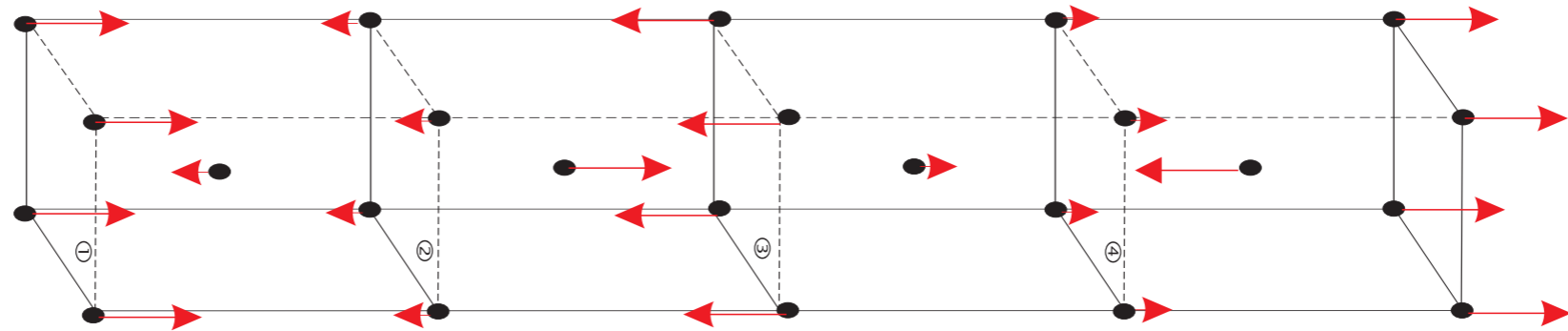
*Reason:* Magnetic Bragg peaks appear at different positions in reciprocal space

# Only amplitudes can be determined

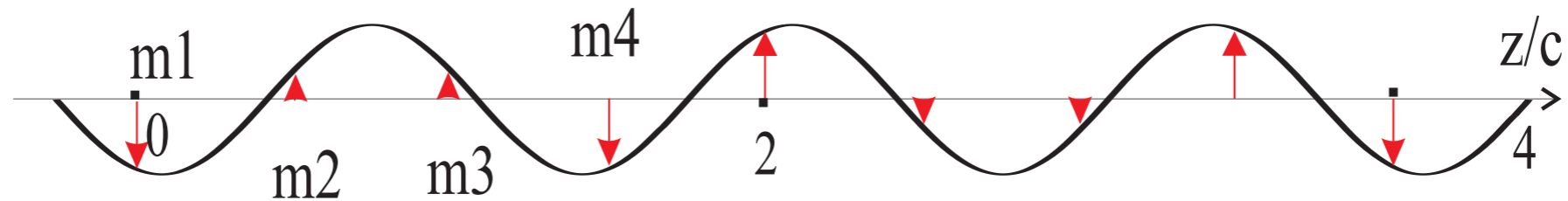
$$\mathbf{S} = \mathbf{S}_0 \cos(2\pi kz + \phi), k = \frac{3}{4}$$

Spin/atom magnetic moment  $\rightarrow$   $\mathbf{S}$

$\mathbf{S}_0$   $\leftarrow$  Amplitude



$$\phi = 7\pi/8$$

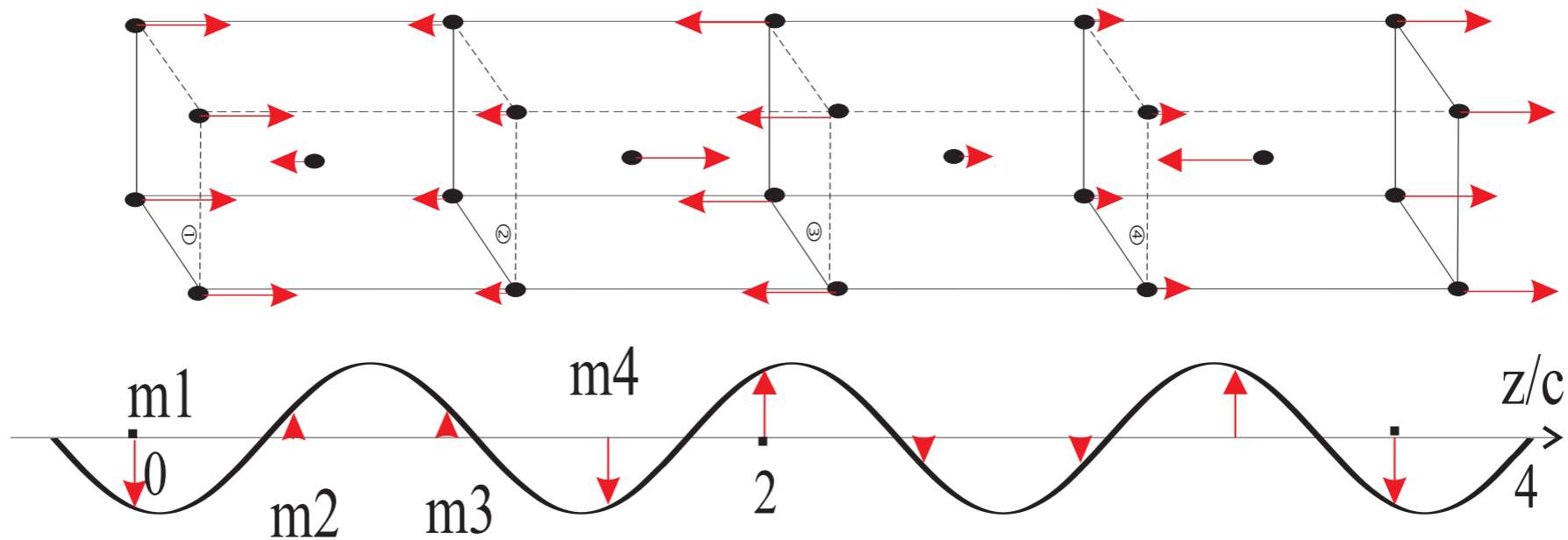




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$$\mathbf{S} = \mathbf{S}_0 \cos(2\pi kz + \phi), k = \frac{3}{4} \quad I \sim S_0^2 + S_0 F \cos(\phi)$$

Spin/atom magnetic moment  $\swarrow$   $\nwarrow$  Amplitude

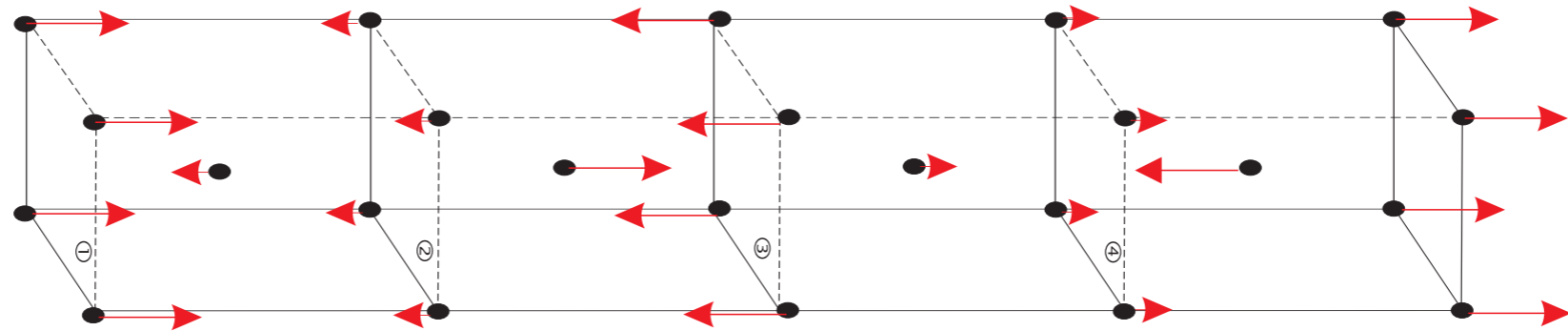
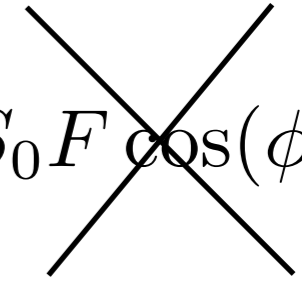


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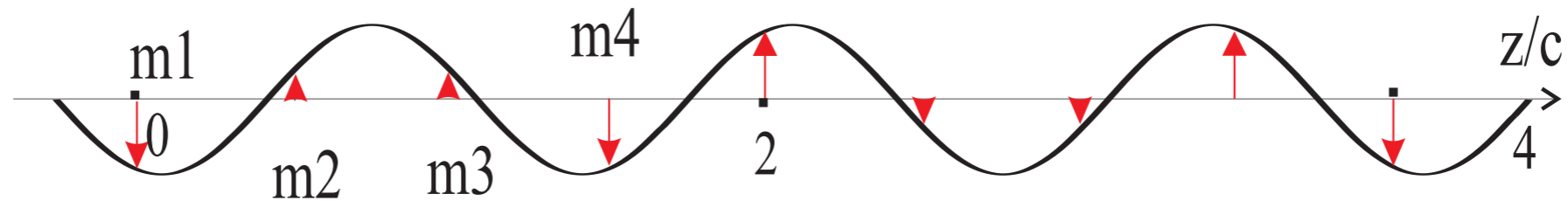
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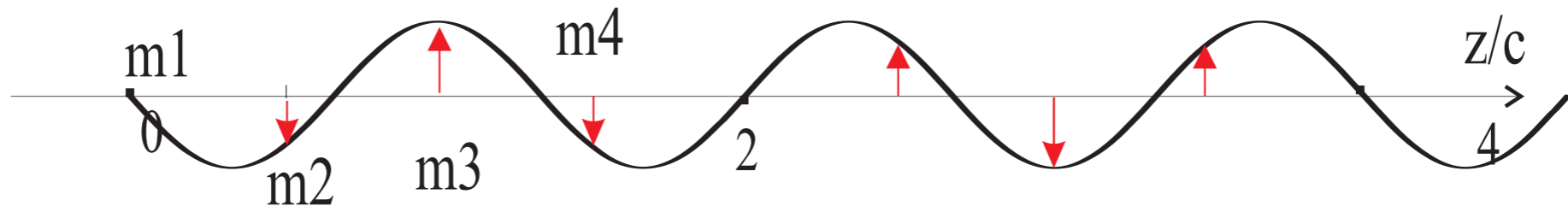
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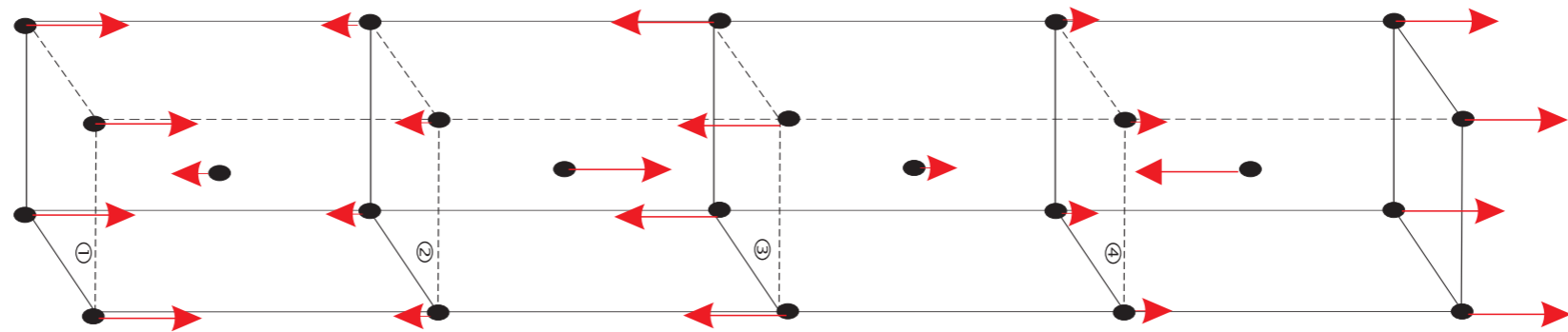


$$\phi = \pi/2$$

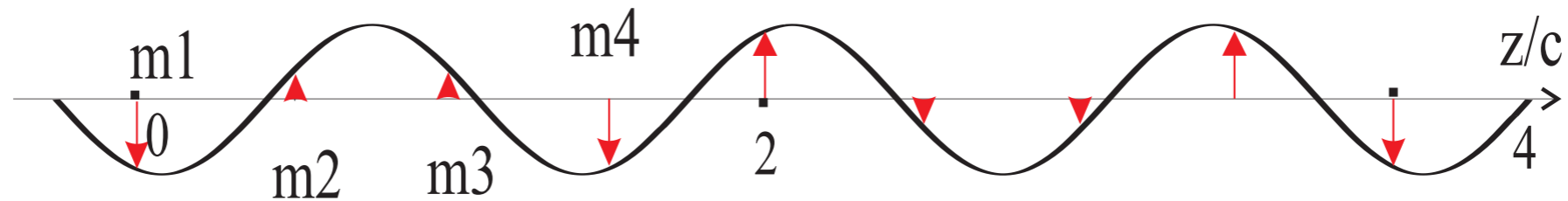


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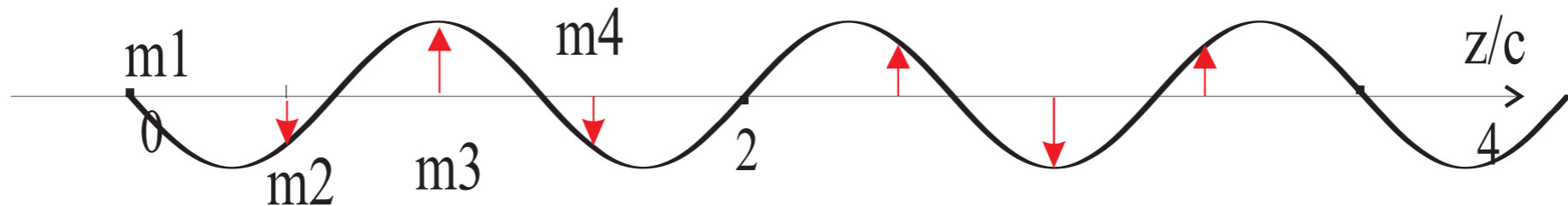
Spin/atom magnetic moment  $\rightarrow$   $\mathbf{S} = \mathbf{S}_0 \cos(2\pi kz + \phi), k = \frac{3}{4}$   $\leftarrow$  Amplitude  $I \sim S_0^2 + S_0 F \cos(\phi)$



$\phi = 7\pi/8$



$\phi = \pi/2$



The phase  $\phi$  is not accessible and the magnetic moments on the atoms cannot be determined.

powder diffraction, + and -

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- + no difficult corrections, e.g. extinction, crystal shape
- peak overlapping/multiplicity in powder data puts a restriction on the determination of spin direction
- small spin components ( $\sim 10^{-1} \mu_B$ ) are difficult to detect

# Powder neutron diffractometers

European Portal for Neutron Scattering

<http://pathfinder.neutron-eu.net>

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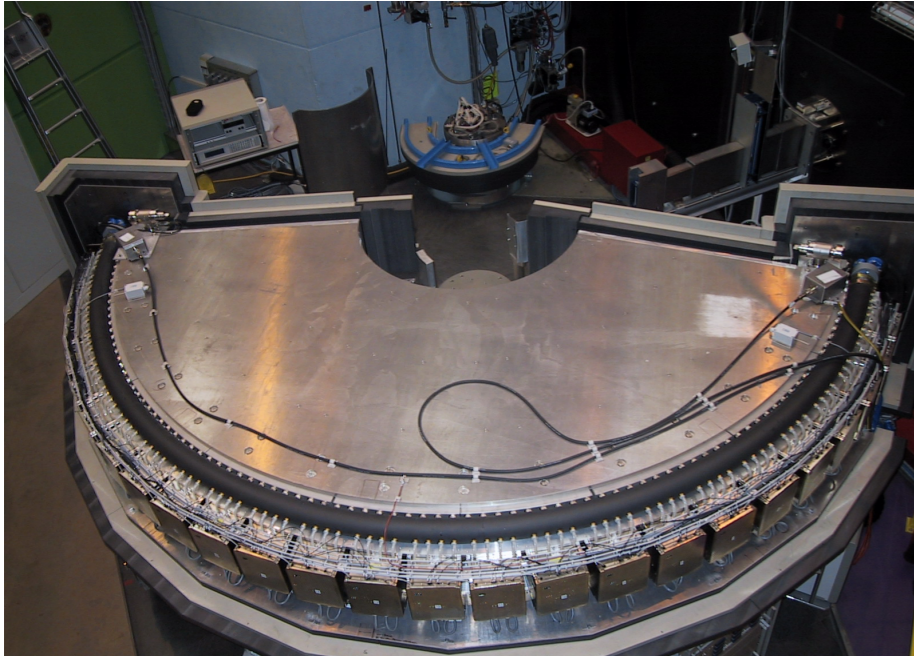
ILL, FR	D20, D2B, DIA
LLB, FR	G41, G42
ISIS, UK	GEM, HRPD, PEARL
FRM-II, DE	SPODI
FLNP/Dubna, RU	HRFD, DN2, DNI2
SINQ/PSI, CH	DMC, HRPT, POLDI

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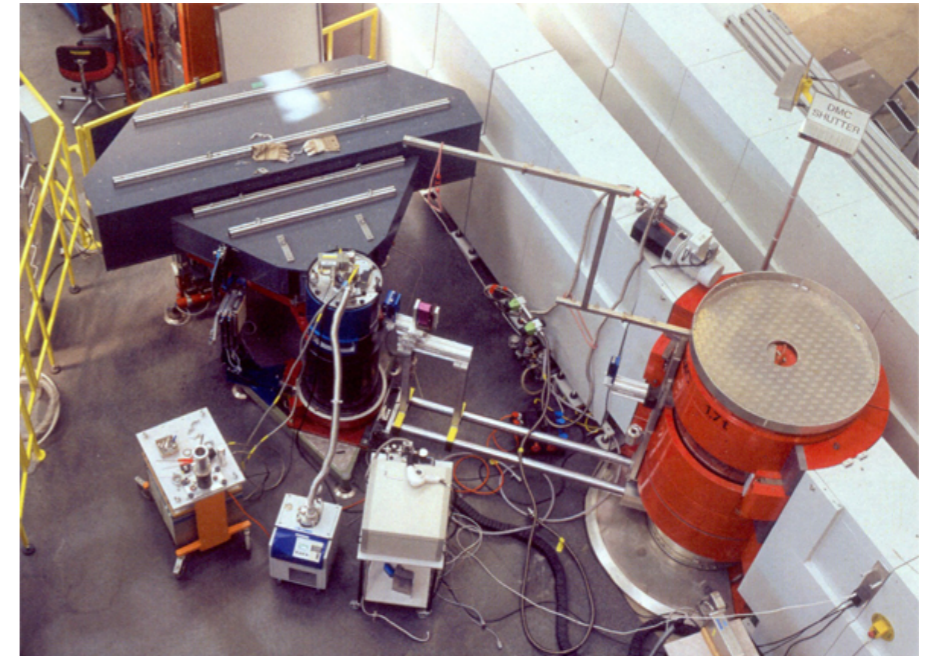
<http://pathfinder.neutron-eu.net>

