

Antiferromagnetic Ordering in the S=1 Quasi-One-Dimensional $\text{PbNi}_{2-x}\text{Mg}_x\text{V}_2\text{O}_8$ (x=0.24)

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1D quantum Heisenberg AF

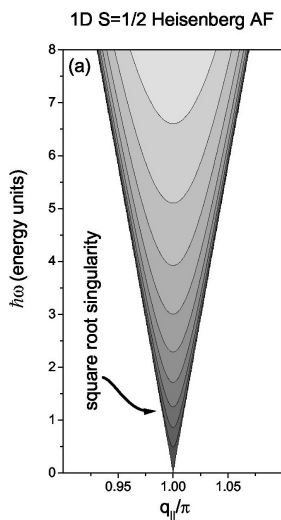
$$H = -J \sum_i \mathbf{S}_i \mathbf{S}_{i+1} + D \sum_i (S_i^z)^2$$

F.D.M.Haldane, *PRL*. **50**, 1153 (1983)

I. Affleck et al, *PRL* **59**, 799 (1987)

▪ $S=1/2$ (or half-odd-integer)

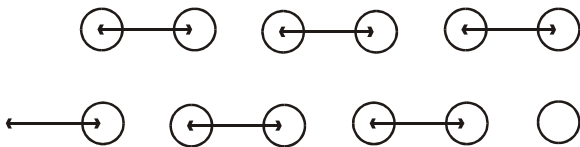
Continuum excitation spectrum $E(k)$



e.g. good 1D:
 Sr_2CuO_3 , Sr_2CuO_3

▪ Arbitrary weak
interchain J_1 suppress QF
 $T_N \sim J_1$, $\mu \sim J_1/J$
e.g. SrCuF_3 $T_N \approx 40\text{K}$,
 $\mu \approx 0.5 \mu_B$

The spin- $1/2$ dimerized Valence Bond states



▪ Ground state is **Resonating Valence Bond /RVB**:

QF restore translational symmetry
and mix in bounds of greater length

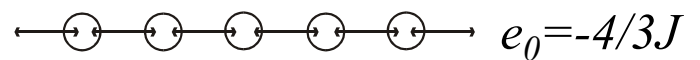
▪ $S=1$ (or integer)

$$E(k) = [v^2(k-\pi)^2 + \Delta^2]^{1/2}$$

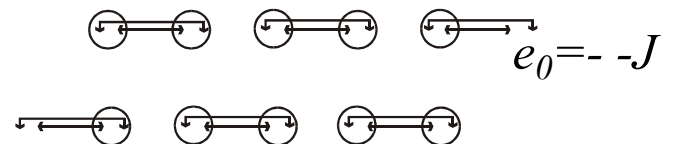
Haldane gap $\Delta \sim 2J e^{-\pi S}$

Singlet ground state

▪ The spin-1 **Valence Bond Solid /VBS**



The spin-1 dimerized states



ESR in $\text{Cu}^{2+}/\text{Ni}^{2+}$ substituted

$[\text{Ni}(\text{C}_2\text{H}_8\text{N}_2)_2(\text{NO}_2)]\text{ClO}_4$ (NENP)
shows spin- $1/2$ end states $\rightarrow$ **VBS**

▪ Energetics of a Haldane gap AF

from MC calculations

$$\Delta_{xx} \approx \Delta - (2/3)D; \Delta_{zz} \approx \Delta + (4/3)D$$

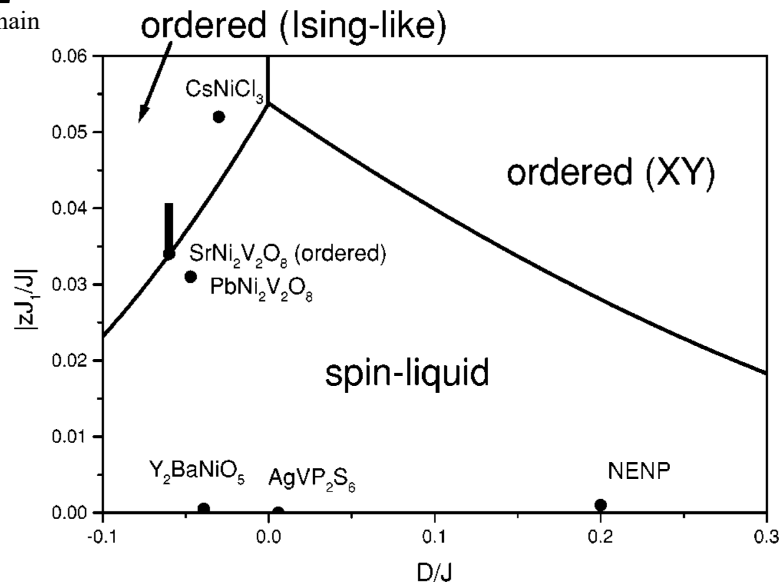
$$\Delta = 0.41J; v = 2.49J; \xi = 2J/\Delta \sim 5$$

$$E_0 = e_0 + A \exp(-L/\xi); e_0 = -1.4J$$

Long range order in a Haldane-gap $S=1$ quasi 1D-AFM

$$H = -J \sum_i \mathbf{S}_i \mathbf{S}_{i+1} + D \sum_i (S_i^z)^2 - J_1 \sum_{\text{interchain}} \mathbf{S}_i \mathbf{S}_{i+1}$$

- Interchain interactions J' and ion anisotropy D promote LRO /Phase diagram from T.Sakai, M.Takahashi PRB **42**, 4537 (1990); A.Zheludev et al., PRB **62**, 8921 (2000)./



- External magnetic field suppress QF, restore gapless spectrum. $\text{Ni}(\text{C}_5\text{D}_{14}\text{N}_2)_2\text{N}_3(\text{PF}_6)$ /NDMAP/ (Pnmn, chains along c) LRO AFM at $>5\text{T}$ /Y.Chen et al., Phys. Rev. Lett. **86**, 1618 (2001)

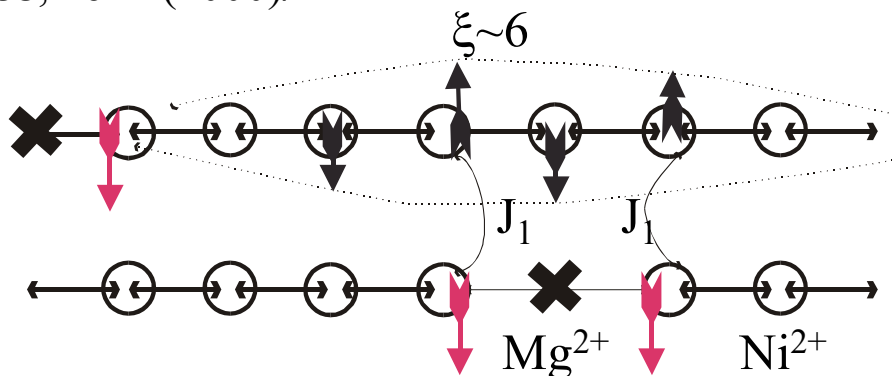
- Spin-vacancy induced LRO** /Theoretically suggested: E.F. Shender and S.A. Kivelson, Phys. Rev. Lett. **66**, 2384 (1991) /

Experimentally observed only in $\text{PbMg}_{2-x}\text{A}_x\text{V}_2\text{O}_8$

/Y. Uchiyama, et al, PRL **83**, 632 (1999); Physica B**284-288**, 1641 (2000)/

Energetics of $\text{PbMg}_2\text{V}_2\text{O}_8$

$\Delta_{zz}=35\text{K}$, $\Delta_{xx}=46\text{K}$, $J=-95\text{K}$, $J_1=-2\text{K}$

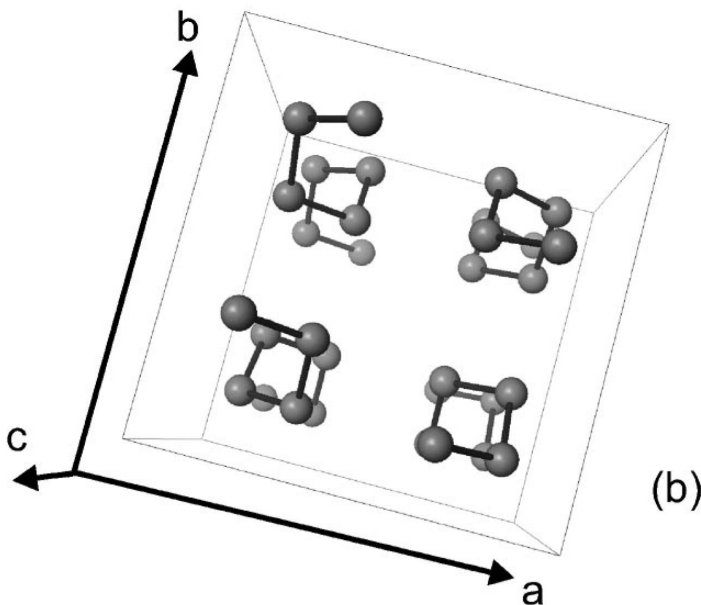
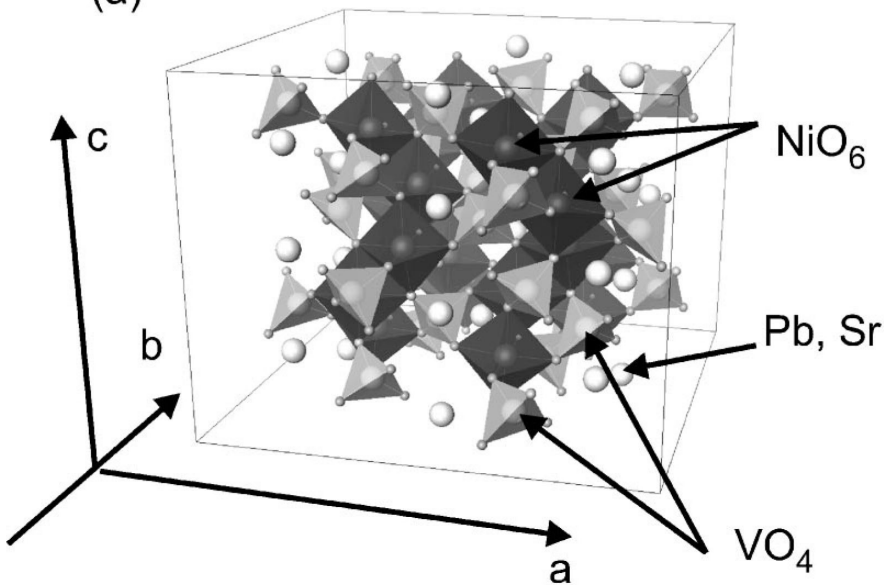