# Powder ND at SINQ/PSI

**HRPT** - High Resolution Powder  
Diffractometer for Thermal Neutrons at SINQ

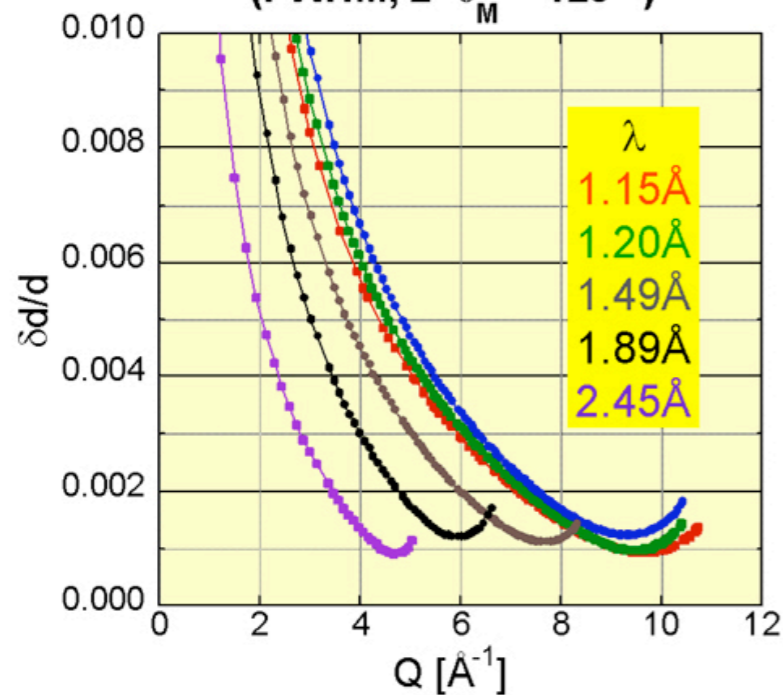


**DMC** - cold neutron powder diffractometer

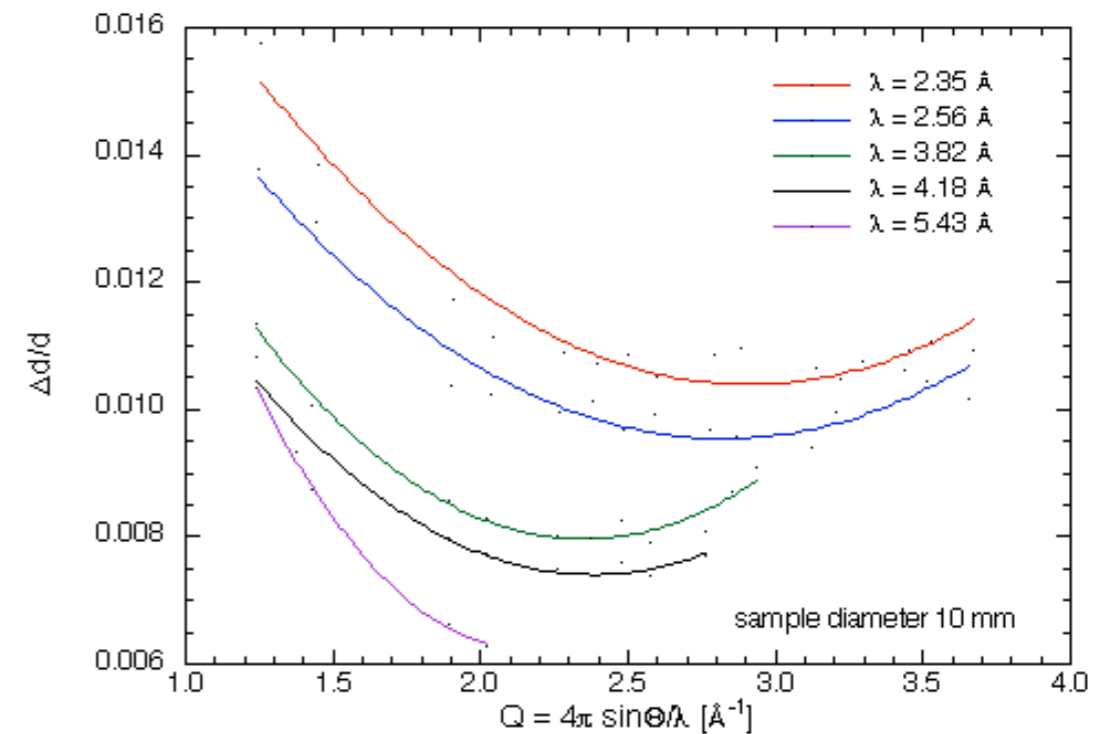


HRPT RESOLUTION FUNCTIONS

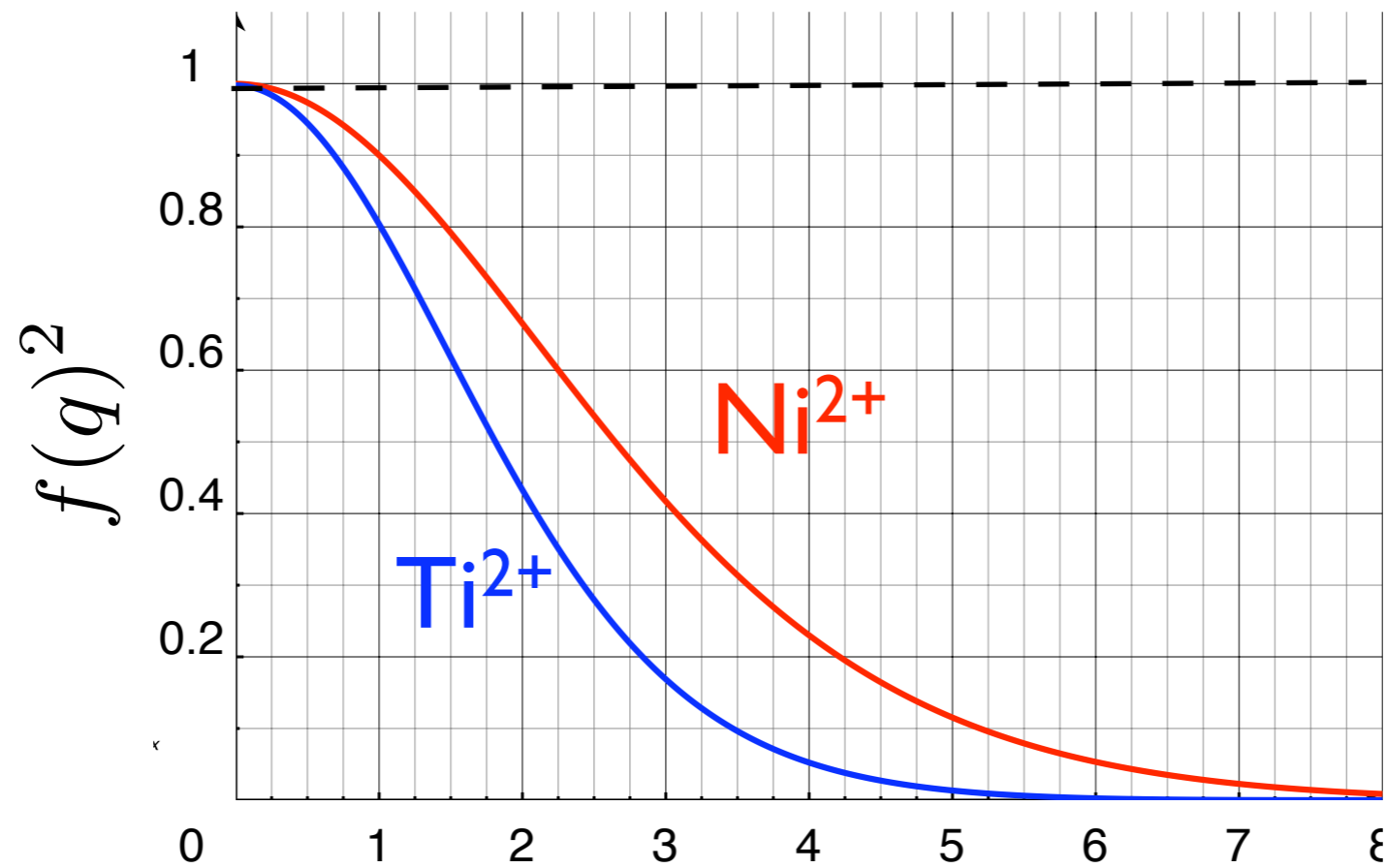
(FWHM,  $2\theta_M = 120^\circ$ )



DMC: experimental resolution functions  $\Delta d/d(Q, \lambda)$

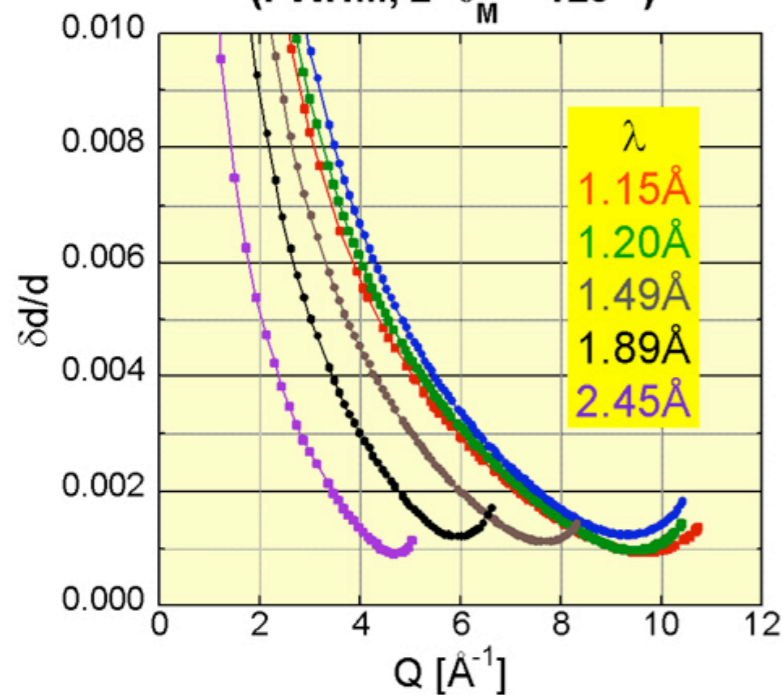


# Powder ND at SINQ/PSI

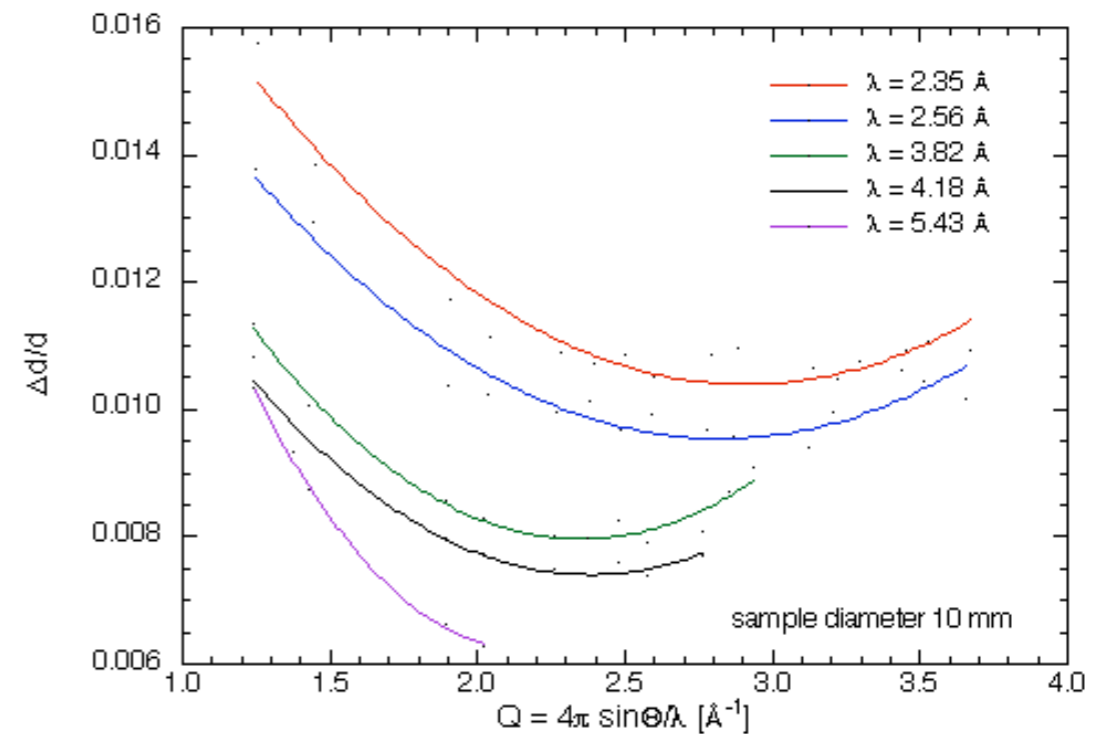


HRPT RESOLUTION FUNCTIONS

(FWHM,  $2\theta_M = 120^\circ$ )



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Examples of magnetic structures.

Propagation vector  $k$

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## Propagation vector **k**



Magnetic moment  
is a real quantity!

$$\mathbf{S}(\mathbf{r}_j) = \text{Re}(\mathbf{S}_0 e^{2\pi i \mathbf{r}_j \mathbf{k}}) = \frac{1}{2} (\mathbf{S}_0 e^{+2\pi i \mathbf{r}_j \mathbf{k}} + c.c.)$$

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☹ Amplitude is complex  $\mathbf{S}_0 = \mathbf{S}_x e^{i\phi_x} + \mathbf{S}_y e^{i\phi_y} + \mathbf{S}_z e^{i\phi_z}$

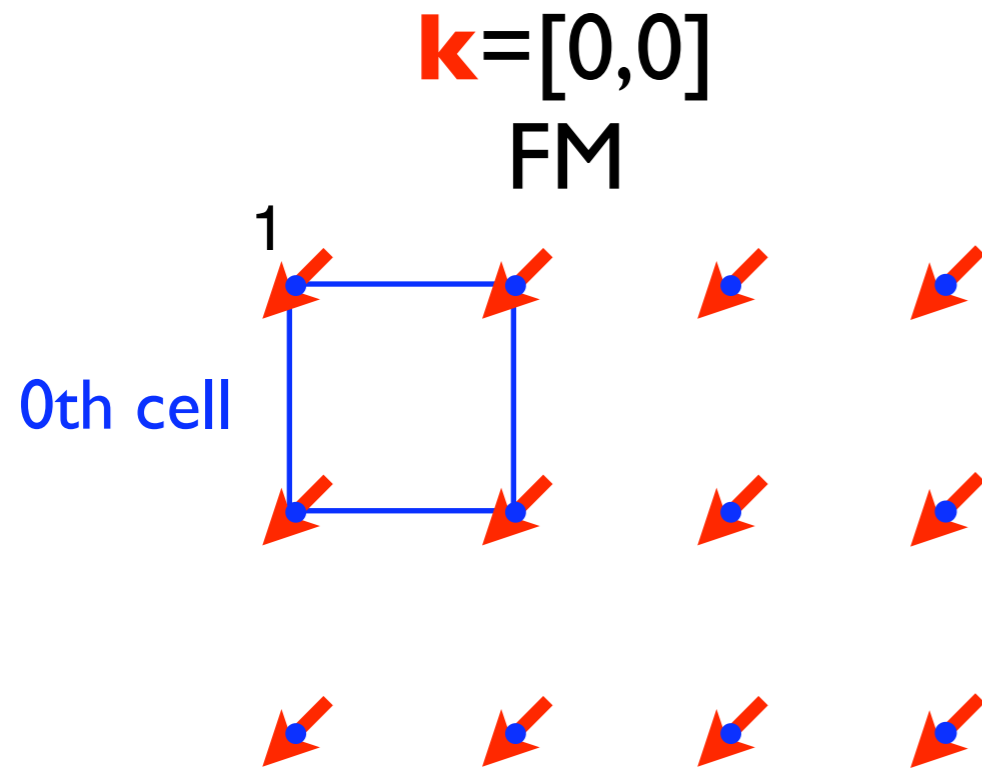


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$$\mathbf{S}_{01} = \mathbf{S}_x + \mathbf{S}_y$$

# Examples of magnetic structures.

## Propagation vector $\mathbf{k}$



Magnetic moment is a real quantity!

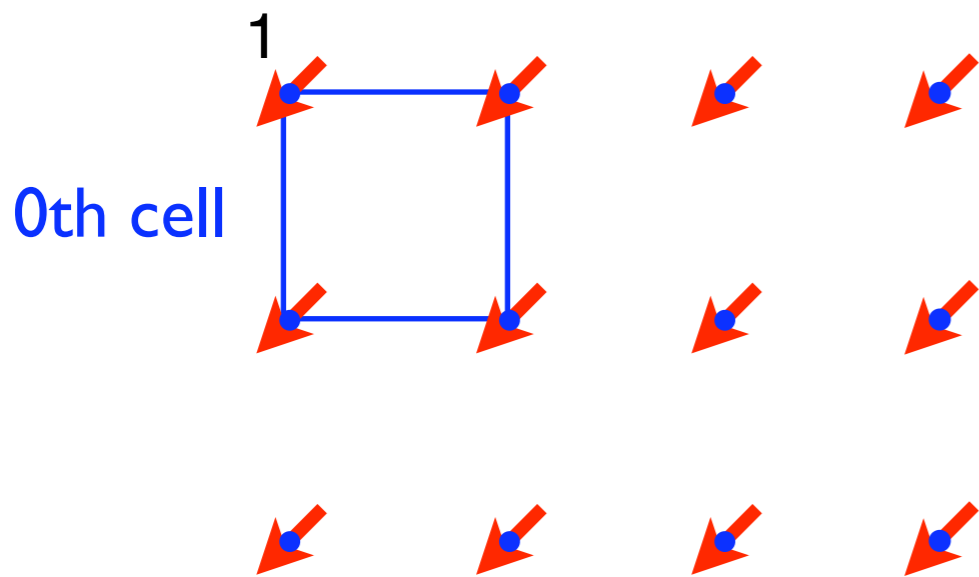
$$\mathbf{S}(\mathbf{r}_j) = \text{Re}(\mathbf{S}_0 e^{2\pi i \mathbf{r}_j \mathbf{k}}) = \frac{1}{2} (\mathbf{S}_0 e^{+2\pi i \mathbf{r}_j \mathbf{k}} + c.c.)$$



Amplitude is complex

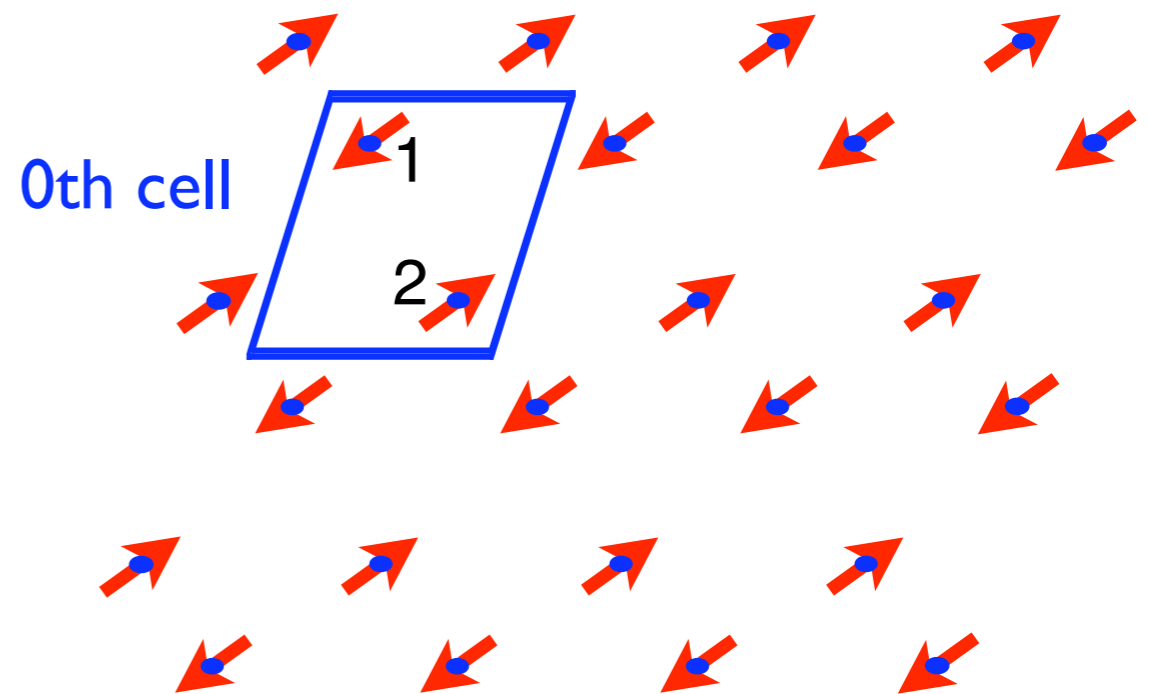
$$\mathbf{S}_0 = \mathbf{S}_x e^{i\phi_x} + \mathbf{S}_y e^{i\phi_y} + \mathbf{S}_z e^{i\phi_z}$$

$\mathbf{k}=[0,0]$   
FM



$$\mathbf{S}_{01} = \mathbf{S}_x + \mathbf{S}_y$$

$\mathbf{k}=[0,0]$   
AFM



$$\mathbf{S}_{01} = \mathbf{S}_x + \mathbf{S}_y$$

$$\mathbf{S}_{02} = -\mathbf{S}_{01}$$

# Examples of magnetic structures.

## Propagation vector $\mathbf{k} \neq 0$

Magnetic moment is a real quantity!

$$\mathbf{S}(\mathbf{r}_j) = \frac{1}{2} (\mathbf{S}_0 e^{+2\pi i \mathbf{r}_j \cdot \mathbf{k}} + \mathbf{S}_0^* e^{-2\pi i \mathbf{r}_j \cdot \mathbf{k}})$$

Amplitude is complex

$$\mathbf{S}_0 = \mathbf{S}_x e^{i\phi_x} + \mathbf{S}_y e^{i\phi_y} + \mathbf{S}_z e^{i\phi_z}$$

# Examples of magnetic structures.

## Propagation vector $\mathbf{k} \neq 0$

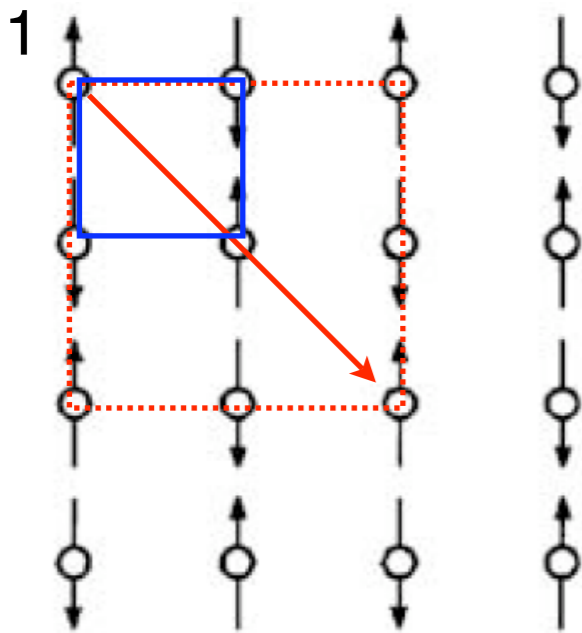
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Amplitude is complex

$$\mathbf{S}_0 = \mathbf{S}_x e^{i\phi_x} + \mathbf{S}_y e^{i\phi_y} + \mathbf{S}_z e^{i\phi_z}$$

$\mathbf{k} = [1/2, 1/2]$  AFM



$$\mathbf{S}_{01} = \mathbf{S}_y$$

# Examples of magnetic structures.

## Propagation vector $\mathbf{k} \neq 0$

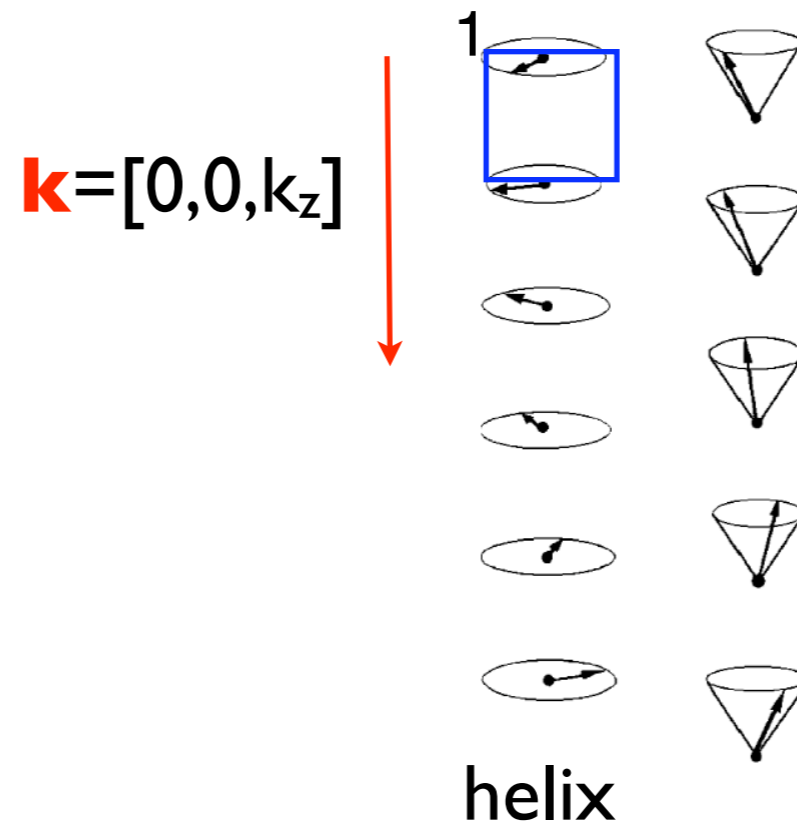
Magnetic moment is a real quantity!

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Amplitude is complex

$$\mathbf{S}_0 = \mathbf{S}_x e^{i\phi_x} + \mathbf{S}_y e^{i\phi_y} + \mathbf{S}_z e^{i\phi_z}$$

modulated (in)commensurate



# Examples of magnetic structures.

## Propagation vector $\mathbf{k} \neq 0$

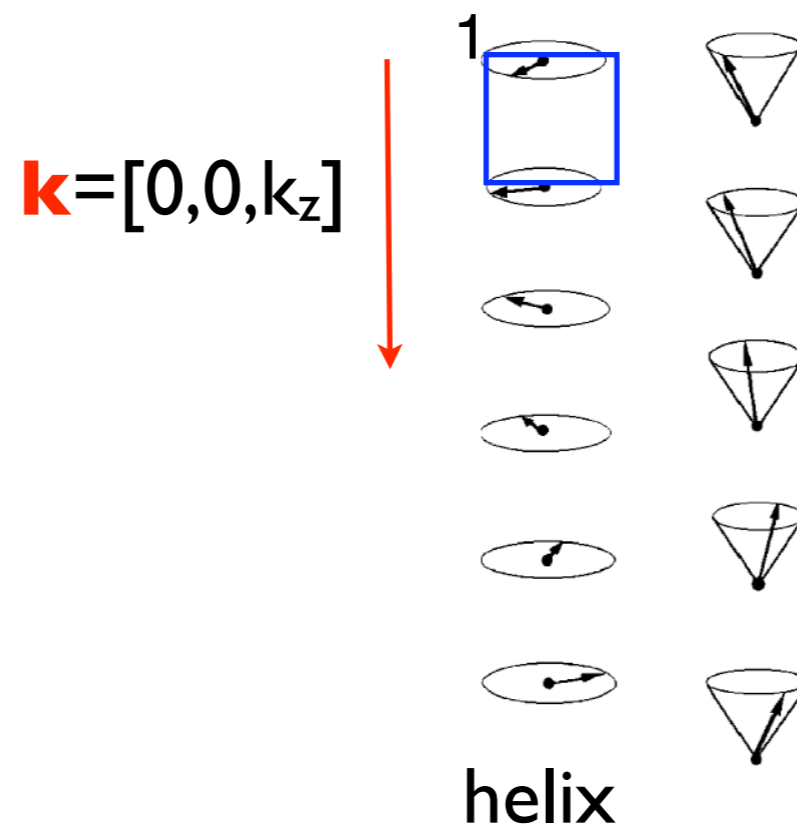
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$$\mathbf{S}(\mathbf{r}_j) = \frac{1}{2} (\mathbf{S}_0 e^{+2\pi i \mathbf{r}_j \cdot \mathbf{k}} + \mathbf{S}_0^* e^{-2\pi i \mathbf{r}_j \cdot \mathbf{k}})$$

Amplitude is complex

$$\mathbf{S}_0 = \mathbf{S}_x e^{i\phi_x} + \mathbf{S}_y e^{i\phi_y} + \mathbf{S}_z e^{i\phi_z}$$

modulated (in)commensurate



$$\mathbf{S}_{01} = \mathbf{S}_x + \mathbf{S}_y e^{\frac{i\pi}{2}} = \mathbf{S}_x + i\mathbf{S}_y$$

# Examples of magnetic structures.

## Propagation vector $\mathbf{k} \neq 0$

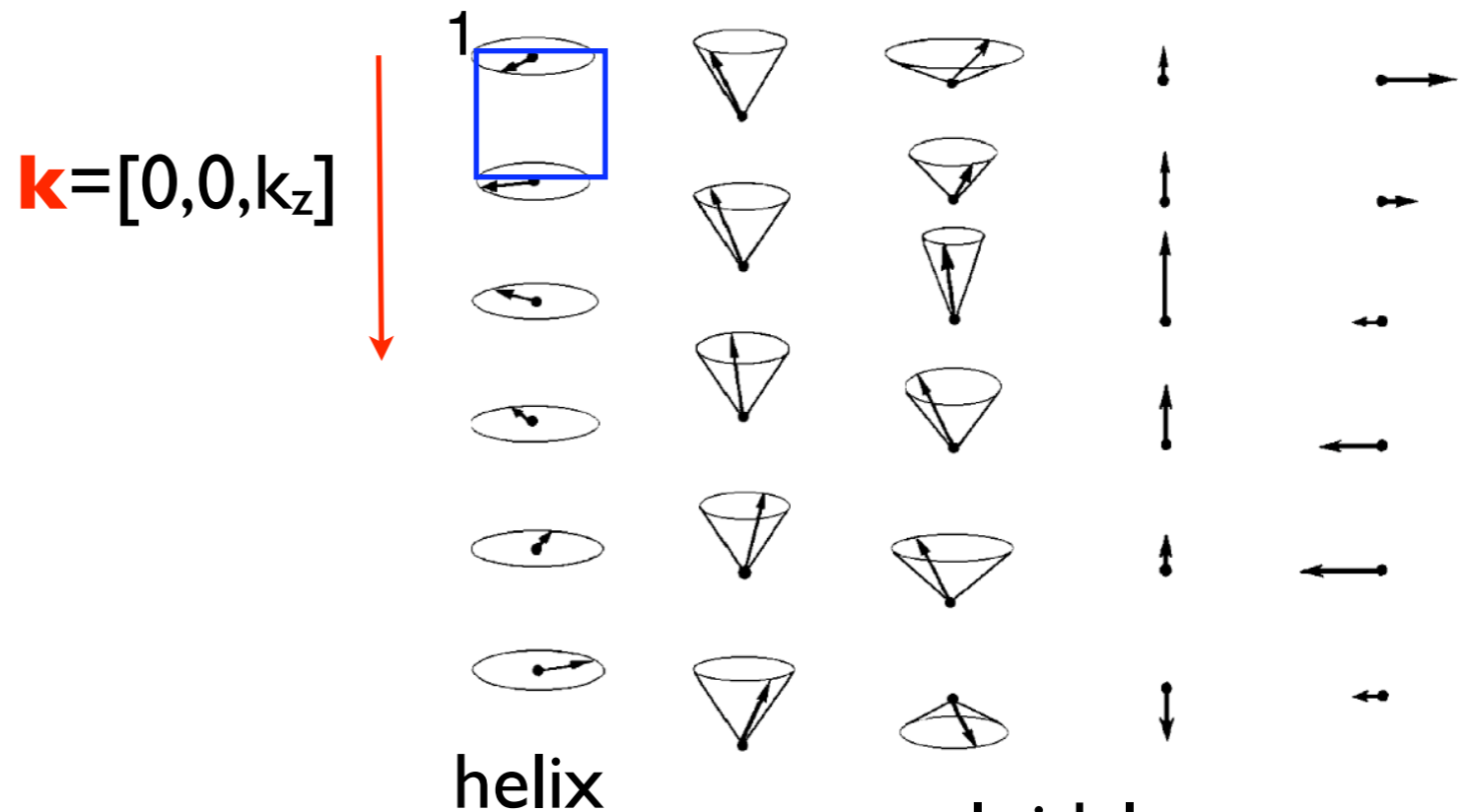
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Amplitude is complex

$$\mathbf{S}_0 = \mathbf{S}_x e^{i\phi_x} + \mathbf{S}_y e^{i\phi_y} + \mathbf{S}_z e^{i\phi_z}$$

modulated (in)commensurate



$$\mathbf{S}_{01} = \mathbf{S}_x + \mathbf{S}_y e^{\frac{i\pi}{2}} = \mathbf{S}_x + i\mathbf{S}_y$$

cycloidal spiral

SDW

$$\mathbf{S}_{01} = \mathbf{S}_x + i\mathbf{S}_y + \mathbf{S}_z e^{i\phi_z}$$

# Example of complex magnetic structure

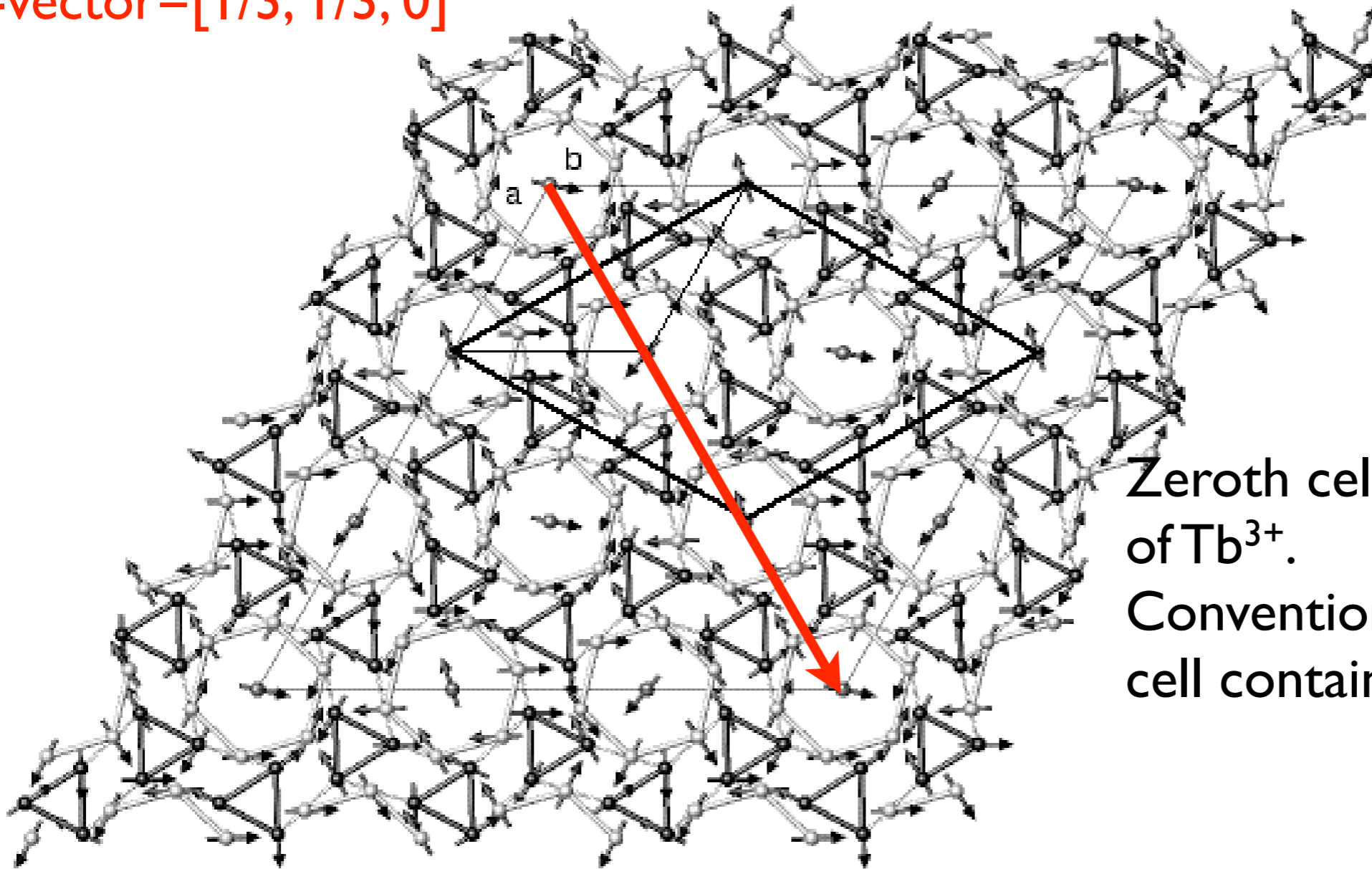


# Example of complex magnetic structure

Antiferromagnetic three sub-lattice ordering in  $\text{Tb}_{14}\text{Au}_{51}$

P6/m

k-vector =  $[\frac{1}{3}, \frac{1}{3}, 0]$



Zeroth cell contains 14 spins of  $\text{Tb}^{3+}$ .

Conventional magnetic unit cell contains 126 spins of  $\text{Tb}^{3+}$ .

# Analysis of magnetic neutron diffraction: computer programs and tutorials/notes

- INDEXING, K-VECTOR: programs distributed with [FullProf Suite](#) [1]
- SYMMETRY: [Baslreps](#) [1], [SARAh](#) [2], [MODY](#) [3]
- SOLUTION: [FullProf](#) [1] (simulated annealing)
- REFINEMENT: [FullProf](#), [GSAS](#) [4]
- Visualization: [FPStudio](#) [1]

## REFERENCES

1. **Juan Rodríguez-Carvajal (ILL) et al**, <http://www.ill.fr/sites/fullprof/>
2. **Andrew S. Wills (UCL)** [http://www.chem.ucl.ac.uk/people/wills/magnetic\\_structures/magnetic\\_structures.html](http://www.chem.ucl.ac.uk/people/wills/magnetic_structures/magnetic_structures.html)
3. **Wieslawa Sikora et al**, <http://www.ftj.agh.edu.pl/~sikora/modyopis.htm>
4. **Bob Von Dreele (ANL) et al**, <http://www.ncnr.nist.gov/programs/crystallography/software/gsas.html>

# Description of magnetic structures

# Description of magnetic structures

## Magnetic symmetry

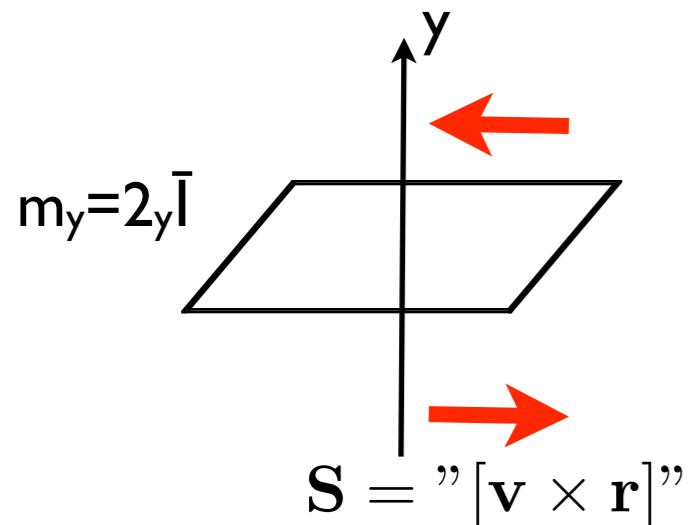
165 | 3D magnetic Shubnikov ( $Sh$ ) space groups.  
Derived from 230 space groups  $G$  and an  
additional element: **spin inversion operator  $R$** .  $Sh$   
groups contain additional 'antieklements'  $g'=(g \cdot R)$ ,  
 $g \in G$  (except  $I$ ) e.g.  $Pn\bar{1}m'$

 **R** = time reversal  
changes **S** to **-S** 

# Description of magnetic structures

## Magnetic symmetry

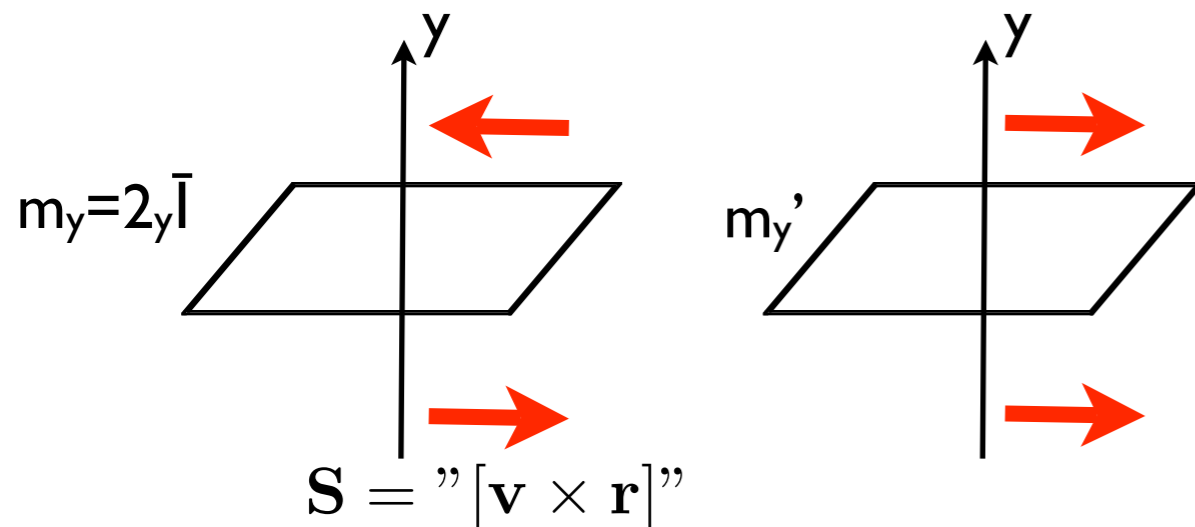
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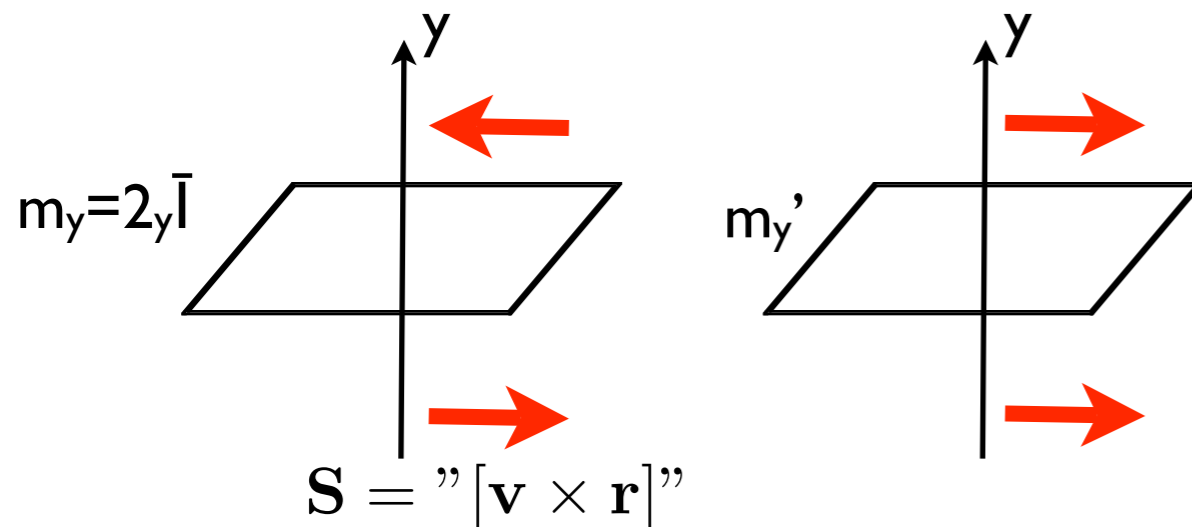
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Disadvantages:

- *Sh* group is not necessarily made from the parent  $G$ . Thus, it is not an ultimate practical tool for obtaining all allowed spin configurations

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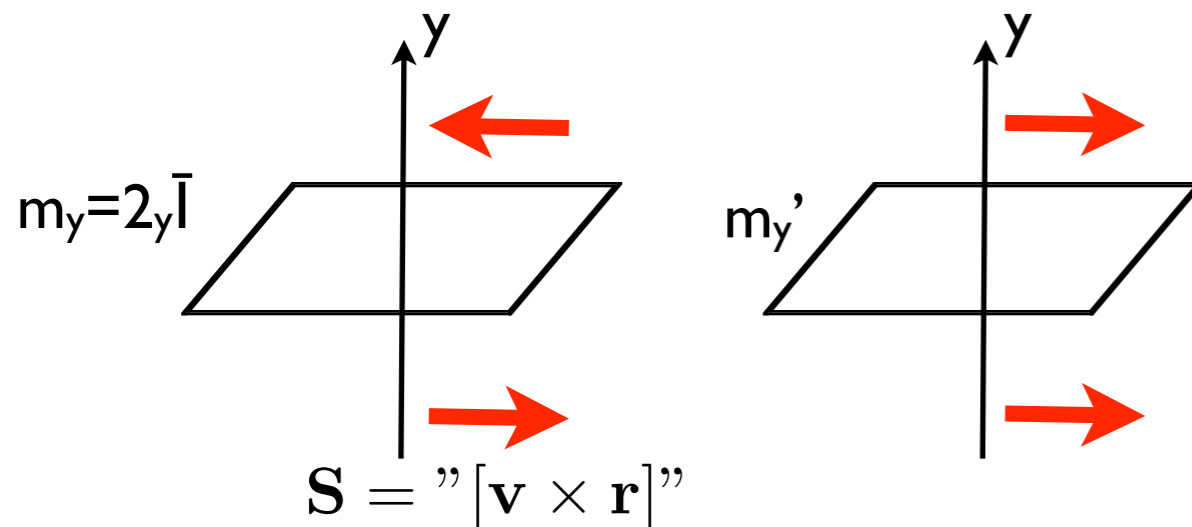
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For example:

$\text{CrCl}_2$  space group:  $Pn\bar{m}$

Possible *Sh* groups derived from the parent space group are:  $Pn\bar{m}$   $Pn'n'm$ ,  $Pn\bar{m}'$ ,  $Pn'n'm'$ ,  $Pn\bar{m}'$ ,  $Pn'n'm'$



Disadvantages:

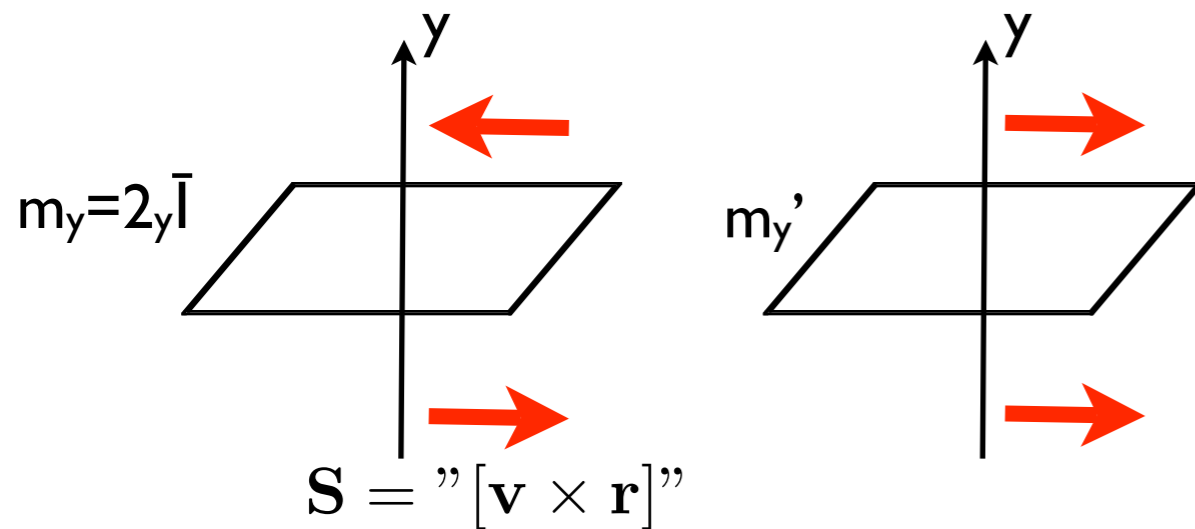
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# Description of magnetic structures

## Magnetic symmetry

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For example:

CrCl<sub>2</sub> space group: *Pnmm*

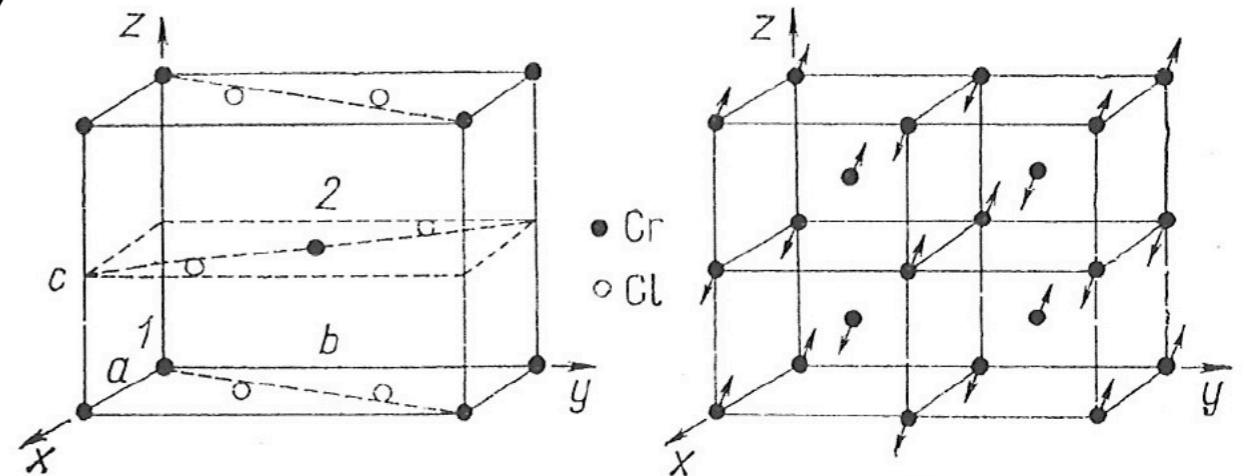
Possible *Sh* groups derived from the parent space group are: *Pnmm* *Pn'n'm*, *Pnmm'*, *Pn'n'm'*, *Pnn'm'*, *Pn'n'm'*

No one describes CrCl<sub>2</sub> magnetic structure

Cr-atoms in 2(a)-position

Cr-spins are antiparallel in 0th cell

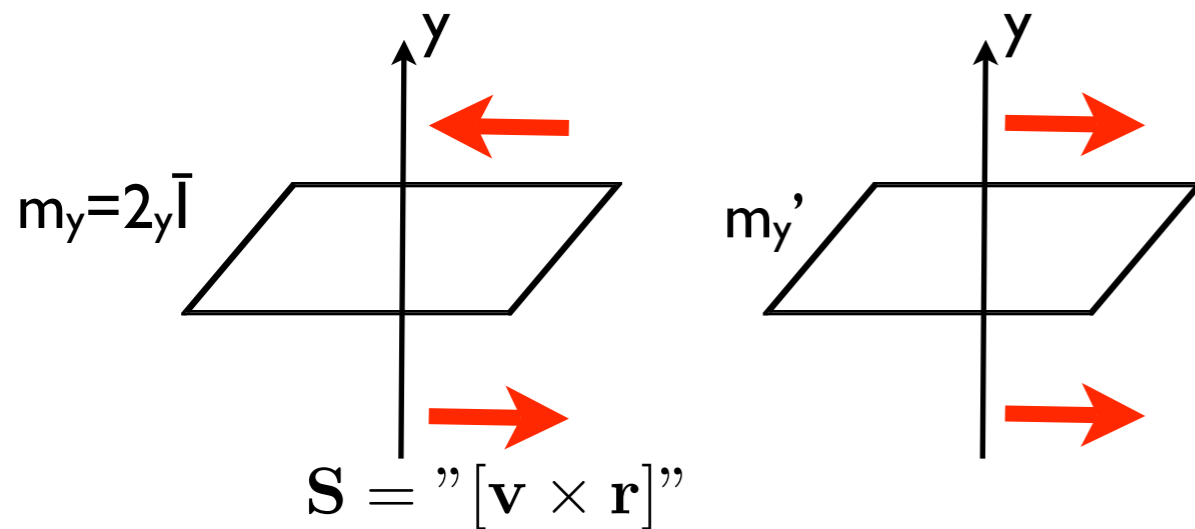
$\mathbf{k}=[0 \ 1/2 \ 1/2]$



# Description of magnetic structures

## Magnetic symmetry

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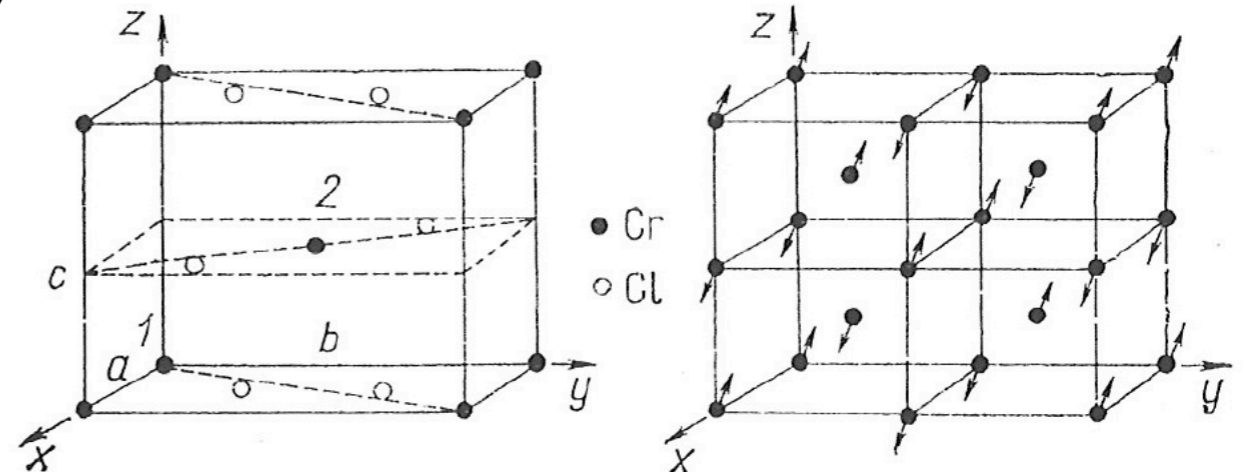
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Cr-atoms in 2(a)-position

Cr-spins are antiparallel in 0th cell

$\mathbf{k}=[0 \ 1/2 \ 1/2]$



One can still find less symmetric *Sh* group

Magnetic symbol

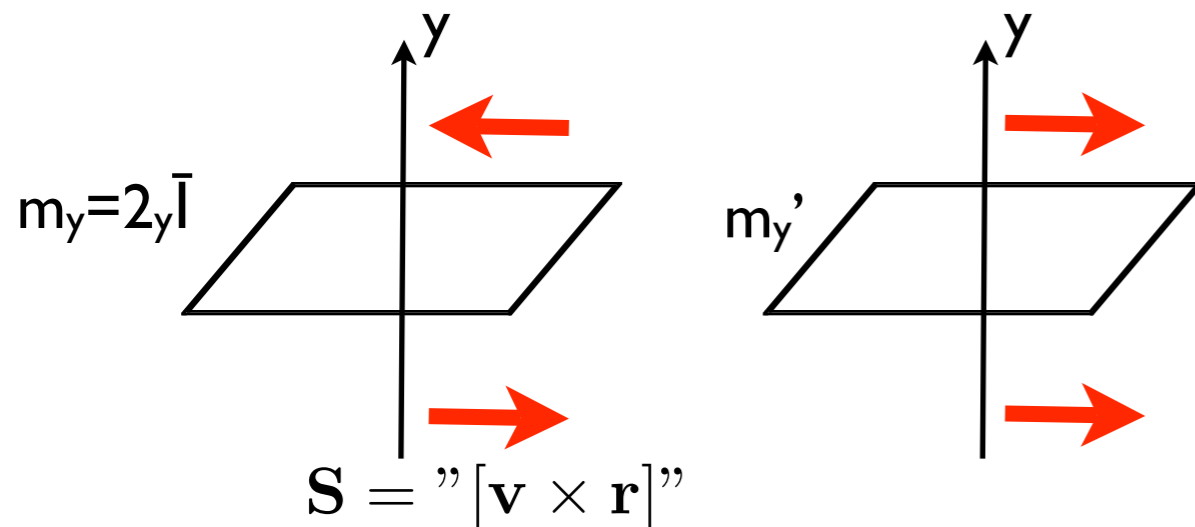
$\{Pnmm; 2(a) \text{ Sh}^7_2 = P_s \bar{1};$

$\mathbf{S}_1=(uvw), \mathbf{S}_2=(-u-v-w)\}$

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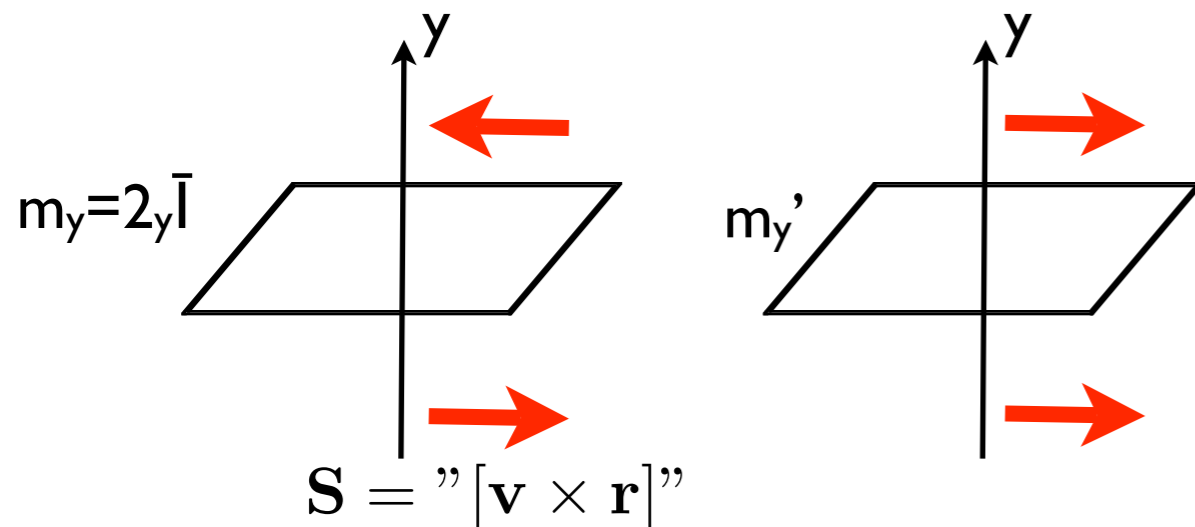
Disadvantages:

- *Sh* group is not necessarily made from the parent  $G$ . Thus, it is not an ultimate practical tool for obtaining all allowed spin configurations
- Do not describe modulated structures. No rotations on non-crystallographic angle - no helix. Linear orthogonal transformations preserve the spin size - no SDW

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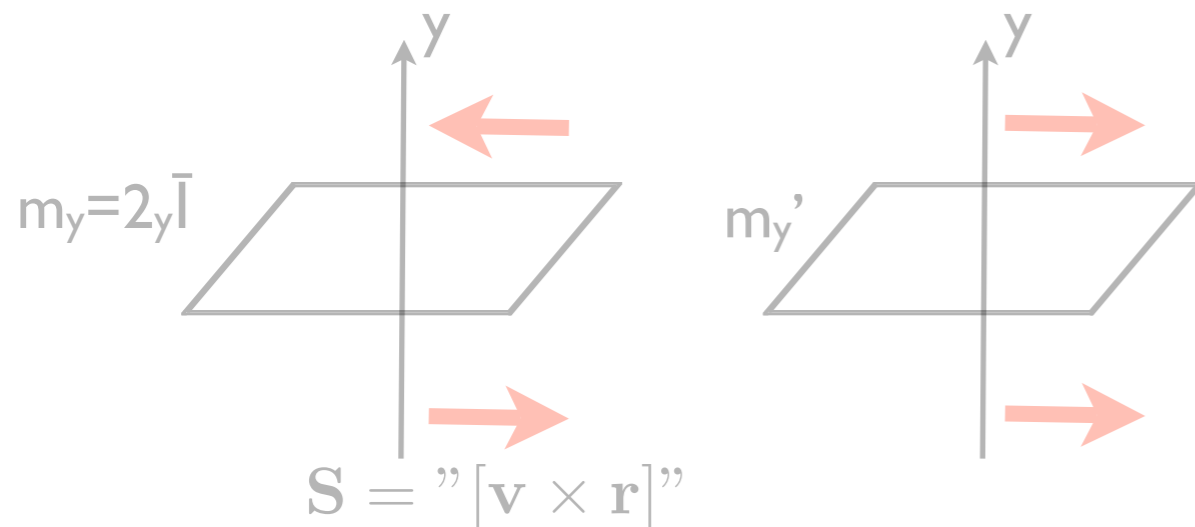
## Representation analysis

A universal technique of finding all possible symmetry adapted spin configurations for the given space group  $G$  and the propagation vector  $\mathbf{k}$ .

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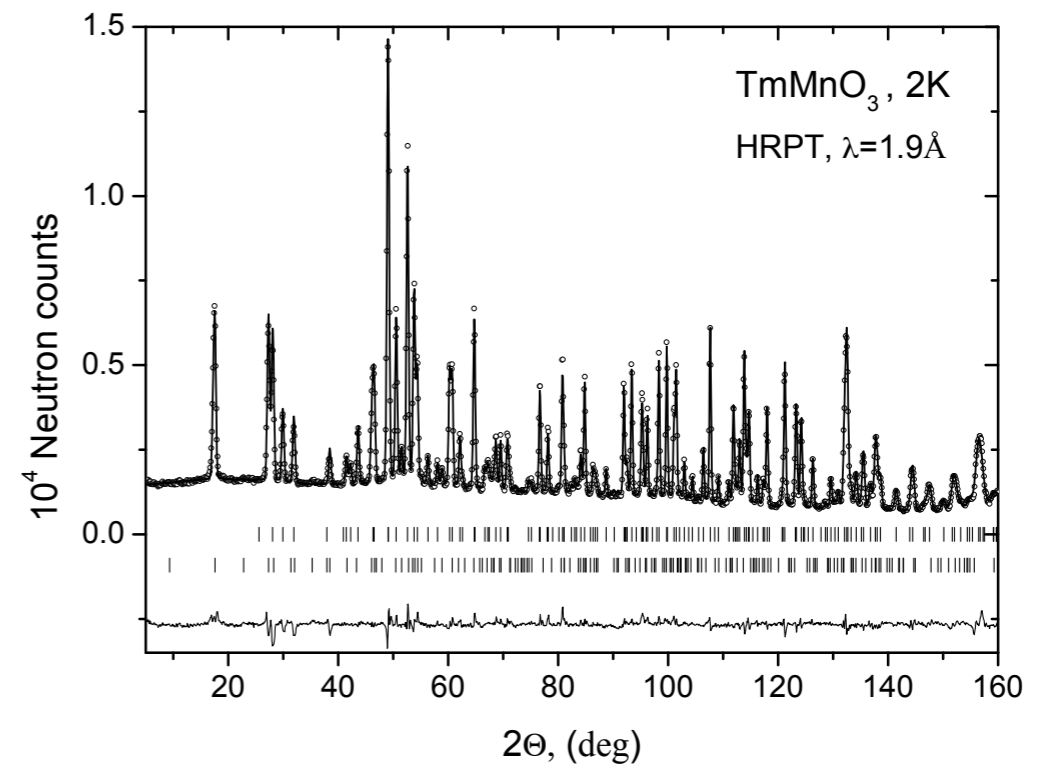
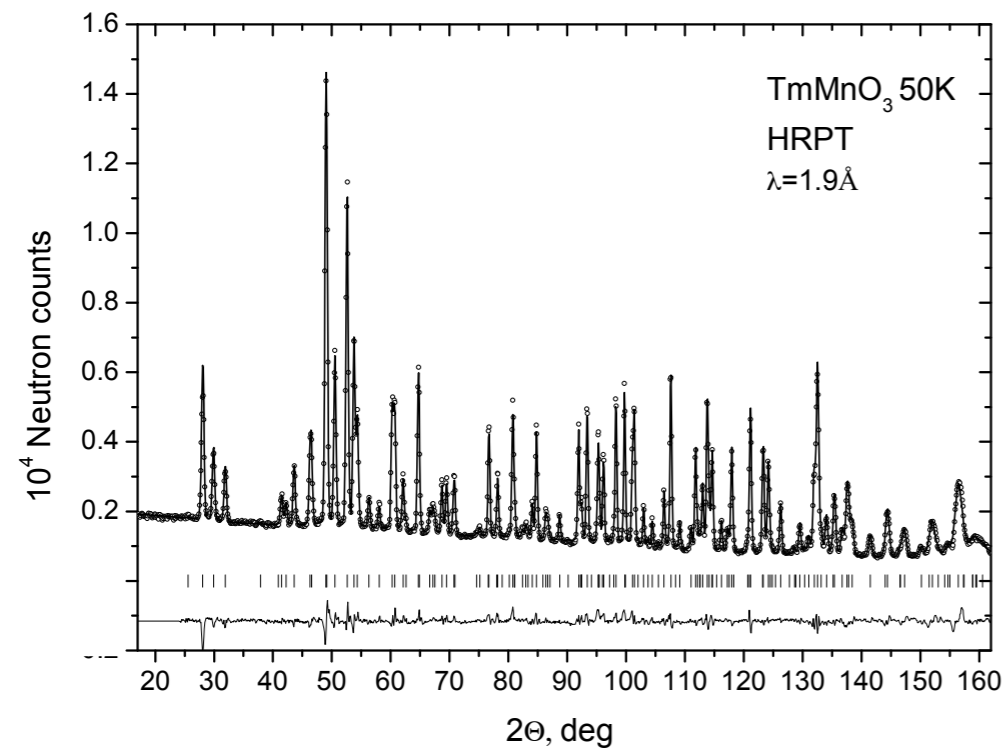
A universal technique of finding all possible symmetry adapted spin configurations for the given space group  $G$  and the propagation vector  $\mathbf{k}$ .