AFM in $\text{PbNi}_{2-x}\text{Mg}_x\text{V}_2\text{O}_8$ ($x=0.24$)

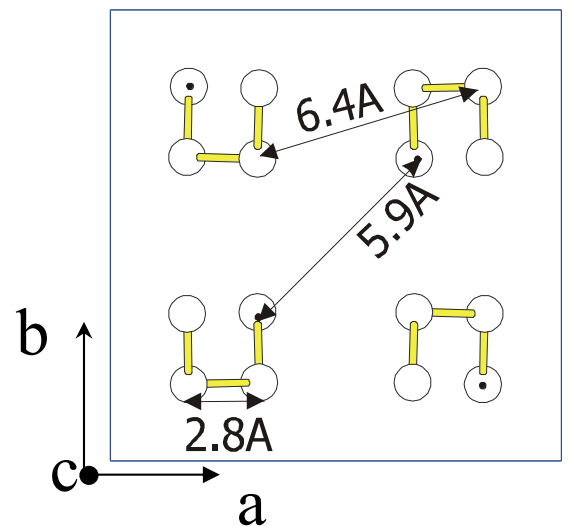
Crystal structure of $\text{PbNi}_2\text{V}_2\text{O}_8$

a tetragonal crystal structure ($a=b=12.2 \text{ \AA}$, $c=8.4 \text{ \AA}$, space group $I4_1cd$) isomorphous to $\text{SrNi}_2\text{V}_2\text{O}_8$

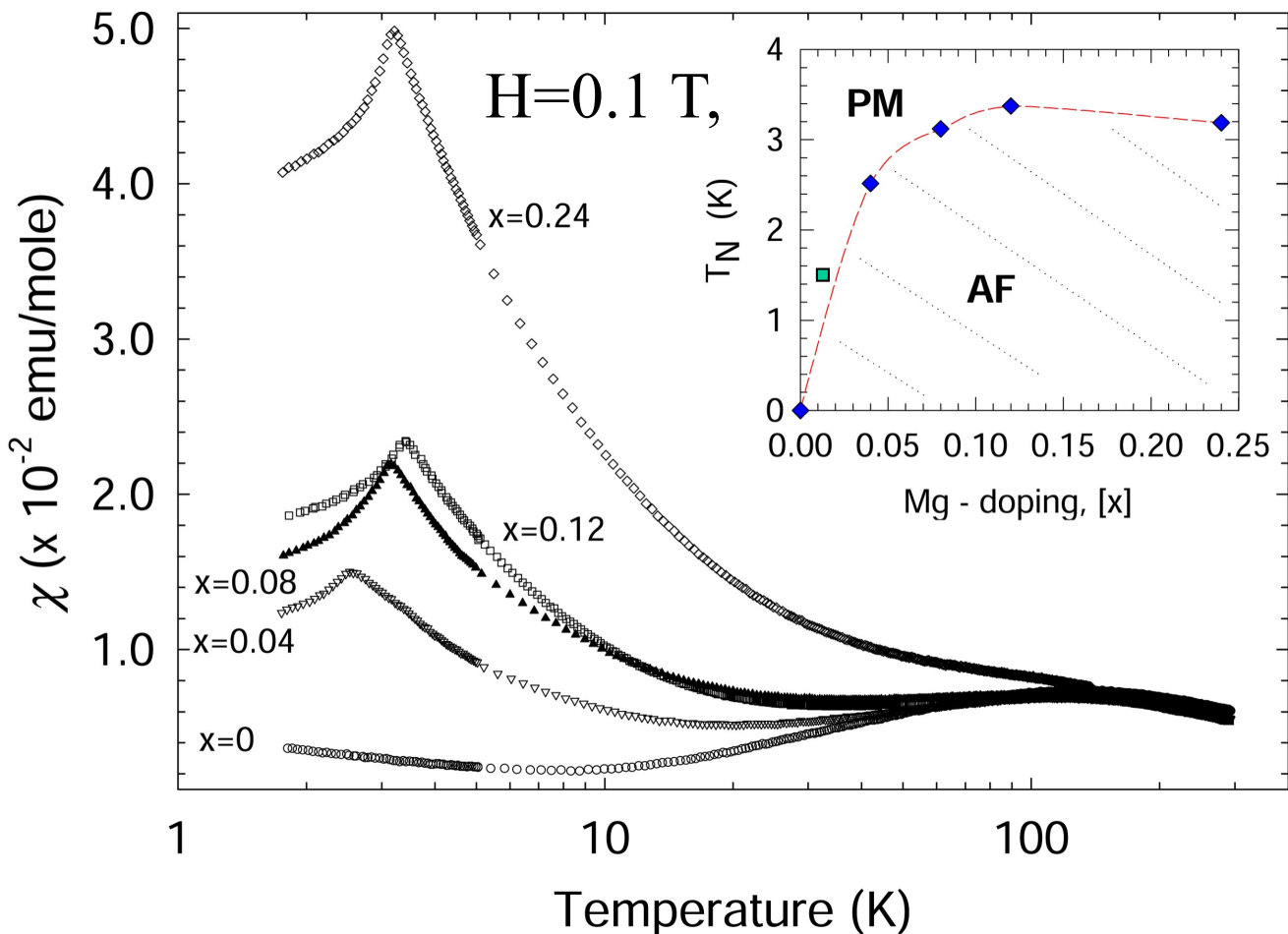
(a)



Ni-Ni distances



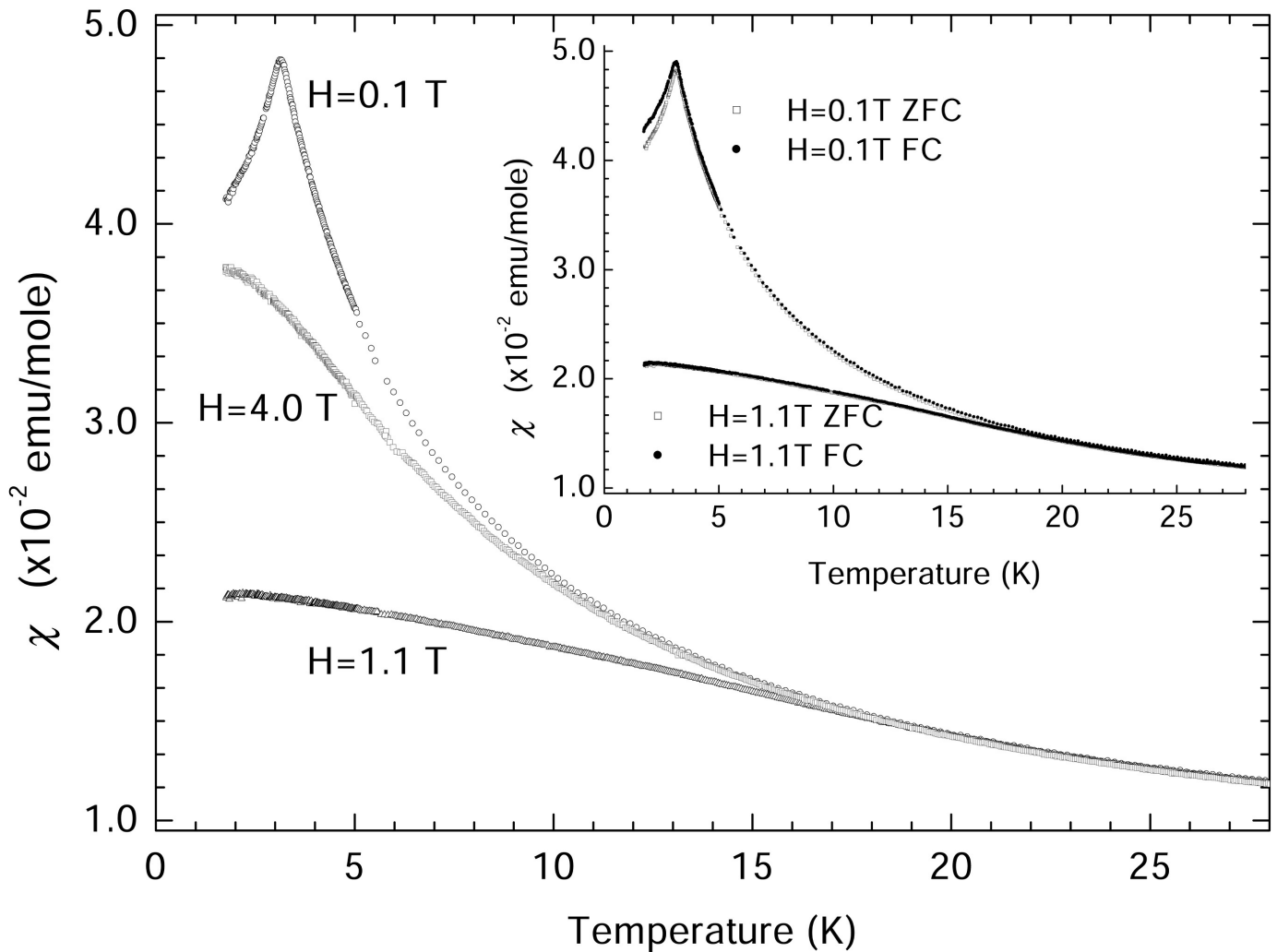
$\chi_{dc}(T)$ and T - x phase diagram of $\text{PbNi}_{2-x}\text{Mg}_x\text{V}_2\text{O}_8$



- Low- T ($6 < T < 40\text{K}$):
 $\chi(T) \cong \chi(0) + C e^{-\Delta/kT}$, $\Delta \cong \text{HaldaneGap} = 30.2(2) \text{ K}^*$
- High- T ($160 < T < 300 \text{ K}$):
 $\chi_{dc}(T) = N(g\mu)^2/kT f(J/T)$ “a series expansion for a $S=1$ 1D-HAF”
 $\rightarrow J = -90.4(3) \text{ K}^*$. $\Delta/|J| = 0.31$

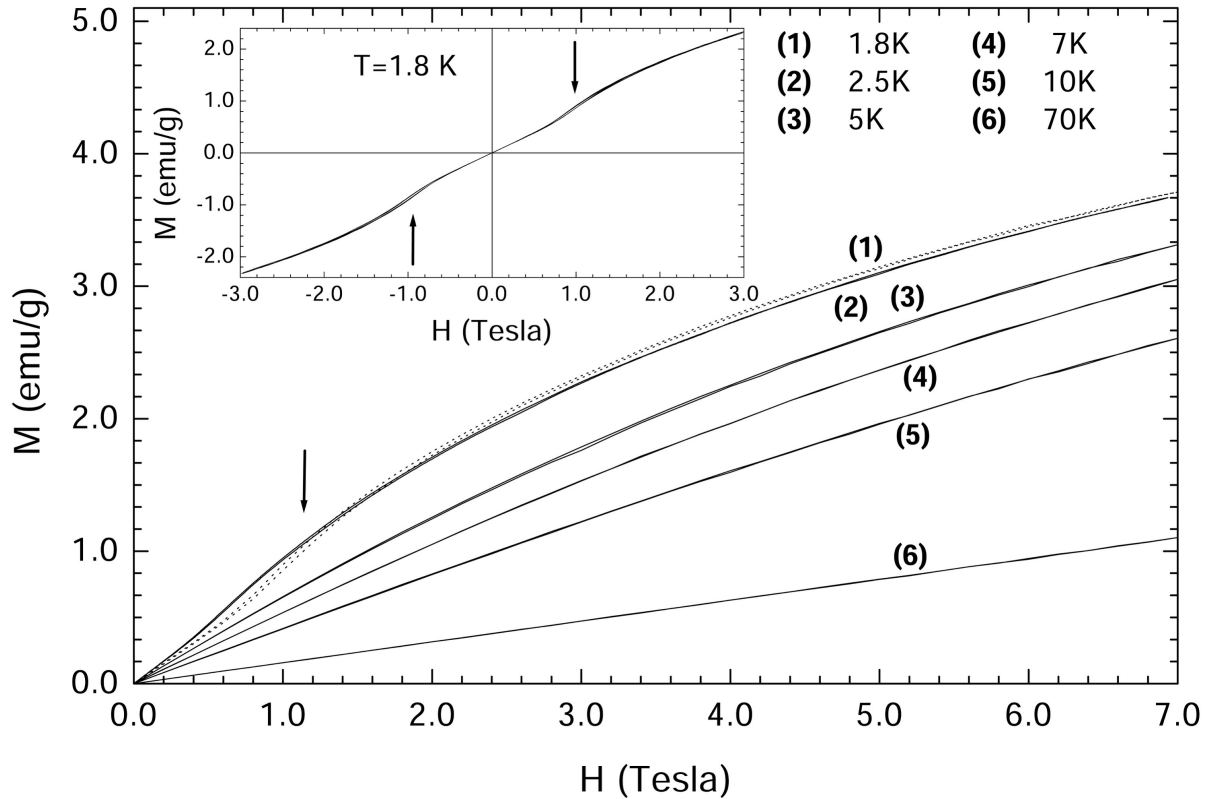
*In accord with Y. Uchiyama, et al, *Phys. Rev. Lett.* **83**, 632 (1999)

Suppressing of the AF transition by external magnetic field



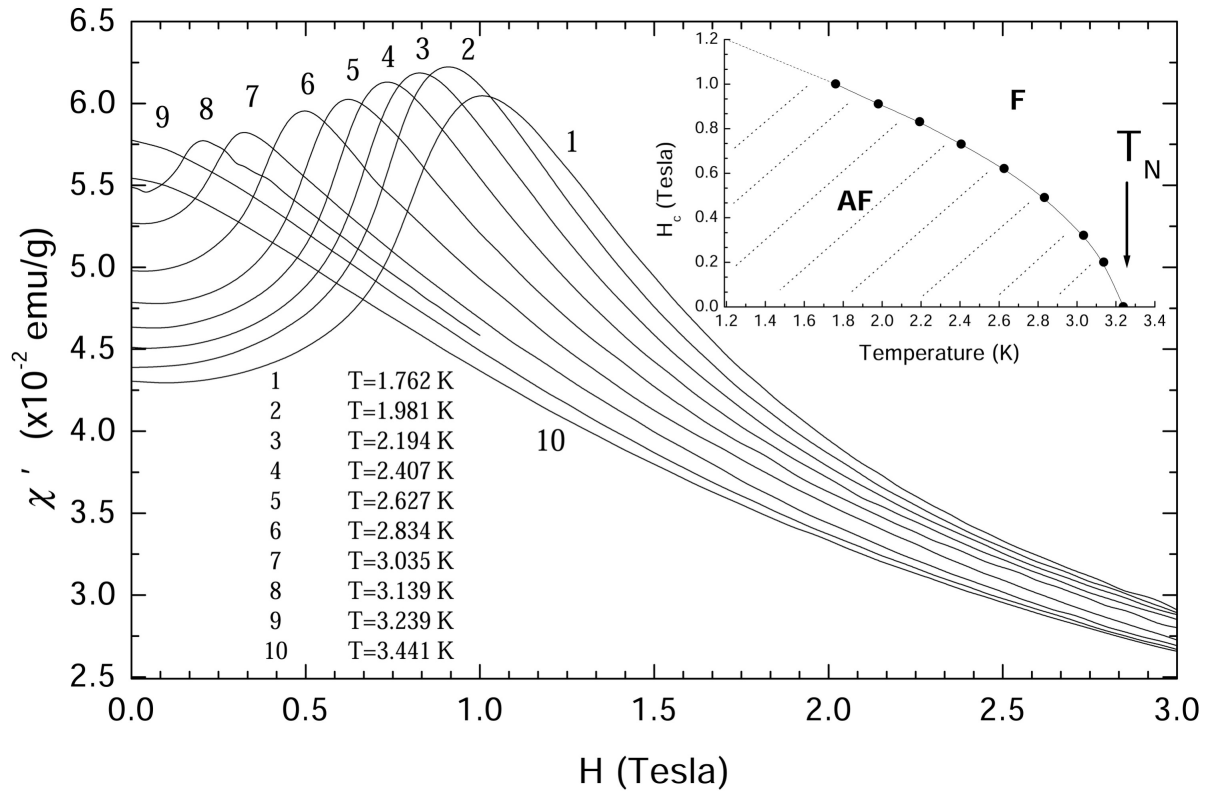
Temperature dependence of the dc susceptibility of $\text{PbNi}_{1.76}\text{Mg}_{0.24}\text{V}_2\text{O}_8$ for ZFC conditions in different magnetic fields. *Inset:* χ_{dc} vs. T , measured under ZFC and FC conditions in 0.1 and 1.1 Tesla.

Metamagnetic behaviour of $\text{PbNi}_{2-x}\text{Mg}_x\text{V}_2\text{O}_8$ ($x=0.24$)



Magnetisation, M as a function of the applied magnetic field, H for various temperatures above and below $T_N=3.2\text{K}$. *Inset:* Full hysteresis loop for $T=1.8\text{K}$.