**Case study. Antiferromagnetic order  
in orthorhombic  $\text{TmMnO}_3$**

# Step 1

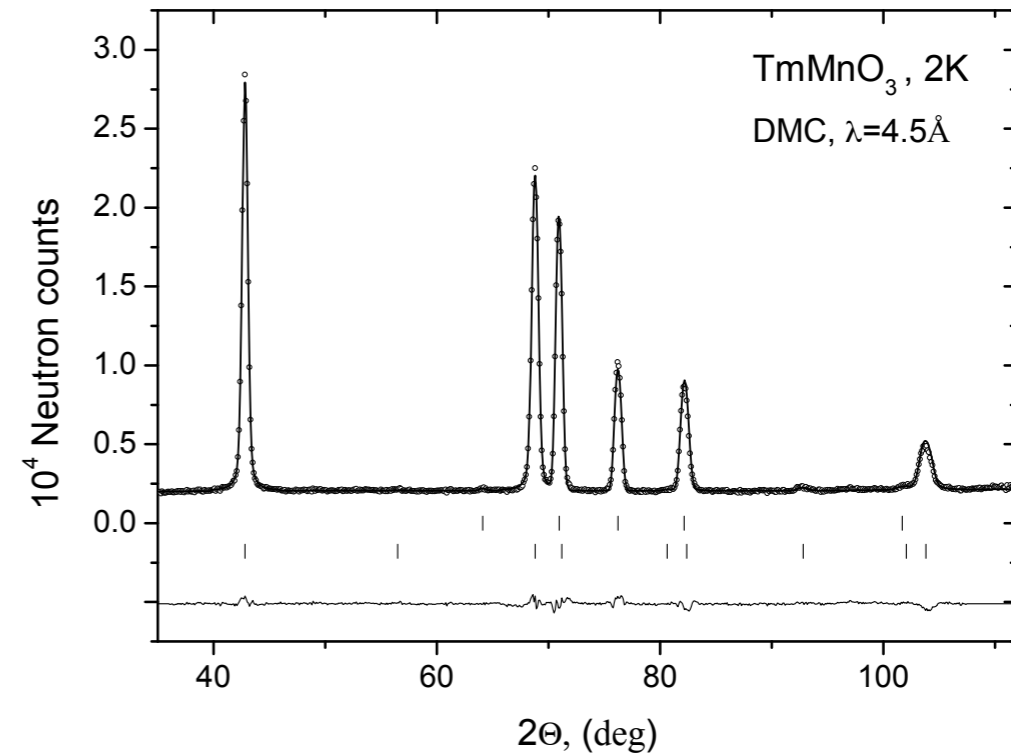
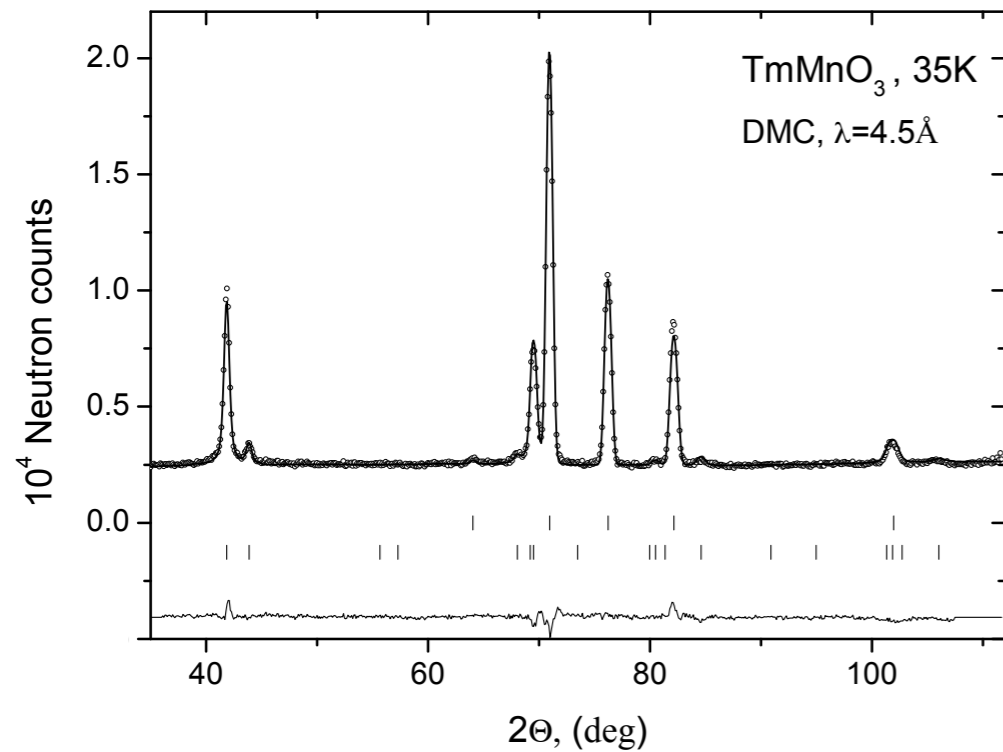
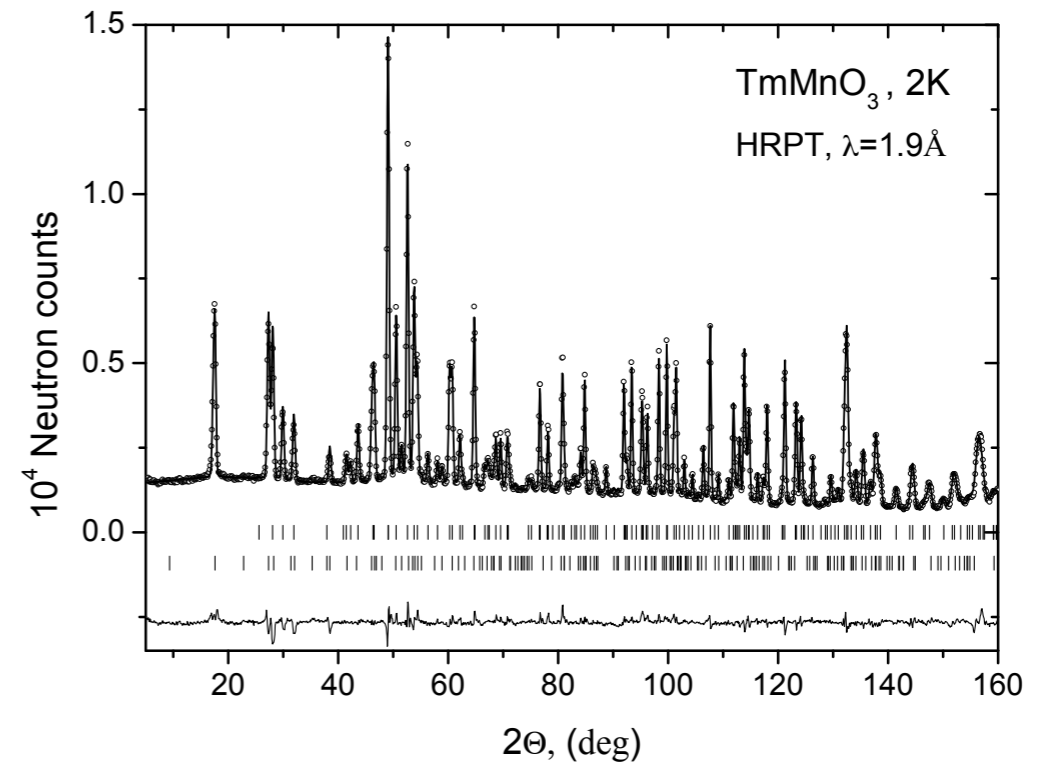
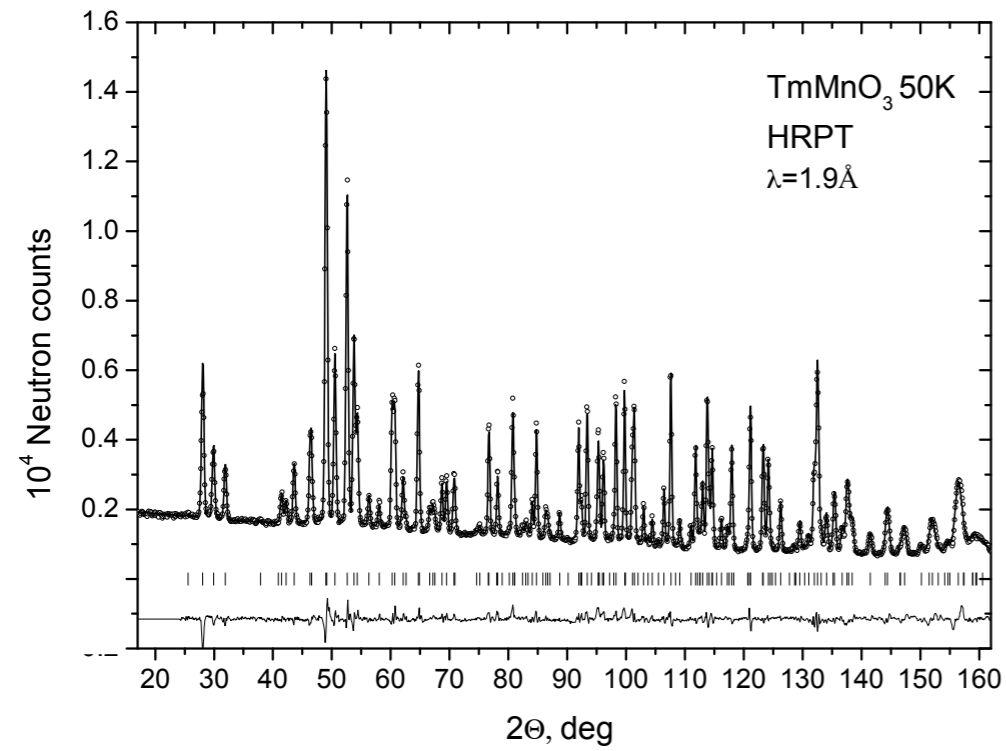
**Experiment. q-range/resolution.**