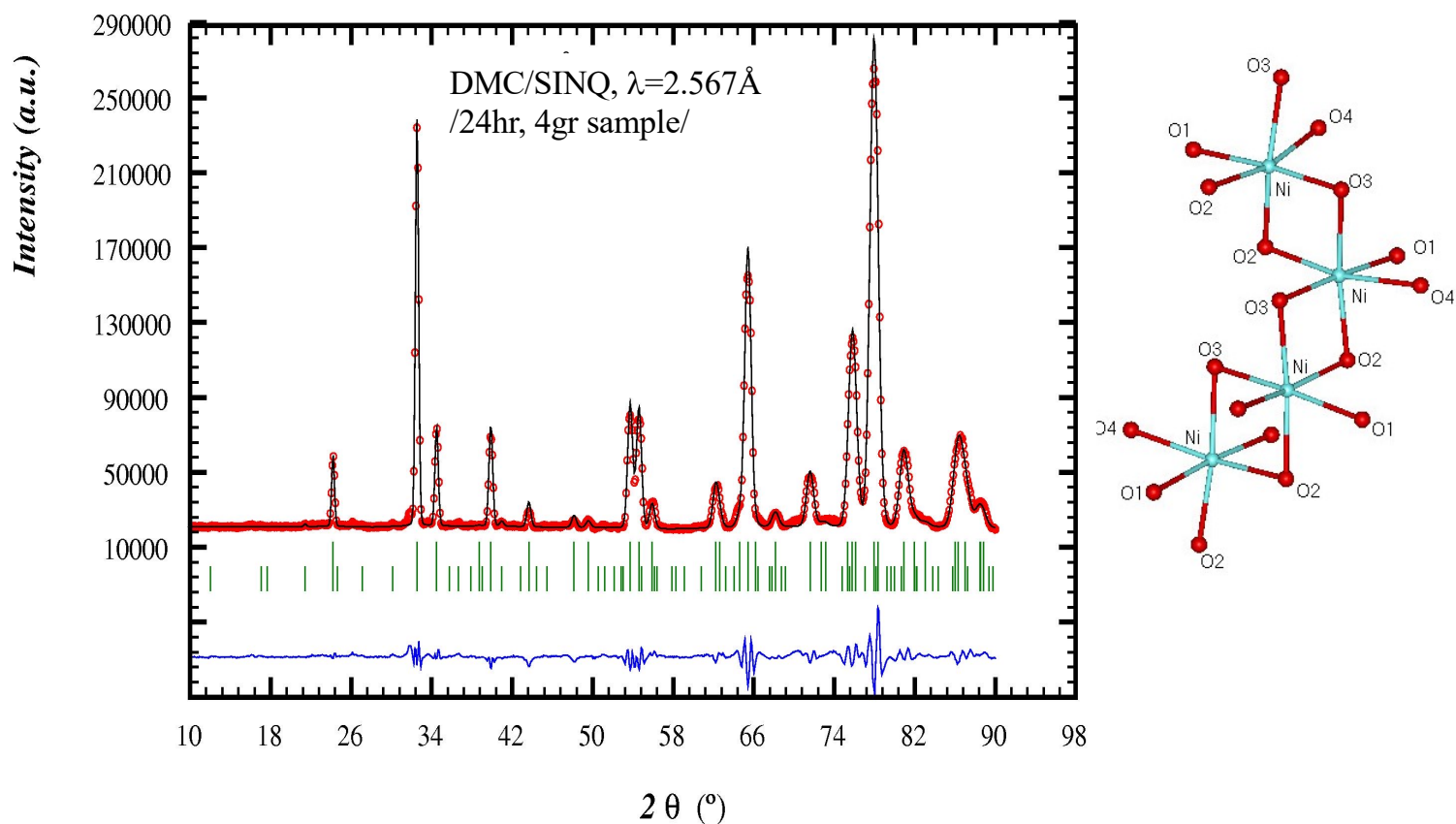
H - T phase diagram of $\text{PbNi}_{2-x}\text{Mg}_x\text{V}_2\text{O}_8$ ($x=0.24$)



Dc magnetic field (H) dependence of the real part of the ac susceptibility, $\chi'(H)$, measured ($h_{ac}=1$ Oe, $f=1$ kHz) at several temperatures below T_N . *Inset*: Phase diagram in the H - T plane.

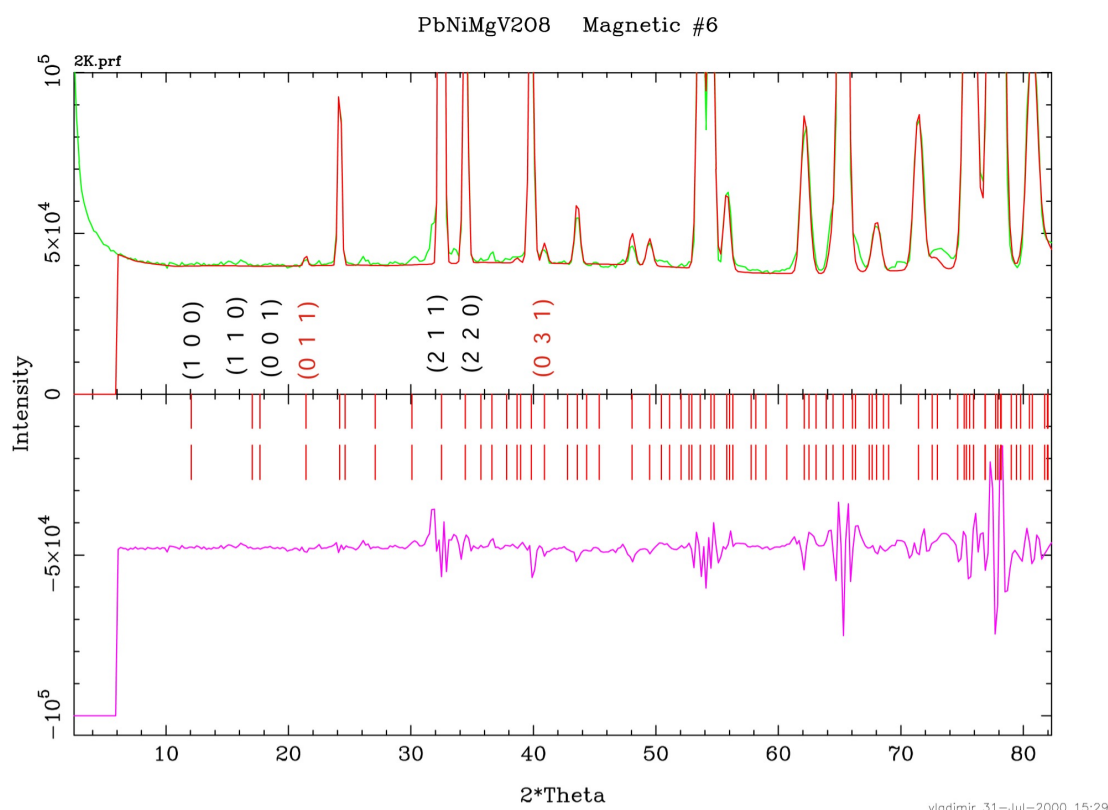
Neutron diffraction pattern.

Crystal structure.

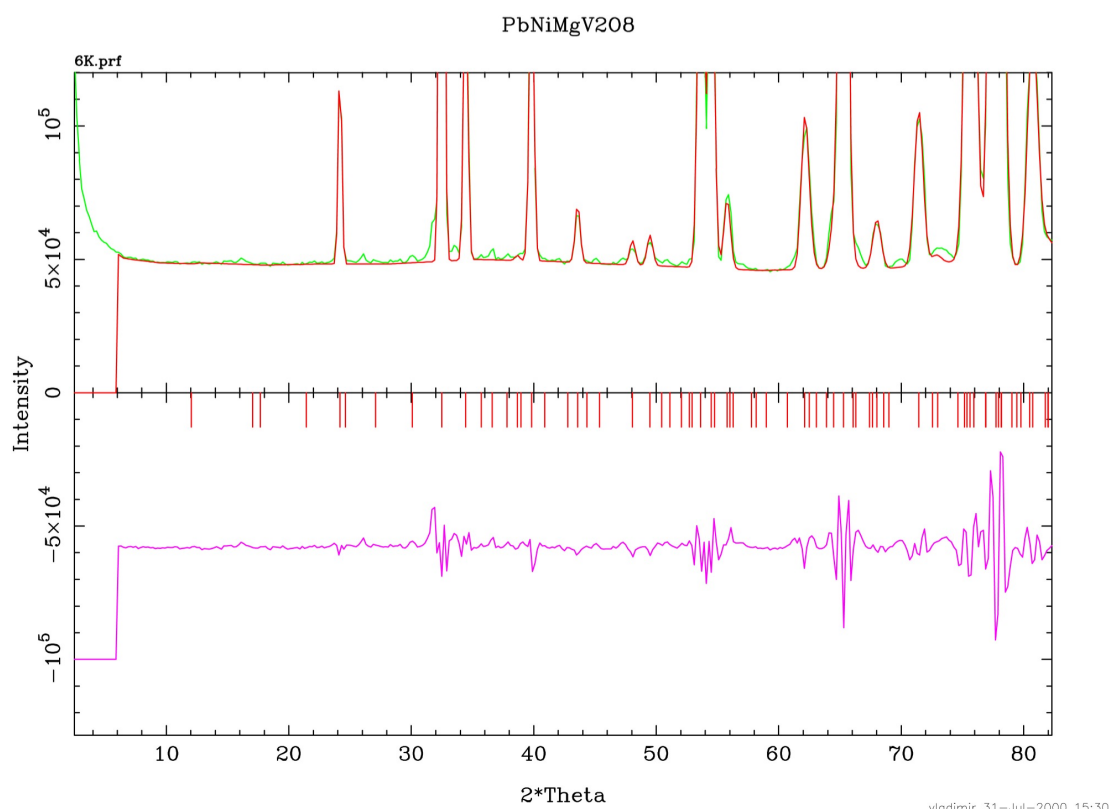


	site	x	y	z
Pb	8a	0	0	0
Ni	16b	0.3262(9)	0.3303(7)	0.1714(14)
Mg	16b	0.3262(9)	0.3303(7)	0.1714(14)
V	16b	0.2578	0.0797	0.1011
O(1)	16b	0.155(1)	0.495(2)	-0.097(2)
O(2)	16b	0.338(2)	0.655(1)	0.434(2)
O(3)	16b	0.161(2)	0.689(1)	0.671(2)
O(4)	16b	0.330(1)	0.499(1)	0.199(2)
Space group: $I4_1cd$. $a=b= 12.2448(7)\text{\AA}$, $c= 8.3592(7)\text{\AA}$				

PbNi_{2-x}Mg_xV₂O₈ (x=0.24), DMC $\lambda=2.567\text{\AA}$



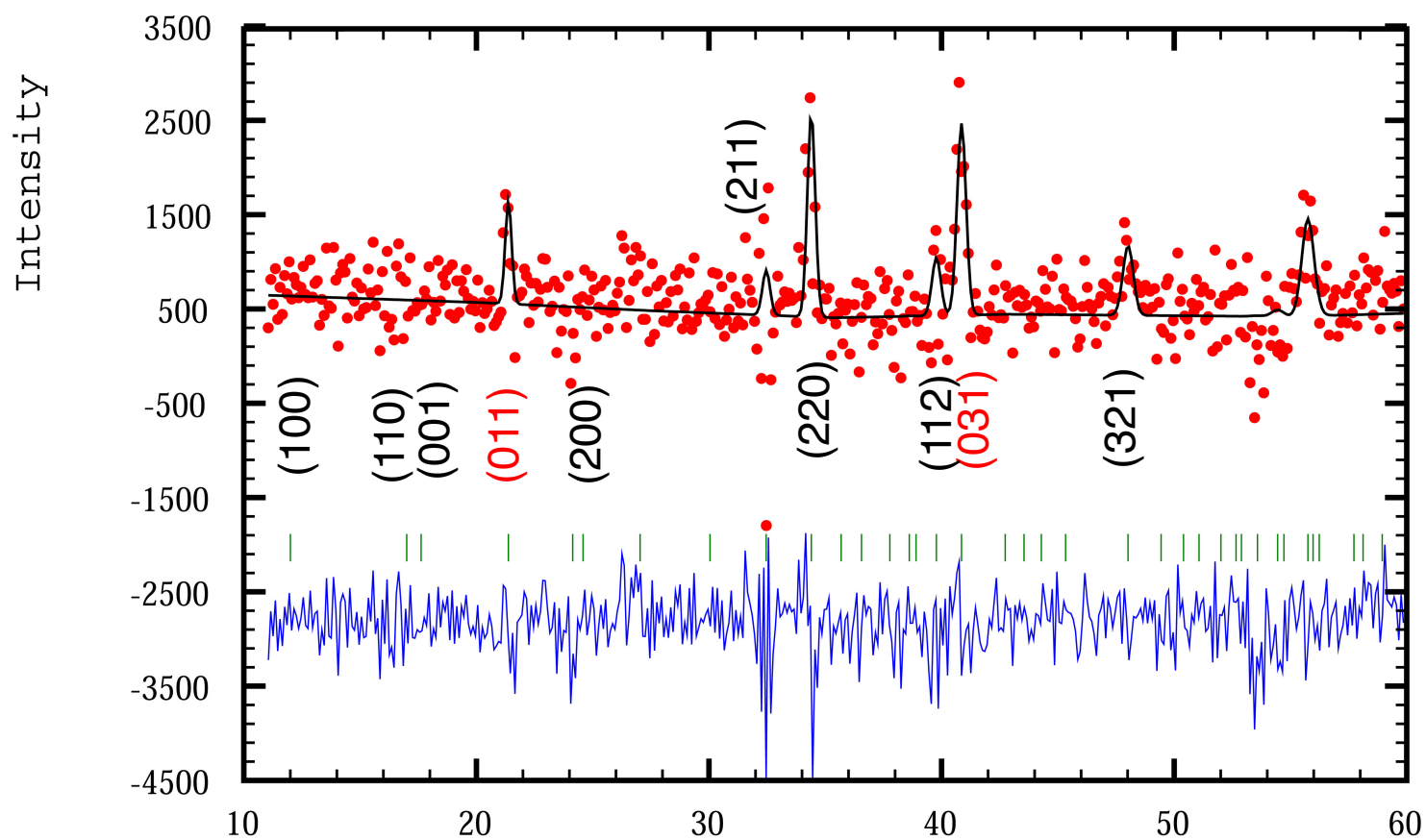
T=2K



T=4K

AFM in PbNi_{2-x}Mg_xV₂O₈ (x=0.24)

Difference pattern “2K-4K”



$k=0$

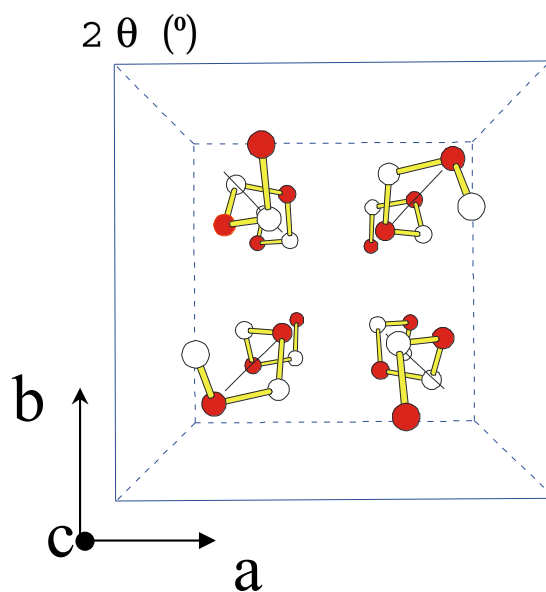
$h+k+l=2n$ present

$h+k+l=2n+1$ extinct

$\mu_{\text{Ni}} \parallel \mathbf{c}$. AFM spirals

$\chi^2=1.55$

Bad $\chi^2=2.7$

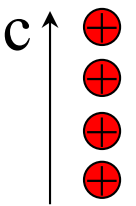
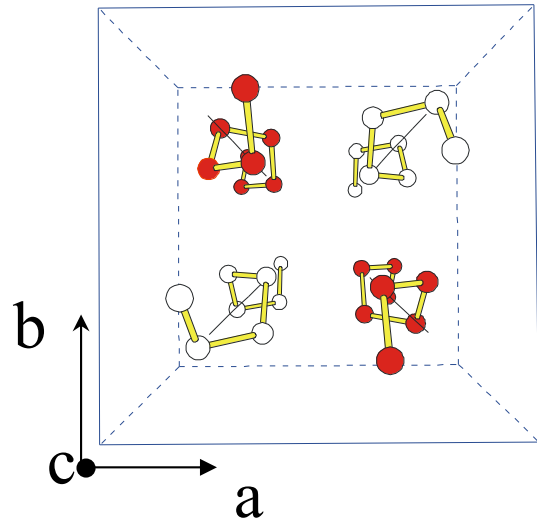
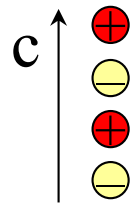
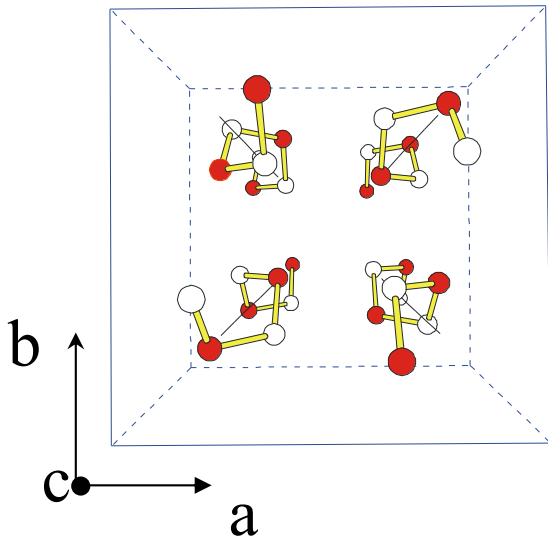


Spin Configurations

$I4_1cd$ sp.gr., 16 Ni^{2+}

✓ I. AF-chains (spirals)

2. F-chains (spirals)

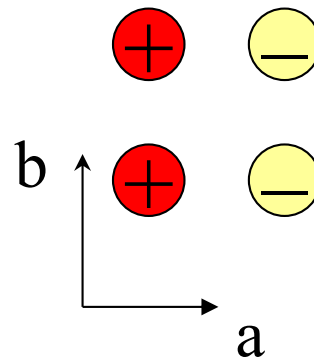
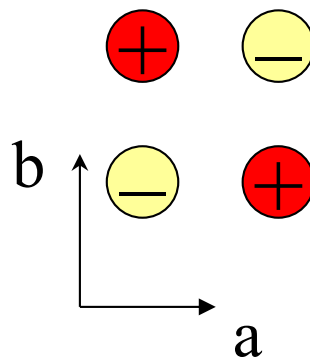
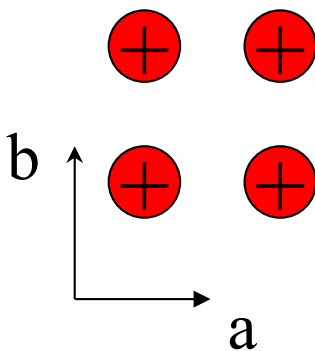