# Patterns, 1.9Å HRPT and 4.5Å DMC



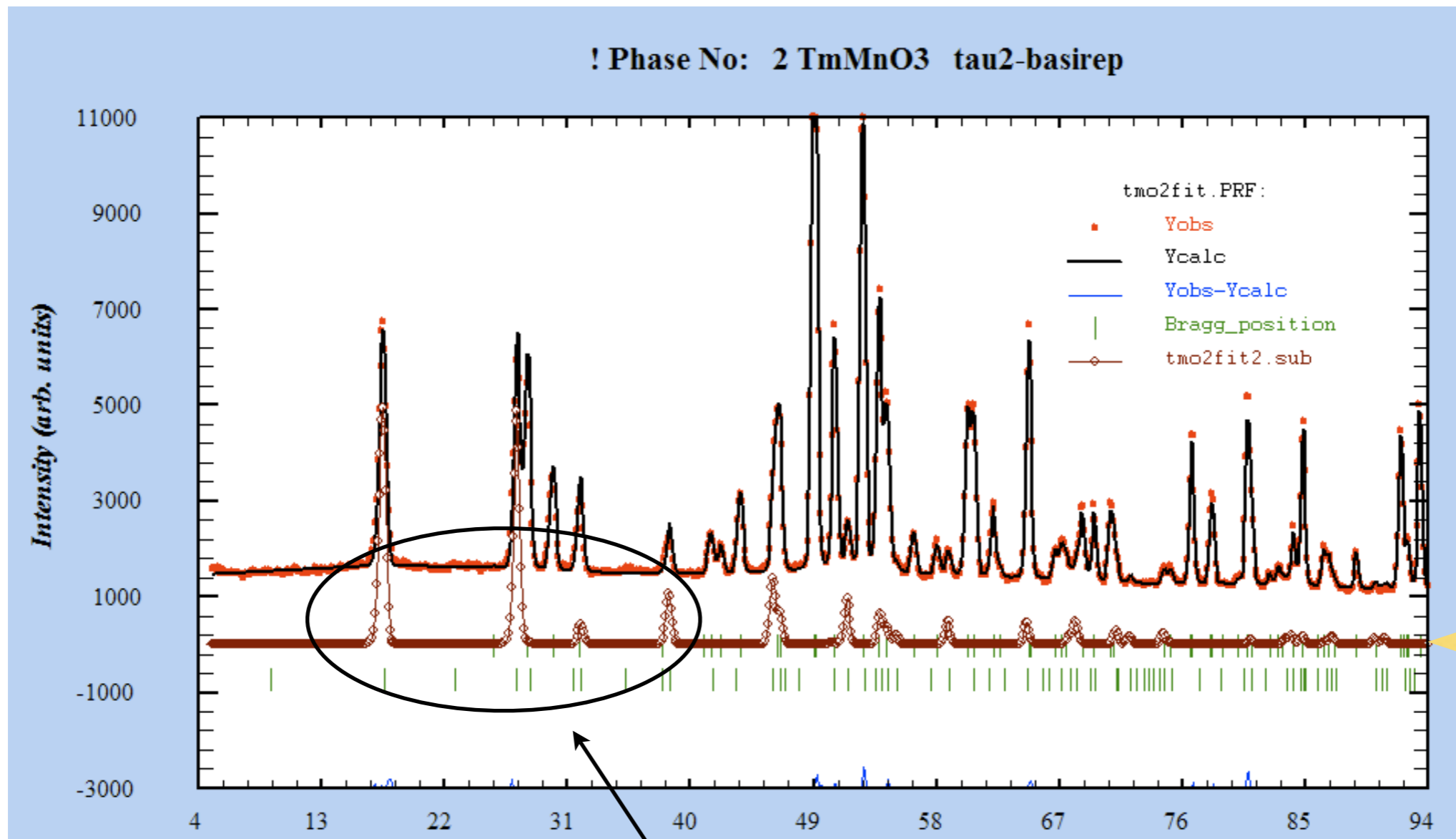


# Patterns, 1.9Å HRPT and 4.5Å DMC



# cf. resolution/q-range

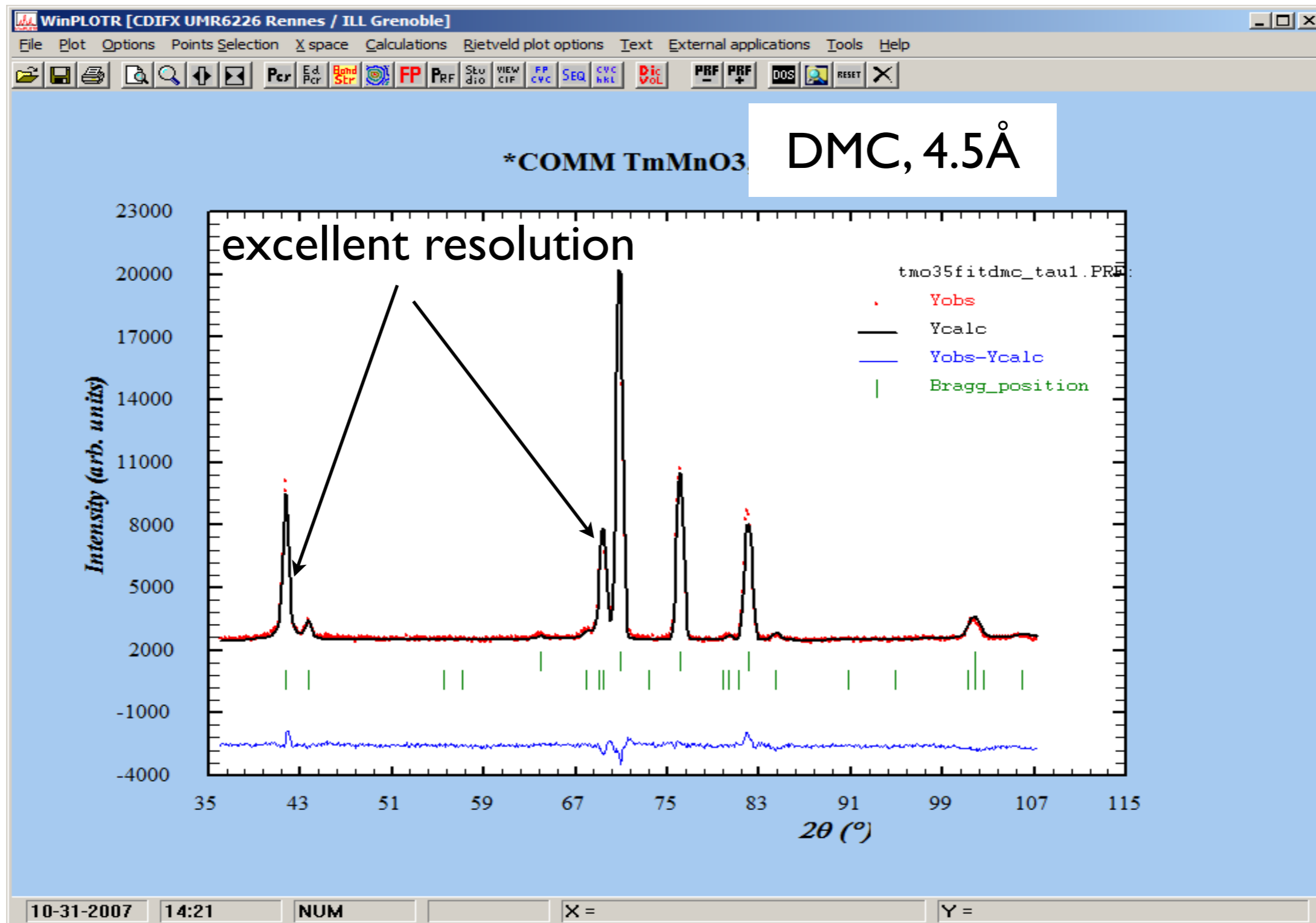
HRPT 1.9Å



magnetic contribution

DMC range at 4.5Å

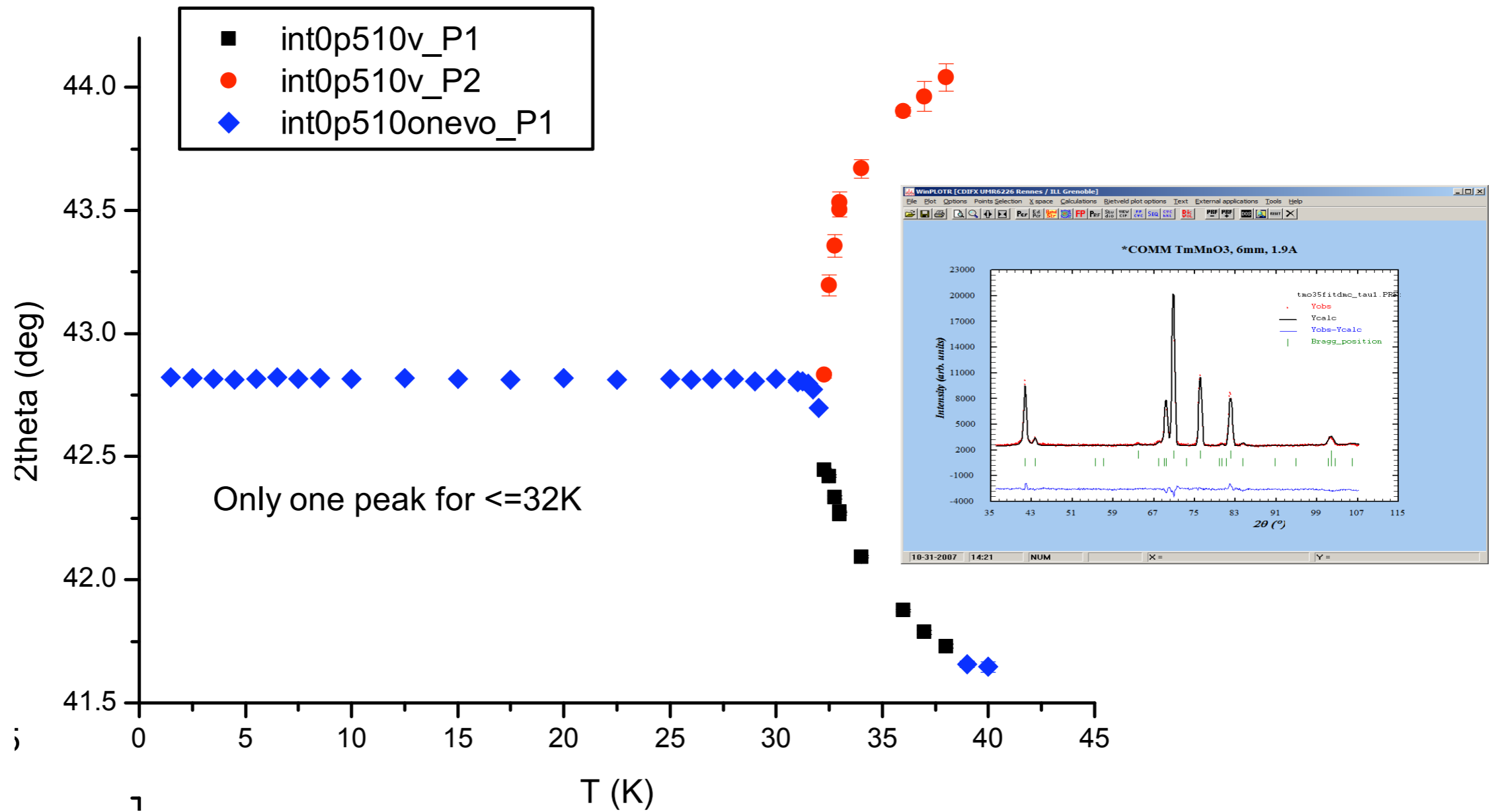
# Cf. resolution/q-range



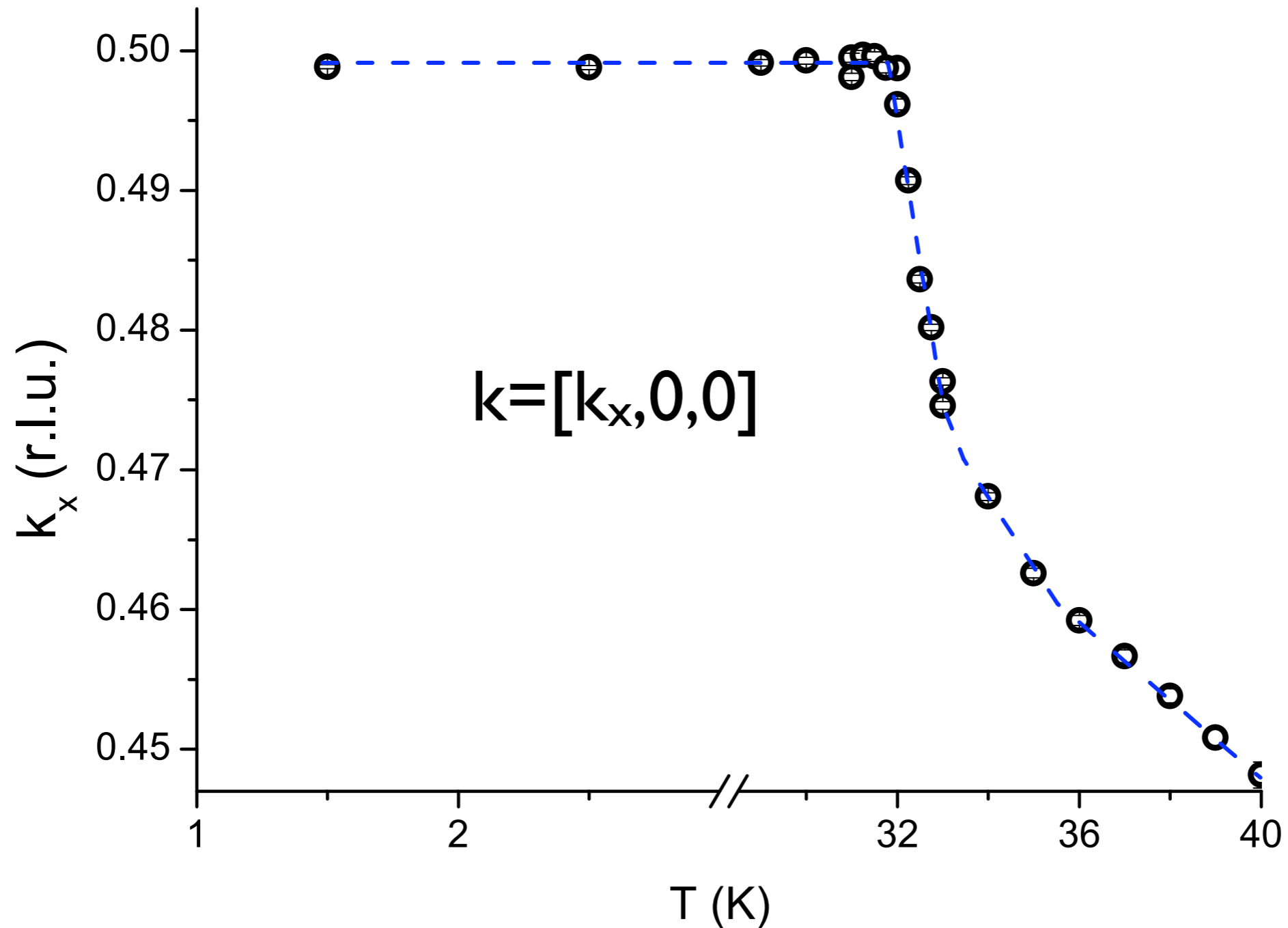
## Step 2

**Finding the propagation vector of  
magnetic structure (k-vector).  
Le Bail profile matching fit.**

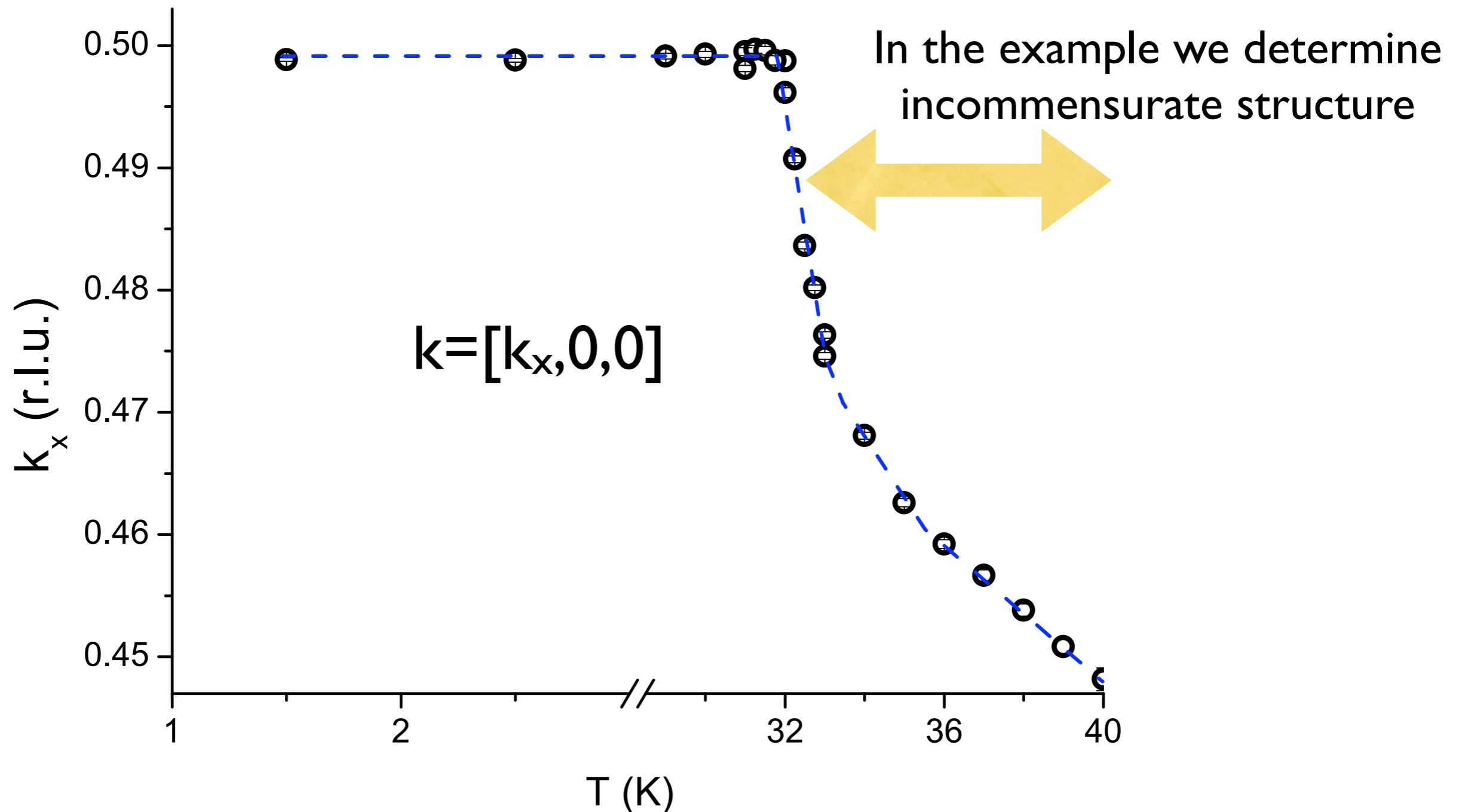
# T-dependence of Bragg peak positions



# Refining the propagation k-vector from profile matching fit



# Refining the propagation k-vector from profile matching fit



# Step 3

**Symmetry analysis.**

**Classifying possible magnetic structures**



# Classifying possible magnetic structures

## k-vector group

Group  $G$ :  $Pnma$ , no.62: 8 symmetry operators

- |               |                            |                            |                                  |
|---------------|----------------------------|----------------------------|----------------------------------|
| (1) $1$       | (2) $2(0, 0, \frac{1}{2})$ | (3) $2(0, \frac{1}{2}, 0)$ | (4) $2(\frac{1}{2}, 0, 0)$       |
| (5) $\bar{1}$ | (6) $a$                    | (7) $m$                    | (8) $n$                          |
| $0, 0, 0$     | $x, y, \frac{1}{4}$        | $x, \frac{1}{4}, z$        | $x, \frac{1}{4}, \frac{1}{4}$    |
|               | $\frac{1}{4}, 0, z$        | $0, y, 0$                  | $n(0, \frac{1}{2}, \frac{1}{2})$ |
|               |                            |                            | $\frac{1}{4}, y, z$              |

# Classifying possible magnetic structures

## **k**-vector group

Group  $G$ :  $Pnma$ , no.62: 8 symmetry operators

(1) $1$	(2) $2(0, 0, \frac{1}{2}) \quad \frac{1}{4}, 0, z$	(3) $2(0, \frac{1}{2}, 0) \quad 0, y, 0$	(4) $2(\frac{1}{2}, 0, 0) \quad x, \frac{1}{4}, \frac{1}{4}$
(5) $\bar{1} \quad 0, 0, 0$	(6) $a \quad x, y, \frac{1}{4}$	(7) $m \quad x, \frac{1}{4}, z$	(8) $n(0, \frac{1}{2}, \frac{1}{2}) \quad \frac{1}{4}, y, z$

Little group  $G_k$ ,  $k=[0.45, 0, 0]=[q, 0, 0]$

Little group of propagation vector  $G_k$  contains only the elements of  $G$  that do not change **k**

# Classifying possible magnetic structures

## k-vector group

Group  $G$ :  $Pnma$ , no.62: 8 symmetry operators

(1) 1	(2) $2(0, 0, \frac{1}{2}) \quad \frac{1}{4}, 0, z$	(3) $2(0, \frac{1}{2}, 0) \quad 0, y, 0$	(4) $2(\frac{1}{2}, 0, 0) \quad x, \frac{1}{4}, \frac{1}{4}$
(5) 1 0,0,0	(6) $a \quad x, y, \frac{1}{4}$	(7) $m \quad x, \frac{1}{4}, z$	(8) $n(0, \frac{1}{2}, \frac{1}{2}) \quad \frac{1}{4}, y, z$

Little group  $G_k$ ,  $k=[0.45, 0, 0]=[q, 0, 0]$

Little group of propagation vector  $G_k$  contains only the elements of  $G$  that do not change  $\mathbf{k}$   
 $P2_1ma$  ( $Pmc2_1$ , 26)

	(1) $x, y, z$	(4) $x + \frac{1}{2}, \bar{y} + \frac{1}{2}, \bar{z} + \frac{1}{2}$	(7) $x, \bar{y} + \frac{1}{2}, z$	(6) $x + \frac{1}{2}, y, \bar{z} + \frac{1}{2}$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_z \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$

# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d

# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d

Permutation representation

# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d

### Permutation representation

element  $g_2$  changes  
atomic position:

$$a \Rightarrow b$$

$$b \Rightarrow a$$

$$c \Rightarrow d$$

$$d \Rightarrow c$$

# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d

### Permutation representation

element  $g_2$  is represented  
by 4x4 matrix

$$\begin{pmatrix} 0100 \\ 1000 \\ 0001 \\ 0010 \end{pmatrix} \begin{pmatrix} a \\ b \\ c \\ d \end{pmatrix} = \begin{pmatrix} b \\ a \\ d \\ c \end{pmatrix}$$

# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d

### Permutation representation

in addition, element  $g_2$  sometimes moves the atom outside of the zerocell.

We have to return the atom back with  $-\mathbf{a}_p$ :

element  $g_2$  is represented by 4x4 matrix

$$\begin{pmatrix} 0100 \\ 1000 \\ 0001 \\ 0010 \end{pmatrix} \begin{pmatrix} a \\ b \\ c \\ d \end{pmatrix} = \begin{pmatrix} b \\ a \\ d \\ c \end{pmatrix}$$

$$\begin{aligned} a &\Rightarrow b \quad (000) \\ b &\Rightarrow a \quad (-100) \\ c &\Rightarrow d \quad (000) \\ d &\Rightarrow c \quad (-100) \end{aligned}$$



# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d

### Permutation representation

in addition, element  $g_2$  sometimes moves the atom outside of the zerocell.

We have to return the atom back with  $-\mathbf{a}_p$ :

element  $g_2$  is represented by 4x4 matrix

$$\begin{pmatrix} 0100 \\ b000 \\ 0001 \\ 00b0 \end{pmatrix} \begin{pmatrix} a \\ b \\ c \\ d \end{pmatrix} = \begin{pmatrix} b \\ a \\ d \\ c \end{pmatrix}$$

$$b = e^{2\pi i(\mathbf{k}\mathbf{a}_p)} \simeq e^{-0.9\pi i}$$

$$a \Rightarrow b \quad (000)$$

$$b \Rightarrow a \quad (-100)$$

$$c \Rightarrow d \quad (000)$$

$$d \Rightarrow c \quad (-100)$$

$$S(\mathbf{r}_j) = S_0 e^{2\pi i \mathbf{r}_j \mathbf{k}}$$

# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d
<b>Permutation representation</b>				
4x4 matrices (P)	$\begin{pmatrix} 1000 \\ 0100 \\ 0010 \\ 0001 \end{pmatrix}$	$\begin{pmatrix} 0100 \\ b000 \\ 0001 \\ 00b0 \end{pmatrix}$	$\begin{pmatrix} 0010 \\ 0001 \\ 1000 \\ 0100 \end{pmatrix}$	$\begin{pmatrix} 0001 \\ 00b0 \\ 0100 \\ b000 \end{pmatrix}$

# Classifying possible magnetic structures

## Magnetic representation

group element	$g^1$	$g^2$	$g^3$	$g^4$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d
	Permutation representation			
4x4 matrices (P)	$\begin{pmatrix} 1000 \\ 0100 \\ 0010 \\ 0001 \end{pmatrix}$	$\begin{pmatrix} 0100 \\ b000 \\ 0001 \\ 00b0 \end{pmatrix}$	$\begin{pmatrix} 0010 \\ 0001 \\ 1000 \\ 0100 \end{pmatrix}$	$\begin{pmatrix} 0001 \\ 00b0 \\ 0100 \\ b000 \end{pmatrix}$

Axial vector (spin) representation

# Classifying possible magnetic structures

## Magnetic representation

group element	$g^1$	$g^2$	$g^3$	$g^4$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d
Permutation representation				
4x4 matrices (P)	$\begin{pmatrix} 1000 \\ 0100 \\ 0010 \\ 0001 \end{pmatrix}$	$\begin{pmatrix} 0100 \\ b000 \\ 0001 \\ 00b0 \end{pmatrix}$	$\begin{pmatrix} 0010 \\ 0001 \\ 1000 \\ 0100 \end{pmatrix}$	$\begin{pmatrix} 0001 \\ 00b0 \\ 0100 \\ b000 \end{pmatrix}$

### Axial vector (spin) representation

For instance:  
 rotational part of element  $g^2$ :  $R(g^2)$  changes  
 atomic spin direction:

$$\begin{pmatrix} S_x \\ S_y \\ S_z \end{pmatrix} = \begin{pmatrix} S_x \\ -S_y \\ -S_z \end{pmatrix}$$

# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d
Permutation representation				
4x4 matrices (P)	$\begin{pmatrix} 1000 \\ 0100 \\ 0010 \\ 0001 \end{pmatrix}$	$\begin{pmatrix} 0100 \\ b000 \\ 0001 \\ 00b0 \end{pmatrix}$	$\begin{pmatrix} 0010 \\ 0001 \\ 1000 \\ 0100 \end{pmatrix}$	$\begin{pmatrix} 0001 \\ 00b0 \\ 0100 \\ b000 \end{pmatrix}$

### Axial vector (spin) representation

For instance:

rotational part of element  $g_2$ :  $R(g_2)$  changes  
atomic spin direction:

element  $g_2$  is represented  
by 3x3 matrix

$$R(g_2) \times \det(R) \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} S_x \\ S_y \\ S_z \end{pmatrix} = \begin{pmatrix} S_x \\ -S_y \\ -S_z \end{pmatrix}$$

# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
rotation+ translation	$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$	$2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$	$m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d
Permutation representation				
4x4 matrices (P)	$\begin{pmatrix} 1000 \\ 0100 \\ 0010 \\ 0001 \end{pmatrix}$	$\begin{pmatrix} 0100 \\ b000 \\ 0001 \\ 00b0 \end{pmatrix}$	$\begin{pmatrix} 0010 \\ 0001 \\ 1000 \\ 0100 \end{pmatrix}$	$\begin{pmatrix} 0001 \\ 00b0 \\ 0100 \\ b000 \end{pmatrix}$
Axial vector (spin) representation				
3x3 matrices (A)				
$R(g_2) \times \det(R)$	$\begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix}$	$\begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix}$	$\begin{pmatrix} \bar{1}00 \\ 010 \\ 00\bar{1} \end{pmatrix}$	$\begin{pmatrix} \bar{1}00 \\ 0\bar{1}0 \\ 001 \end{pmatrix}$

# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	a	b	c	d
spin	$S_1$	$S_2$	$S_3$	$S_4$

### Permutation representation

4x4 matrices (P)	$\begin{pmatrix} 1000 \\ 0100 \\ 0010 \\ 0001 \end{pmatrix}$	$\begin{pmatrix} 0100 \\ b000 \\ 0001 \\ 00b0 \end{pmatrix}$	$\begin{pmatrix} 0010 \\ 0001 \\ 1000 \\ 0100 \end{pmatrix}$	$\begin{pmatrix} 0001 \\ 00b0 \\ 0100 \\ b000 \end{pmatrix}$
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### Axial vector (spin) representation