Arrangements of spirals

(1)

(2) ✓

(3)



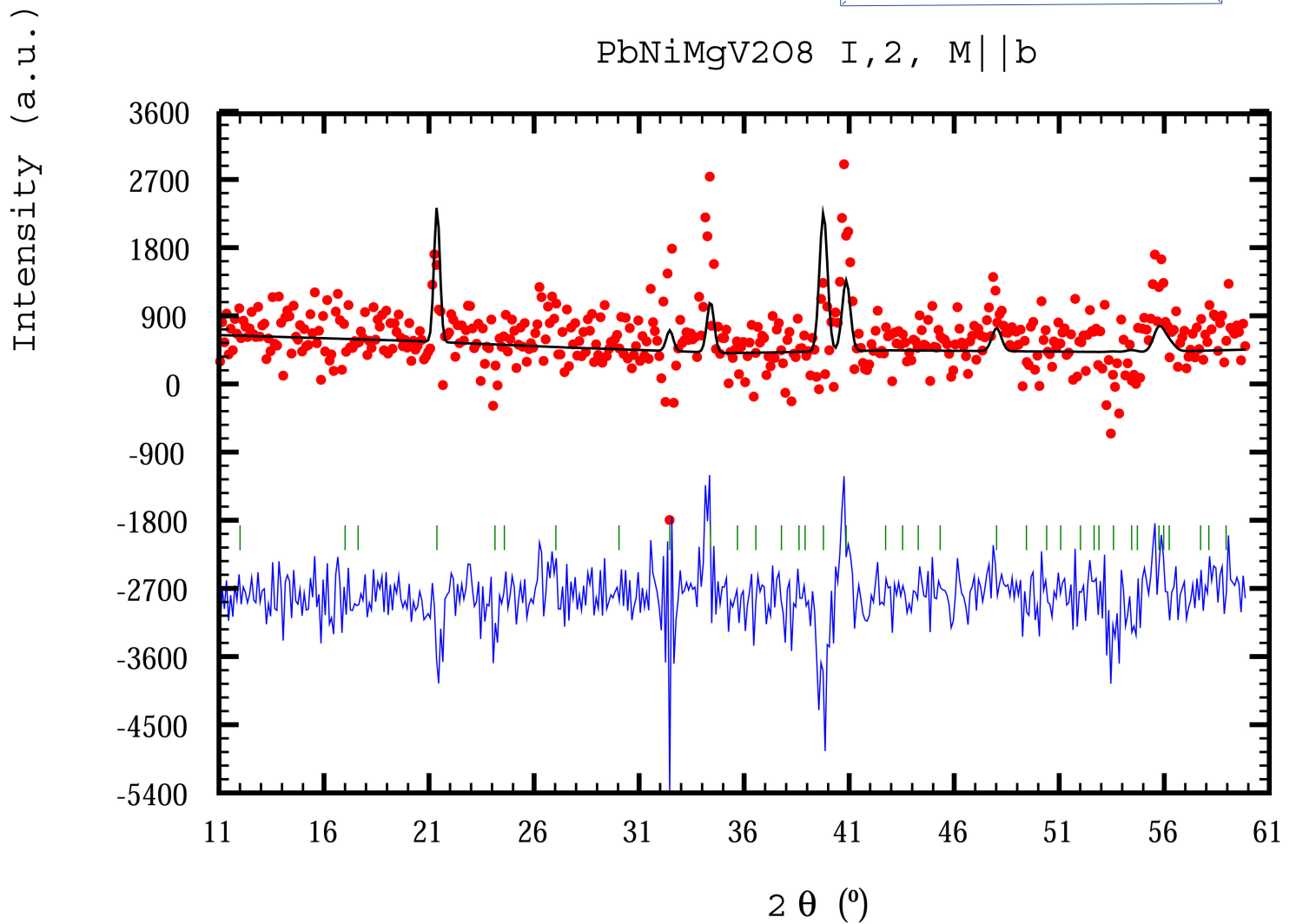
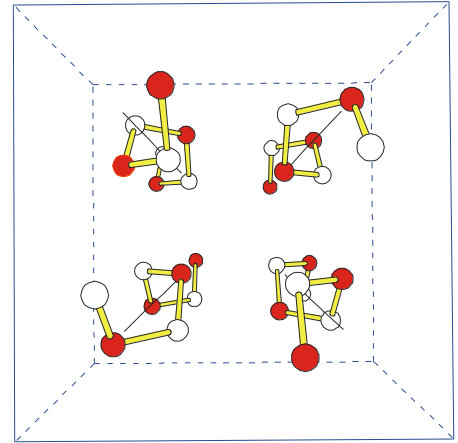
III. Helix structure (90° spin rotation along the spiral)

Spin direction in AFM chains

Illustration of sensitivity to the magnetic moment direction.

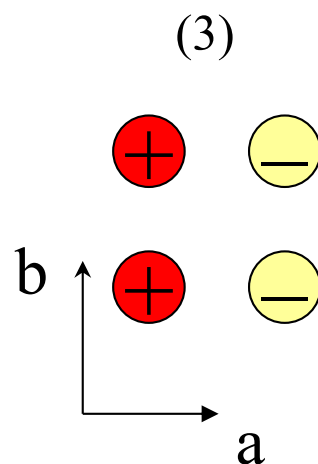
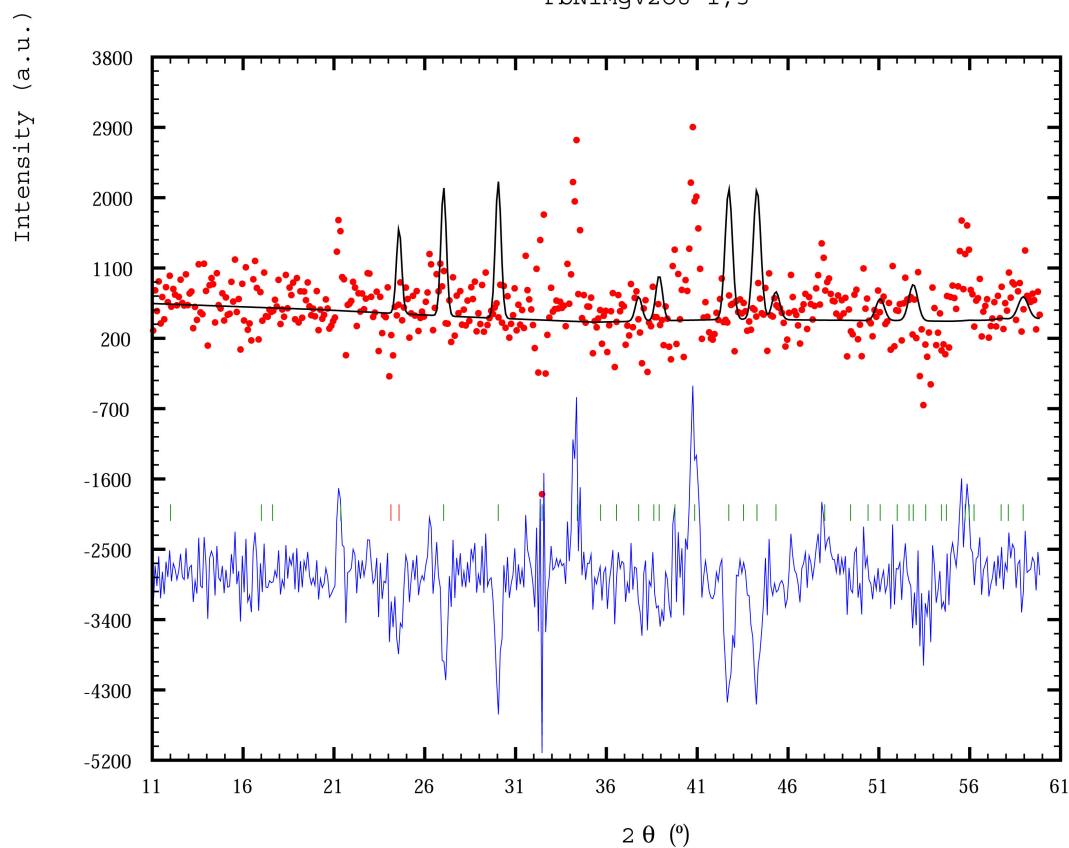
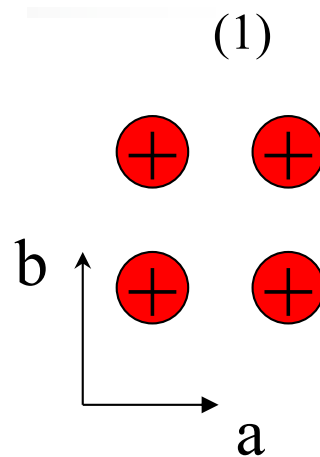
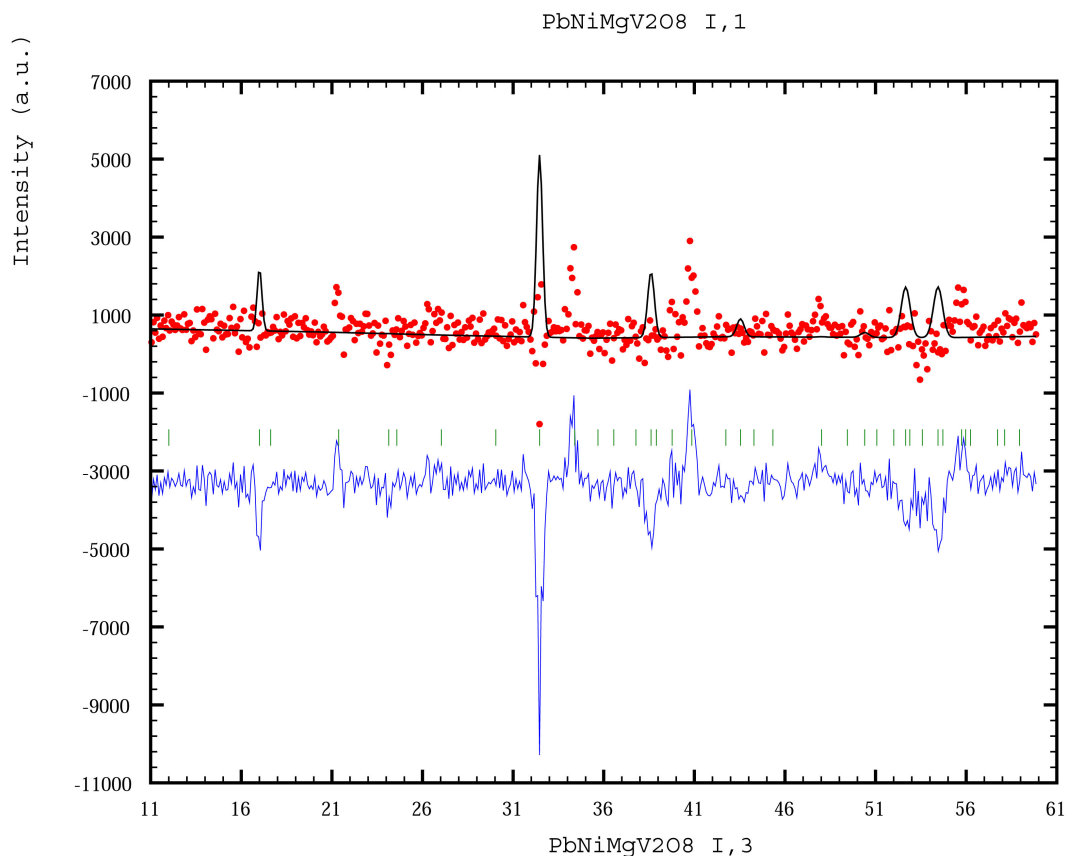
AF-chains, $\mu_{\text{Ni}} \parallel (\mathbf{ab})$ type 2.