3x3 matrices (A) $R(g_2) \times \det(R)$	$\begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix}$	$\begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix}$	$\begin{pmatrix} \bar{1}00 \\ 010 \\ 00\bar{1} \end{pmatrix}$	$\begin{pmatrix} \bar{1}00 \\ 0\bar{1}0 \\ 001 \end{pmatrix}$
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# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$	
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$	
position number	<b>a</b>	<b>b</b>	<b>c</b>	<b>d</b>	
spin	<b><math>S_1</math></b>	<b><math>S_2</math></b>	<b><math>S_3</math></b>	<b><math>S_4</math></b>	
	Permutation representation				
4x4 matrices (P)	$\begin{pmatrix} 1000 \\ 0100 \\ 0010 \\ 0001 \end{pmatrix}$	$\begin{pmatrix} 0100 \\ b000 \\ 0001 \\ 00b0 \end{pmatrix}$	$\begin{pmatrix} 0010 \\ 0001 \\ 1000 \\ 0100 \end{pmatrix}$	$\begin{pmatrix} 0001 \\ 00b0 \\ 0100 \\ b000 \end{pmatrix}$	$\begin{pmatrix} a \\ b \\ c \\ d \end{pmatrix}$
	Axial vector (spin) representation				
3x3 matrices (A) $R(g_2) \times \det(R)$	$\begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix}$	$\begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix}$	$\begin{pmatrix} \bar{1}00 \\ 010 \\ 00\bar{1} \end{pmatrix}$	$\begin{pmatrix} \bar{1}00 \\ 0\bar{1}0 \\ 001 \end{pmatrix}$	$\begin{pmatrix} S_x \\ S_y \\ S_z \end{pmatrix}$



# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	$a$	$b$	$c$	$d$
spin	$S_1$	$S_2$	$S_3$	$S_4$

Vector spaces

	Permutation representation			
4x4 matrices (P)	$\begin{pmatrix} 1000 \\ 0100 \\ 0010 \\ 0001 \end{pmatrix}$	$\begin{pmatrix} 0100 \\ b000 \\ 0001 \\ 00b0 \end{pmatrix}$	$\begin{pmatrix} 0010 \\ 0001 \\ 1000 \\ 0100 \end{pmatrix}$	$\begin{pmatrix} 0001 \\ 00b0 \\ 0100 \\ b000 \end{pmatrix}$
3x3 matrices (A) $R(g_2) \times \det(R)$	$\begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix}$	$\begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix}$	$\begin{pmatrix} \bar{1}00 \\ 010 \\ 00\bar{1} \end{pmatrix}$	$\begin{pmatrix} \bar{1}00 \\ 0\bar{1}0 \\ 001 \end{pmatrix}$

$$\begin{pmatrix} a \\ b \\ c \\ d \end{pmatrix}$$

$$\begin{pmatrix} S_x \\ S_y \\ S_z \end{pmatrix}$$

Magnetic representation

direct (tensor) product

$$P \otimes A$$

12x12 matrices

# Classifying possible magnetic structures

## Magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$
position number	$a$	$b$	$c$	$d$
spin	$S_1$	$S_2$	$S_3$	$S_4$

Vector spaces

$$\begin{pmatrix} a \\ b \\ c \\ d \end{pmatrix}$$

$$\begin{pmatrix} S_x \\ S_y \\ S_z \end{pmatrix}$$

Permutation representation

4x4 matrices (P)	$\begin{pmatrix} 1000 \\ 0100 \\ 0010 \\ 0001 \end{pmatrix}$	$\begin{pmatrix} 0100 \\ b000 \\ 0001 \\ 00b0 \end{pmatrix}$	$\begin{pmatrix} 0010 \\ 0001 \\ 1000 \\ 0100 \end{pmatrix}$	$\begin{pmatrix} 0001 \\ 00b0 \\ 0100 \\ b000 \end{pmatrix}$
------------------	--	--	--	--

Axial vector (spin) representation

3x3 matrices (A) $R(g_2) \times \det(R)$	$\begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix}$	$\begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix}$	$\begin{pmatrix} \bar{1}00 \\ 010 \\ 00\bar{1} \end{pmatrix}$	$\begin{pmatrix} \bar{1}00 \\ 0\bar{1}0 \\ 001 \end{pmatrix}$
---	---	---	---	---

Magnetic representation

direct (tensor) product  $P \otimes A$   
12x12 matrices

e.g. for group element  $g_2$

$$\begin{pmatrix} 0100 \\ b000 \\ 0001 \\ 00b0 \end{pmatrix} \otimes \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & \bar{1} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \bar{1} & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \bar{b} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & \bar{b} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{1} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{1} \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{b} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{b} & 0 & 0 & 0 \end{pmatrix}$$

# Classifying possible magnetic structures

## Reducing magnetic representation

group element

rotation+ translation

$$\begin{array}{cccc}
 & g_1 & g_2 & g_3 & g_4 \\
 E & \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} & 2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix} & m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix} & m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}
 \end{array}$$

Matrix of magnetic representation acts on  
12 dimensional vector

4x3=12 spin components

E.g. for the element  $g_2$   
 $M=P \otimes A$

$$\begin{pmatrix}
 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & \bar{1} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & \bar{1} & 0 & 0 & 0 & 0 & 0 & 0 \\
 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & \bar{b} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & \bar{b} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{1} & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{1} \\
 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{b} & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{b} & 0 & 0 & 0
 \end{pmatrix}
 \begin{pmatrix}
 s_{x1} \\
 s_{y1} \\
 s_{z1} \\
 s_{x2} \\
 s_{y2} \\
 s_{z2} \\
 s_{x3} \\
 s_{y3} \\
 s_{z3} \\
 s_{x4} \\
 s_{y4} \\
 s_{z4}
 \end{pmatrix}$$

# Classifying possible magnetic structures

## Reducing magnetic representation

group element

rotation+ translation

$$\begin{array}{cccc}
 g_1 & & g_2 & & g_3 & & g_4 \\
 E & \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} & 2_x & \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix} & m_y & \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix} & m_y & \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}
 \end{array}$$

Matrix of magnetic representation acts on  
12 dimensional vector

4x3=12 spin components

E.g. for the element  $g_2$   
 $M=P \otimes A$

$$\begin{pmatrix}
 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & \bar{1} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & \bar{1} & 0 & 0 & 0 & 0 & 0 & 0 \\
 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & \bar{b} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & \bar{b} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{1} & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{1} \\
 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{b} & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{b} & 0 & 0 & 0
 \end{pmatrix}
 \begin{pmatrix}
 S_{x1} \\
 S_{y1} \\
 S_{z1} \\
 S_{x2} \\
 S_{y2} \\
 S_{z2} \\
 S_{x3} \\
 S_{y3} \\
 S_{z3} \\
 S_{x4} \\
 S_{y4} \\
 S_{z4}
 \end{pmatrix}$$

# Classifying possible magnetic structures

## Reducing magnetic representation

group element

rotation+ translation

$$E \begin{pmatrix} 100 \\ 010 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \quad 2_x \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix} \quad m_y \begin{pmatrix} 100 \\ 0\bar{1}0 \\ 001 \end{pmatrix} \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix} \quad m_y \begin{pmatrix} 100 \\ 010 \\ 00\bar{1} \end{pmatrix} \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{pmatrix}$$

Matrix of magnetic representation acts on  
12 dimensional vector

4x3=12 spin components

E.g. for the element  $g_2$   
 $M=P \otimes A$

$$\begin{pmatrix} 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & \bar{1} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \bar{1} & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \bar{b} & 0 & 0 & 0 & 0 & a & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & \bar{b} & 0 & 0 & a & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{1} & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{1} & 0 \\ 0 & 0 & a & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & a & 0 & 0 & 0 & 0 & 0 & \bar{b} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \bar{b} & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} S_{x1} \\ S_{y1} \\ S_{z1} \\ S_{x2} \\ S_{y2} \\ S_{z2} \\ S_{x3} \\ S_{y3} \\ S_{z3} \\ S_{x4} \\ S_{y4} \\ S_{z4} \end{pmatrix}$$

**Magnetic representation is reducible!**

# Classifying possible magnetic structures

## Reducing magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
Mn-position	$0, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, 0$

Magnetic representation is reducible to a block-diagonal shape that is a direct sum of irreducible square matrices  $\tau_1, \tau_2, \dots$  (dimensions can be from 1 to 6.)

$$\tau_1 \oplus \tau_2 \oplus \tau_3 \oplus \dots = \begin{pmatrix} \tau_1 & 0 & 0 & \dots & 0 \\ 0 & \tau_2 & 0 & \dots & 0 \\ 0 & 0 & \tau_3 & \dots & 0 \\ \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot \\ 0 & 0 & 0 & & \end{pmatrix} \begin{pmatrix} S_{\tau_1} \\ S_{\tau_2} \\ S_{\tau_3} \\ \cdot \\ \cdot \end{pmatrix}$$

# Classifying possible magnetic structures

## Reducing magnetic representation

group element	$g_1$	$g_2$	$g_3$	$g_4$
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Each of these matrices  $\tau_1, \tau_2, \dots$  acts only on a subspace of the 12 spin components.  $S_{\tau_1}, S_{\tau_2}, \dots$  are vectors with dimension of matrix  $\tau_1, \tau_2$

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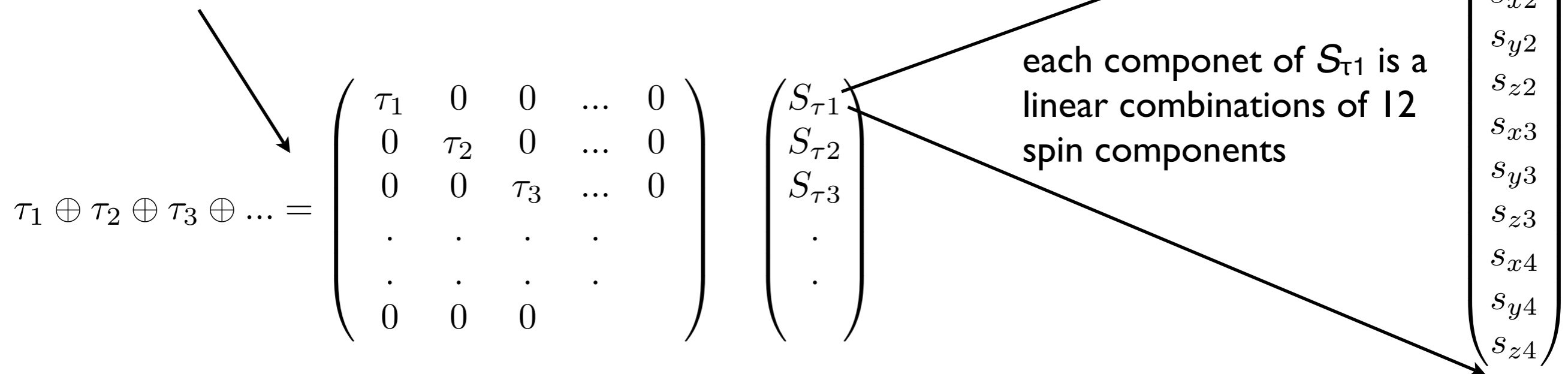
# Classifying possible magnetic structures

## Reducing magnetic representation

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**Landau theory of phase transitions  
says that only one irreducible  
representation is needed to describe  
the structure**

**Why the Landau theory does work for magnetic  
phase transition is a separate topic.**

# Classifying possible magnetic structures basis vectors/functions $S_{\tau_1}, S_{\tau_2}, S_{\tau_3}, \dots$

$Pnma, k=[0.45,0,0]$

Mn in (4a)-position

Magnetic representation is reduced to four  
one-dimensional irreps

$$3\tau_1 \oplus 3\tau_2 \oplus 3\tau_3 \oplus 3\tau_4$$

	$g_1$	$g_2$	$g_3$	$g_4$
$\tau_1$	1	$a$	1	$a$
$\tau_2$	1	$a$	-1	$-a$
$\tau_3$	1	$-a$	1	$-a$
$\tau_4$	1	$-a$	-1	$a$

$$a = e^{\pi i k_x}$$

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$\tau_4$	1	$-a$	-1	$a$

Mn-position

$$0, 0, \frac{1}{2}$$

$$\frac{1}{2}, \frac{1}{2}, 0$$

$$0, \frac{1}{2}, \frac{1}{2}$$

$$\frac{1}{2}, 0, 0$$

1

2

3

4

$$S'_{\tau_3} = +1\mathbf{e}_{1x} - a^*\mathbf{e}_{2x} - 1\mathbf{e}_{3x} + a^*\mathbf{e}_{4x}$$

$$S''_{\tau_3} = +1\mathbf{e}_{1y} + a^*\mathbf{e}_{2y} + 1\mathbf{e}_{3y} + a^*\mathbf{e}_{4y}$$

$$S'''_{\tau_3} = +1\mathbf{e}_{1z} + a^*\mathbf{e}_{2z} - 1\mathbf{e}_{3z} - a^*\mathbf{e}_{4z}$$

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$Pnma, k=[0.45,0,0]$

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Mn-position

$$0, 0, \frac{1}{2}$$

$$\frac{1}{2}, \frac{1}{2}, 0$$

$$0, \frac{1}{2}, \frac{1}{2}$$

$$\frac{1}{2}, 0, 0$$

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$$a = e^{\pi i k_x}$$

Assuming that the phase transition goes according  
to one irreducible representation  $\tau_3$  the spins of all  
four atoms are set only by 3 variables instead of 12!

$$C_1 S'_{\tau_3} + C_2 S''_{\tau_3} + C_3 S'''_{\tau_3}$$

## Steps 3-4 in practice

**Solving/refining the magnetic structure  
by using one irreducible representation**

## Steps 3-4 in practice

**Solving/refining the magnetic structure  
by using one irreducible representation**

- I. construct basis functions for single irreducible representation irrep (use [BasIreps](#), [SARAh](#), [MODY](#))

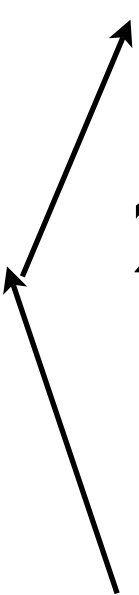
## Steps 3-4 in practice

### Solving/refining the magnetic structure by using one irreducible representation

1. construct basis functions for single irreducible representation irrep (use [BasIreps](#), [SARAh](#), [MODY](#))
2. plug them in the [FULLPROF](#) and try to fit the data. In difficult cases the Monte-Carlo simulated annealing search is required

## Steps 3-4 in practice

### Solving/refining the magnetic structure by using one irreducible representation

1. construct basis functions for single irreducible representation irrep (use **BasIreps**, **SARAh**, **MODY**)
  2. plug them in the **FULLPROF** and try to fit the data. In difficult cases the Monte-Carlo simulated annealing search is required
  3. If the fit is bad go to 1 and choose different irrep. If the fit is good it is still better to sort out all irreps.
- 



# Refinement of the data for $\tau_3$

$$\mathbf{S}(\mathbf{r}) = \frac{1}{2}(C_1 S'_{\tau_3} + C_2 S''_{\tau_3} + C_3 S'''_{\tau_3})e^{2\pi i\mathbf{k}\mathbf{r}} + c.c.$$

$$\mathbf{k}=[0.45,0,0]$$

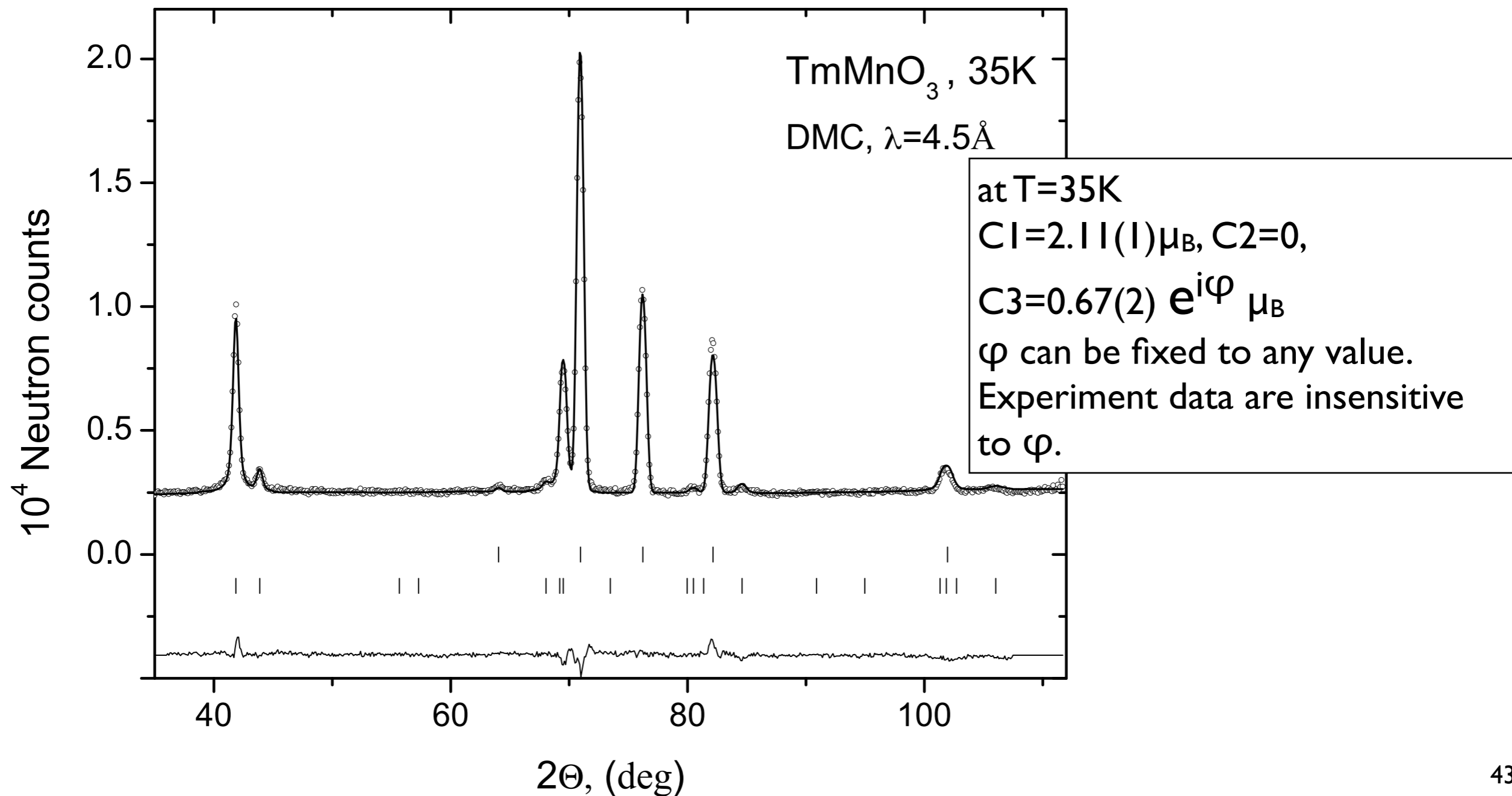
$$S'_{\tau_3} = +1\mathbf{e}_{1x} - a^*\mathbf{e}_{2x} - 1\mathbf{e}_{3x} + a^*\mathbf{e}_{4x}$$

$$S''_{\tau_3} = +1\mathbf{e}_{1y} + a^*\mathbf{e}_{2y} + 1\mathbf{e}_{3y} + a^*\mathbf{e}_{4y}$$

$$S'''_{\tau_3} = +1\mathbf{e}_{1z} + a^*\mathbf{e}_{2z} - 1\mathbf{e}_{3z} - a^*\mathbf{e}_{4z}$$

# Refinement of the data for $\tau_3$

$$\mathbf{S}(\mathbf{r}) = \frac{1}{2}(C_1 S'_{\tau_3} + C_2 S''_{\tau_3} + C_3 S'''_{\tau_3})e^{2\pi i \mathbf{k} \mathbf{r}} + c.c.$$



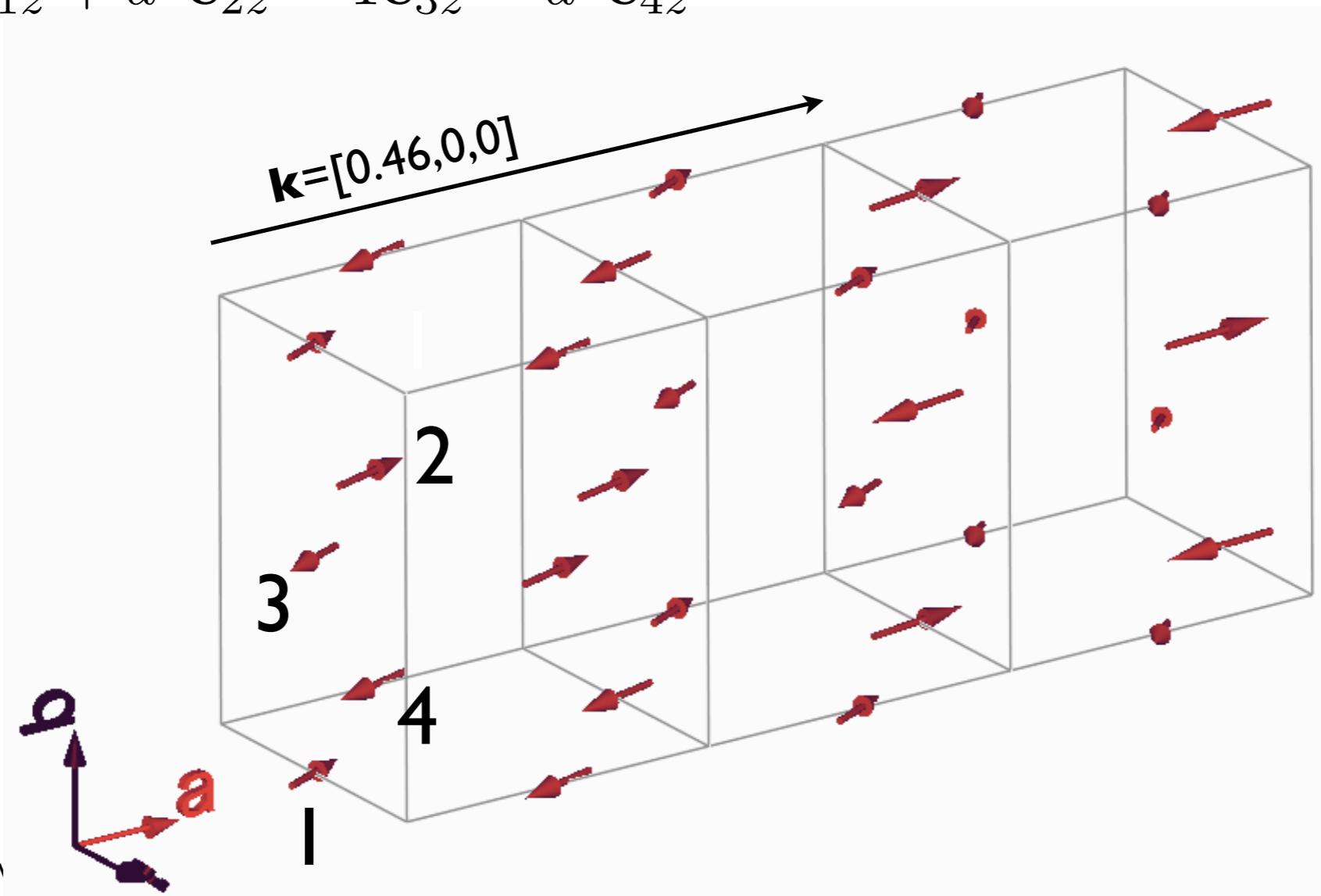
# Visualization of the magnetic structure

a cycloid structure propagating along x-direction

$$\mathbf{S}(\mathbf{r}) = \text{Re} [(C_1 S'_{\tau 3} + |C_3| \exp(i\varphi) S'''_{\tau 3}) \exp(2\pi i \mathbf{k} \mathbf{r})]$$

$$S'_{\tau 3} = +1\mathbf{e}_{1x} - a^* \mathbf{e}_{2x} - 1\mathbf{e}_{3x} + a^* \mathbf{e}_{4x}$$

$$S'''_{\tau 3} = +1\mathbf{e}_{1z} + a^* \mathbf{e}_{2z} - 1\mathbf{e}_{3z} - a^* \mathbf{e}_{4z}$$



# Visualization of the magnetic structure

a cycloid structure propagating along x-direction

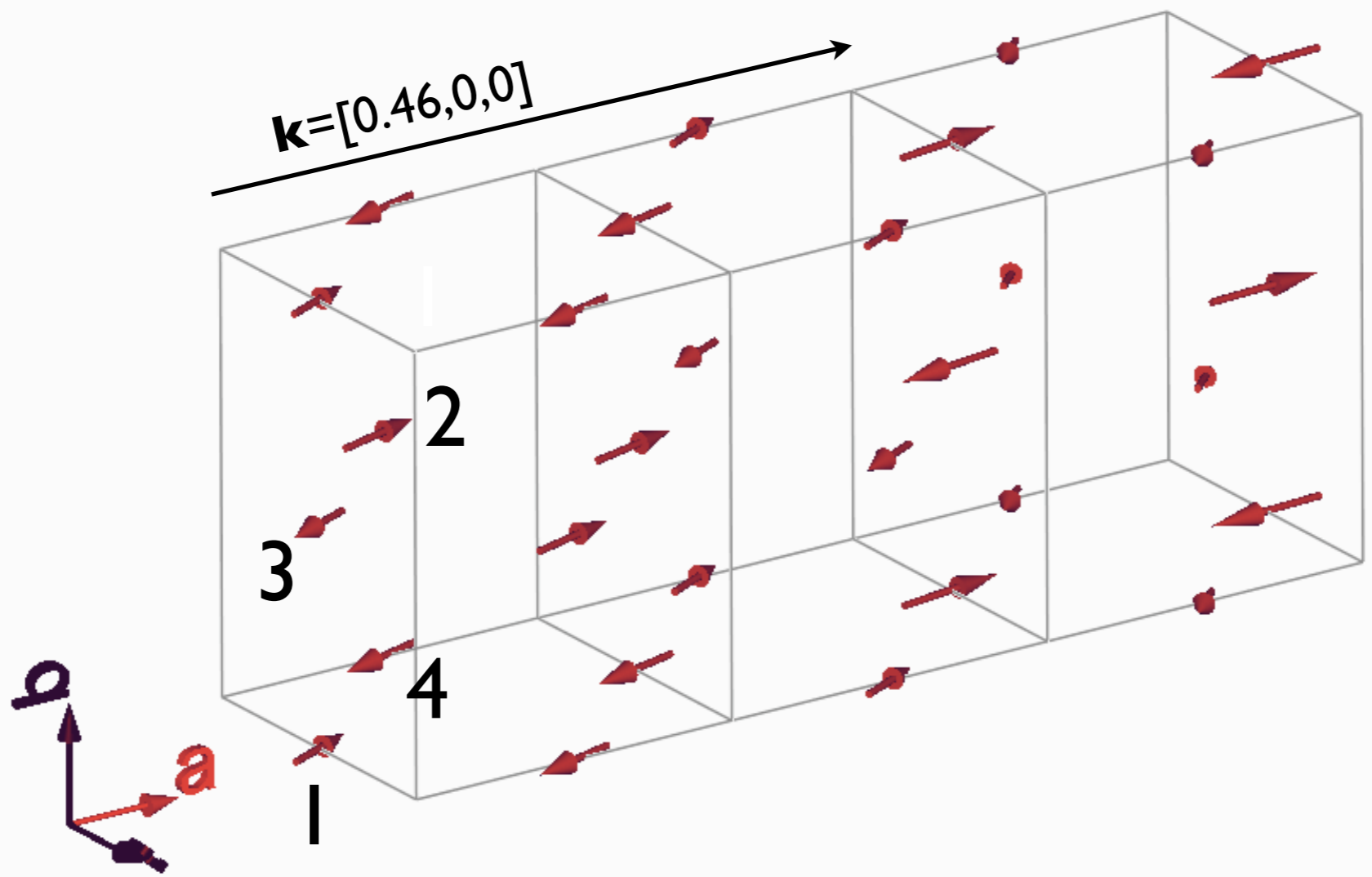
$$\mathbf{S}(\mathbf{r}) = \text{Re} [(C_1 S'_{\tau 3} + |C_3| \exp(i\varphi) S'''_{\tau 3}) \exp(2\pi i \mathbf{k} \cdot \mathbf{r})]$$

$$S'_{\tau 3} = +1\mathbf{e}_{1x} - a^* \mathbf{e}_{2x} - 1\mathbf{e}_{3x} + a^* \mathbf{e}_{4x}$$

$$S'''_{\tau 3} = +1\mathbf{e}_{1z} + a^* \mathbf{e}_{2z} - 1\mathbf{e}_{3z} - a^* \mathbf{e}_{4z}$$

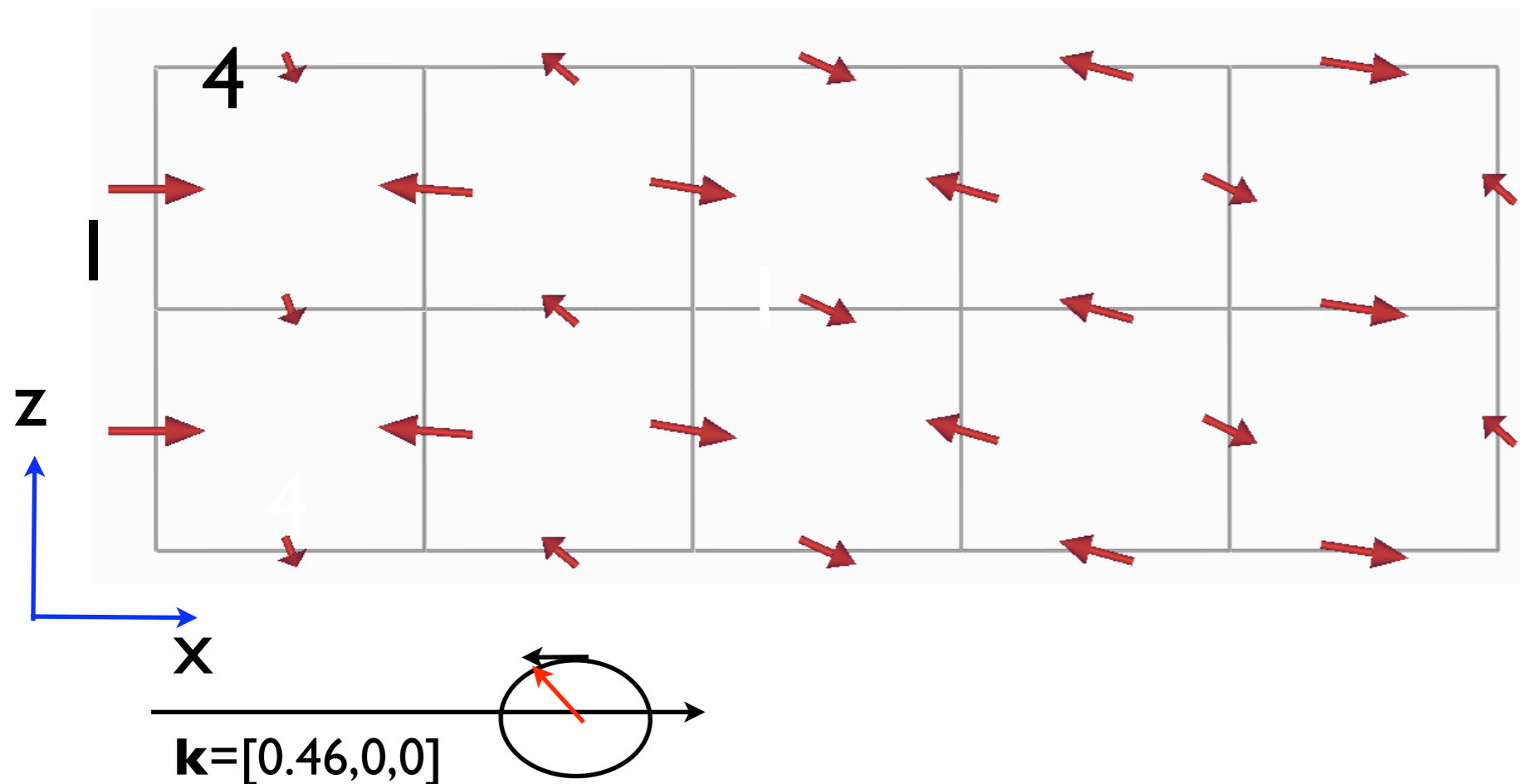
Propagation of the spin, e.g. for atom no. 1

$$\mathbf{S}_1(x) = C_1 \cos(kx) \mathbf{e}_x + |C_3| \cos(kx + \varphi) \mathbf{e}_z$$



# Visualization of the magnetic structure: xz-projection

for arbitrary  $\varphi$ :  
both direction and size of  $S_I$  are changed

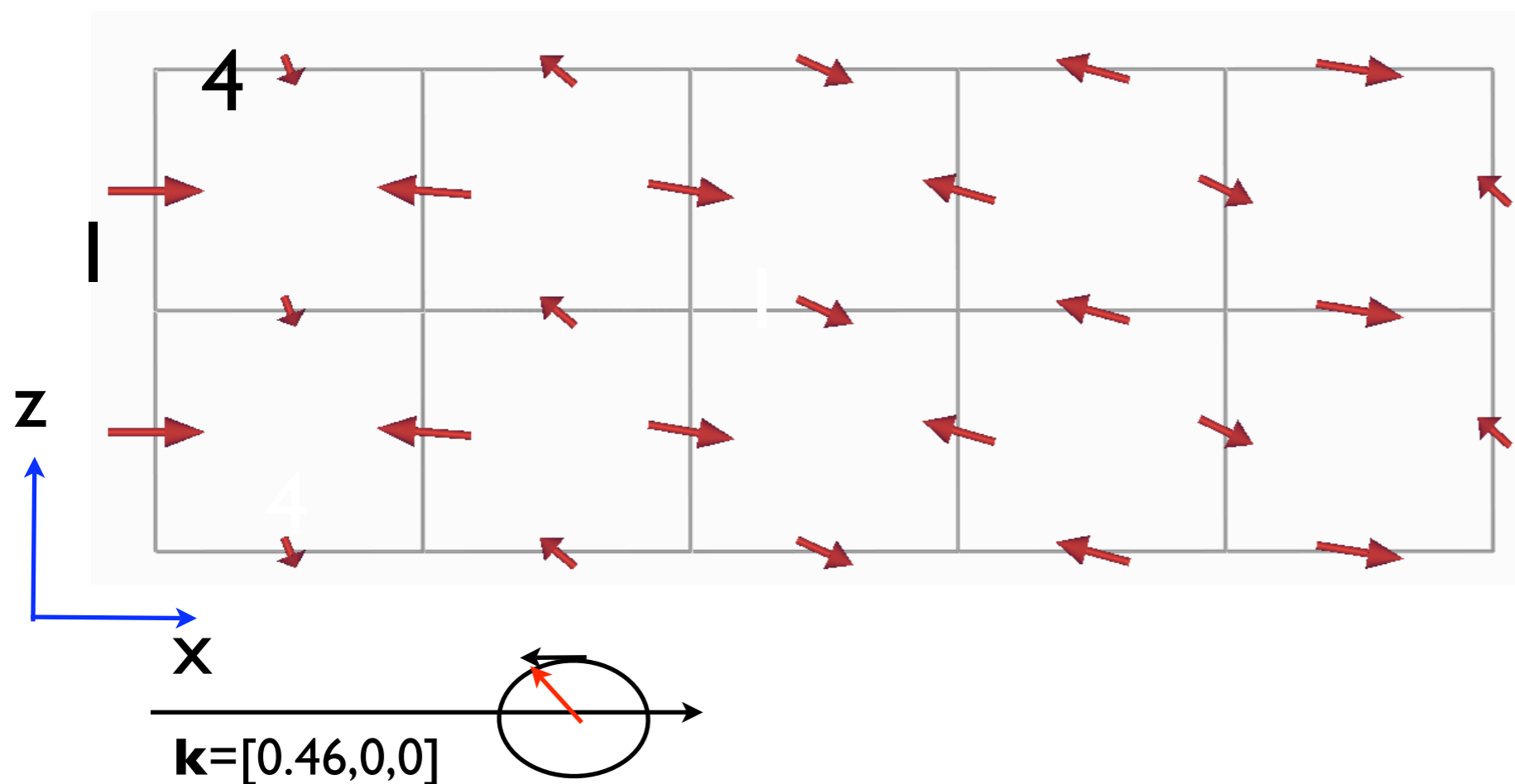


# Visualization of the magnetic structure: xz-projection

for arbitrary  $\varphi$ :  
both direction and size of  $S_I$  are changed

Propagation of the spin, e.g. for atom no. 1  

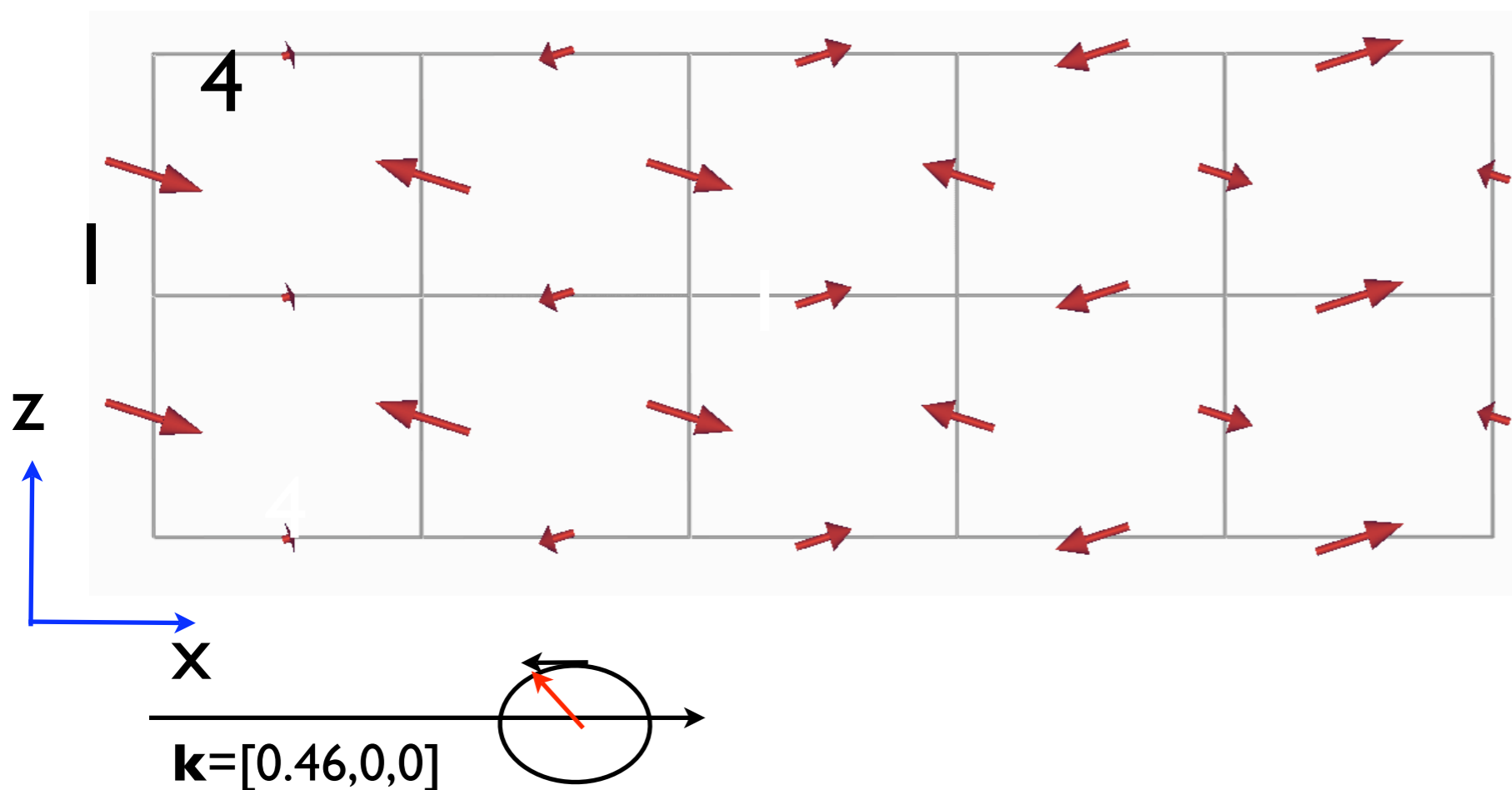
$$\mathbf{S}_1(x) = C_1 \cos(kx) \mathbf{e}_x + |C_3| \cos(kx + \varphi) \mathbf{e}_z$$



# Visualization of the magnetic structure: xz-projection

for  $\varphi=0$ :  
only the size of  $S_l$  are changed

Propagation of the spin, e.g. for atom no. 1  
 $S_1(x) = (C_1 \mathbf{e}_x + |C_3| \mathbf{e}_z) \cos(kx)$



# literature, programs and tutorials/notes

All you need to know about magnetic neutron diffraction.

Symmetry, representation analysis

Yu. A. Izyumov, V. E. Naish and R. P. Ozerov, "*Neutron diffraction of magnetic materials*", New York [etc.]: Consultants Bureau, 1991.

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## COMPUTER PROGRAMS, TUTORIALS

1. Juan Rodríguez-Carvajal (ILL) et al, <http://www.ill.fr/sites/fullprof/>
2. Andrew S. Wills (UCL) [http://www.chem.ucl.ac.uk/people/wills/magnetic\\_structures/magnetic\\_structures.html](http://www.chem.ucl.ac.uk/people/wills/magnetic_structures/magnetic_structures.html)
3. Wieslawa Sikora et al, <http://www.ftj.agh.edu.pl/~sikora/modyopis.htm>

This lecture

<http://sinq.web.psi.ch/sinq/instr/hrpt/praktikum>