Wrong direction $\rightarrow$ bad $\chi^2=2.05$.

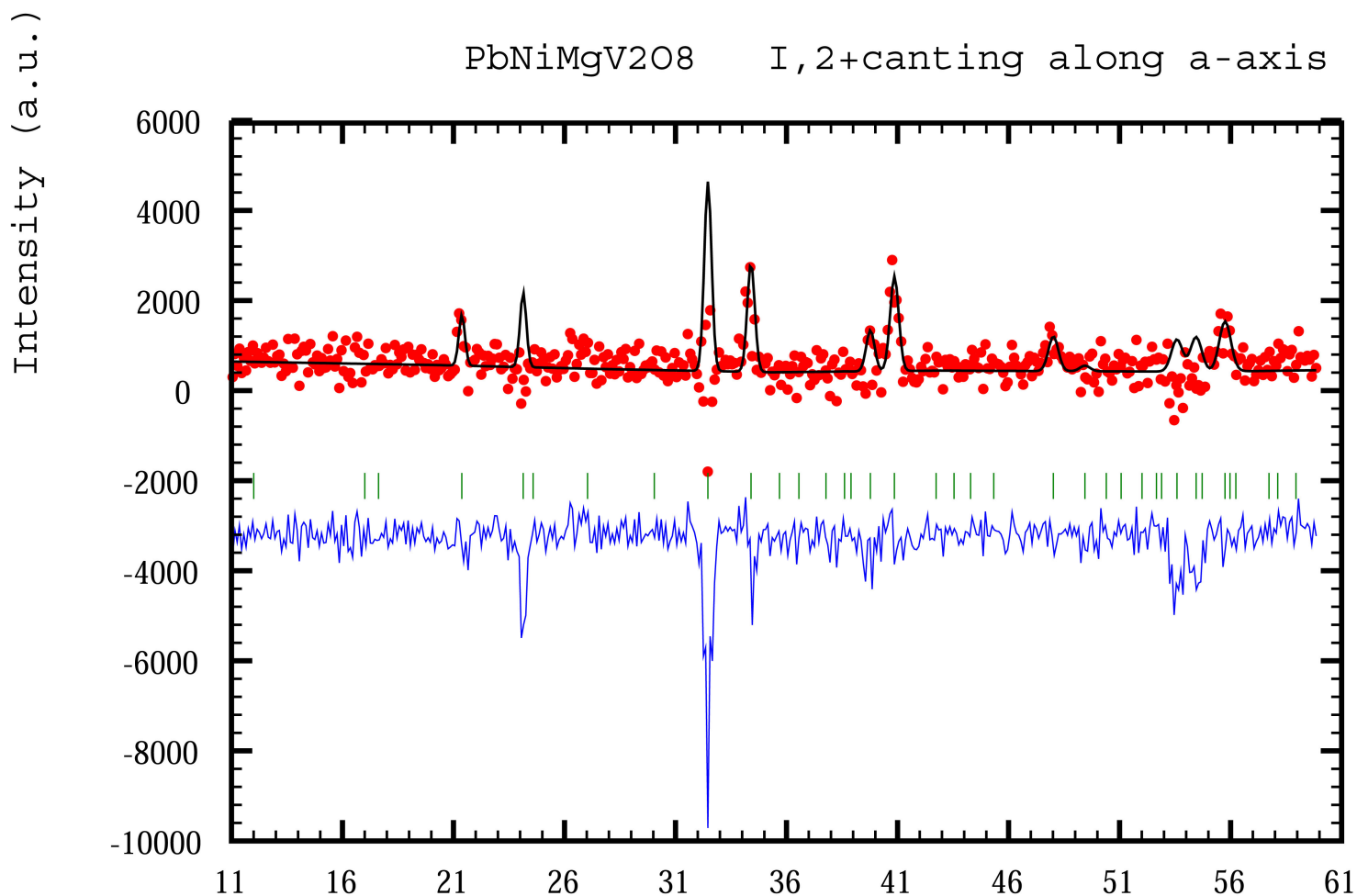


Examples of “wrong” AF-spiral configurations

Theoretical profiles calculated for the same moment (scale factor) obtained from refinement to the “right” spin arrangement.

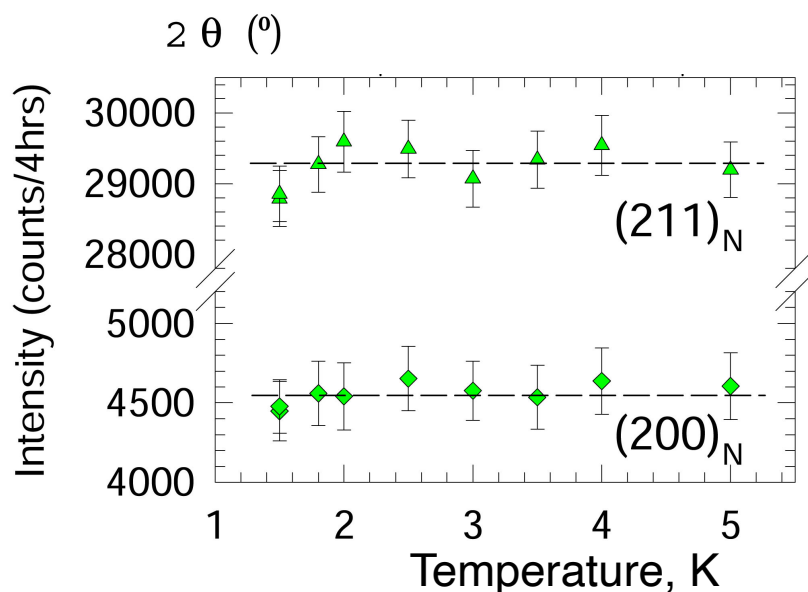


FM-canting of spins

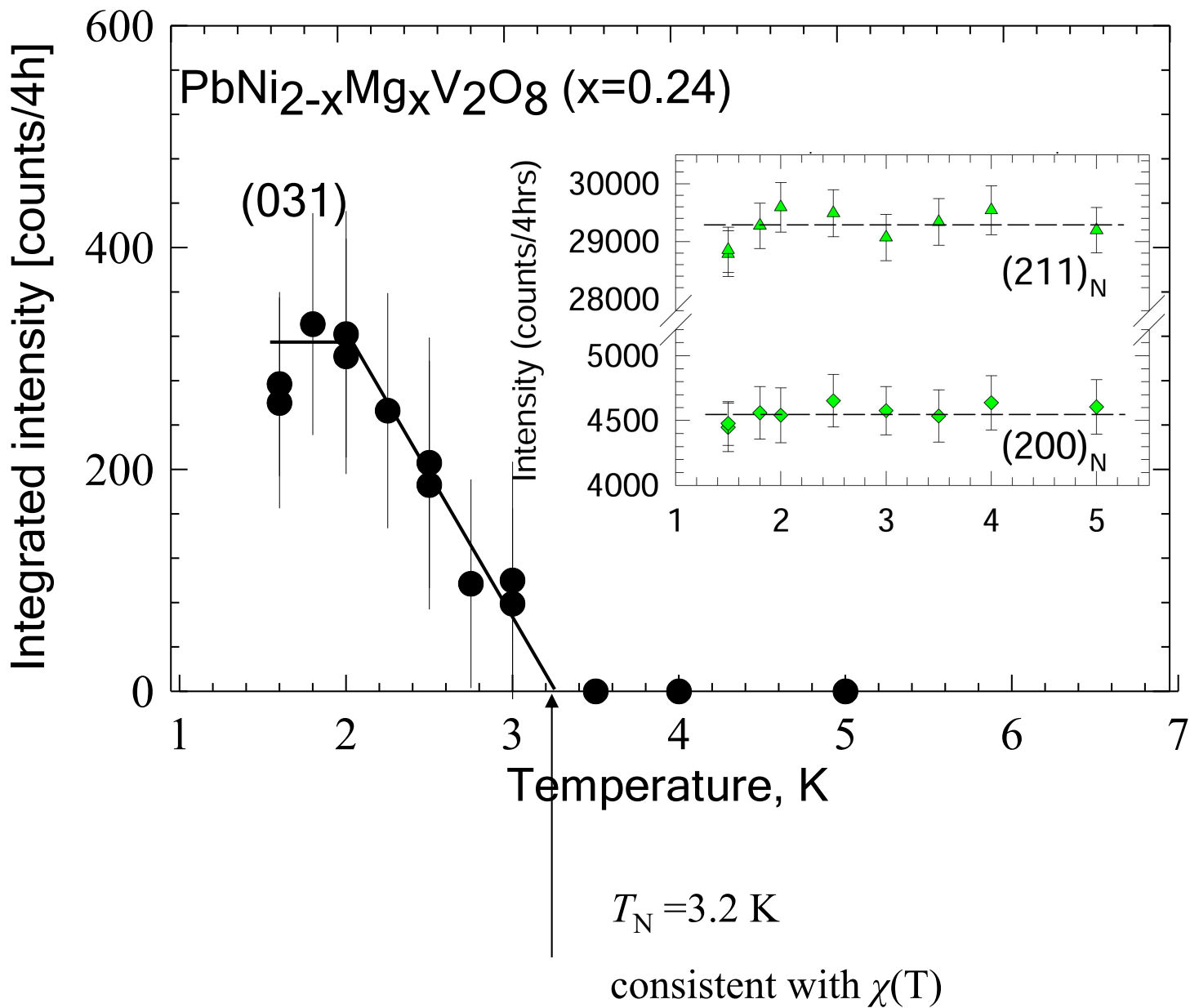


Pattern calculated assuming a FM component along the [100] direction and a moment size equal to half of that found for the purely AF spin AF-chains

statistical estimate of the canting component $\mu_x \sim (0.03 \pm 0.19) \mu_z$.



T-dependence of the integrated Bragg peak intensities

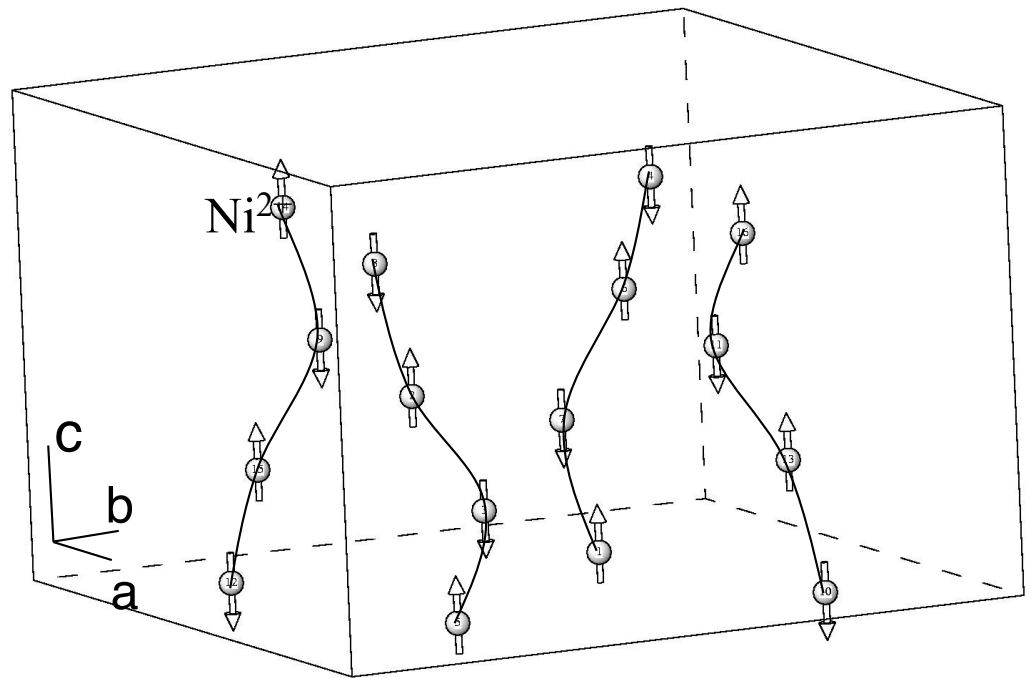


Magnetic Structure of $\text{PbNi}_{2-x}\text{Mg}_x\text{V}_2\text{O}_8$ ($x=0.24$)

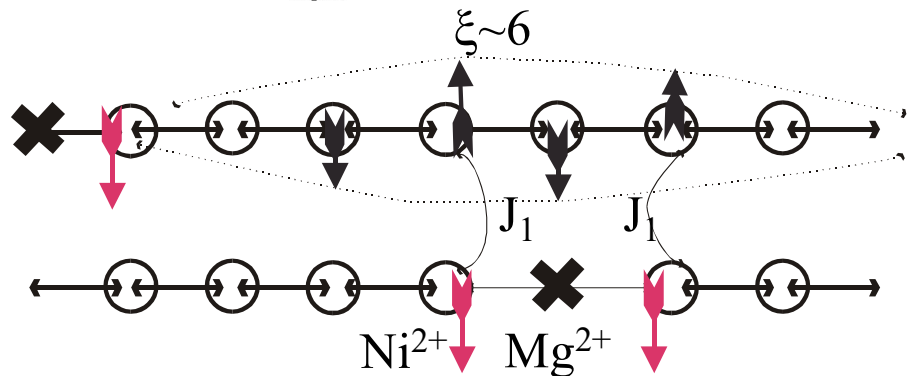
Refined magnetic moment is $0.9(1) \mu_B$ per Ni-site, at 2 K.

Local magnetic structure

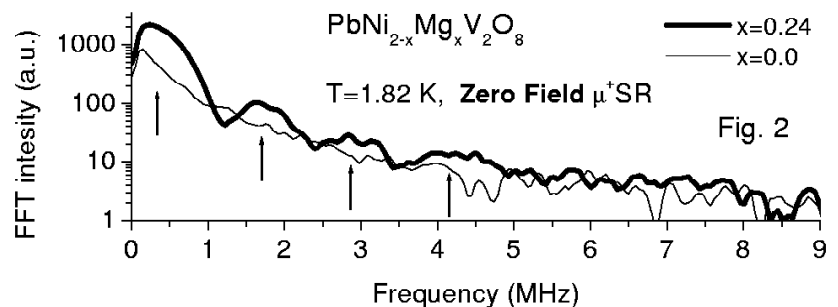
- $2 \times x (=0.48)$ spin-vacancies
- end-chain Ni can carry $[1.76/0.48] \times 0.9 = 3.3 \mu_B$.
- Spin- $1/2$ value $1.73 \mu_B$



Chain length $\langle L \rangle = (2/x) \cong 8$



ZF- μ^+ SR spectra was a convolution of more than one muon local fields (H_{int}). These internal fields take a number of values, $\langle H_{\text{int}} \rangle \sim 30, 120, 220$ G. /A.Lappas et al., unpublished (2001)/



Conclusions

1. Phase diagram $T_N(x)$ of $\text{PbNi}_{2-x}\text{Mg}_x\text{V}_2\text{O}_8$, $0.04 \leq x \leq 0.24$ is confirmed: T_N is increasing rapidly, rising to a maximum at an optimum Mg-doping level of $x=0.12$.
2. Neutron diffraction in the $\text{PbNi}_{1.76}\text{Mg}_{0.24}\text{V}_2\text{O}_8$ probes an average ordered moment size of $0.9(1) \mu_B$ per Ni-site, at 2 K. The moments are antiferromagnetically coupled below $T_N=3.2$ K, both in the chains along the c -axis and within the same (ab)-plane for Ni-sites on nearest-neighbour chains
3. Below T_N , the $x=0.24$ Mg-rich compound exhibits metamagnetic behaviour. $H_c(T=0\text{K})=1.2\text{T}$.

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Magndata $\text{PbNi}_{1.76}\text{Mg}_{0.24}\text{V}_2\text{O}_8$ (#0.1035) BNS, Bilbao